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## One-step Synthesis of Few-layer WS<sub>2</sub> by Pulsed Laser Deposition

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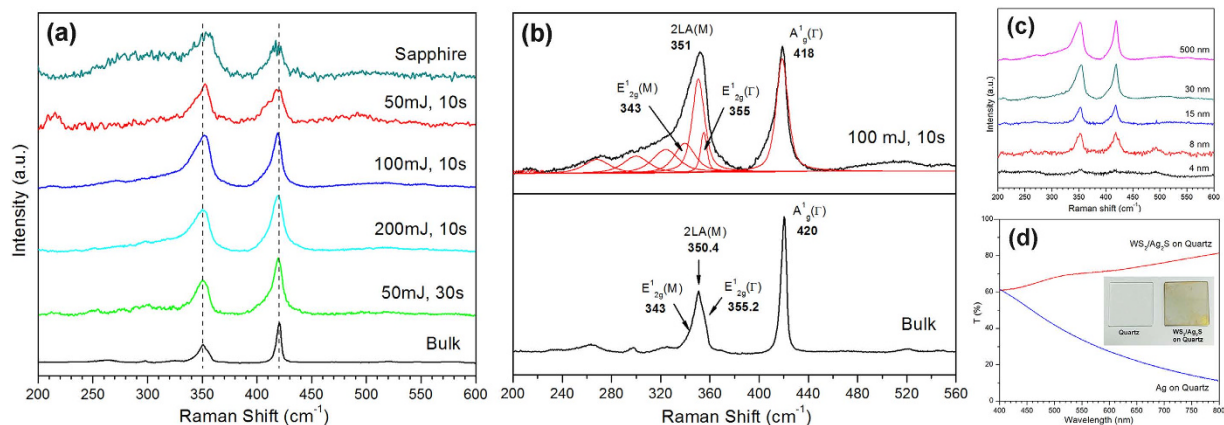
Atomically thin tungsten disulfide (WS<sub>2</sub>) has attracted much attention in recent years due its indirect-to-direct band gap transition, band gap tunability, and giant spin splitting. However, the fabrication of atomically thin WS<sub>2</sub> remains largely underdeveloped in comparison to its structural analogue MoS<sub>2</sub>. Here we report the direct fabrication of highly crystalline few-layer WS<sub>2</sub> on silver substrates by pulse laser deposition at the relatively low temperature of 450 °C. The growth takes places by conventional epitaxy, through the *in-situ* formation of nearly lattice-matching Ag<sub>2</sub>S on the silver surface. Intriguingly, it was observed that the resulting film was composed of not only the usual semiconducting 2H-WS<sub>2</sub> structure but also the less common metallic 1T-WS<sub>2</sub>. Modifications of the synthesis parameters allow for control over the crystalline quality, film thickness and crystal phase composition of the resulting WS<sub>2</sub> film.

Among the many materials that have the capacity to form 2-dimensional (2D) layers, much attention has been devoted to the lamellar transition metal dichalcogenides (TMDC)<sup>1–4</sup>. This class of materials, in which a plane of metal atoms are sandwiched between two planes of chalcogen atoms, exhibits interesting and diverse properties, ranging from insulators to semiconductors to metals<sup>2</sup>. As a typical example of a TMDC, tungsten disulphide (WS<sub>2</sub>) demonstrates a unique combination of structural, electronic, optical, mechanical, chemical, and thermal properties that have been the focus of various studies. Its electronic band gap for example, undergoes an indirect (1.4 eV) to direct (2 eV) transition when its size is reduced from bulk to a single layer<sup>2</sup>. The material has been used in fluorescent emitters<sup>5</sup> and field effect transistors<sup>6</sup>, for photovoltaics<sup>7</sup> and photocatalysis<sup>8</sup>, and is regarded as highly valuable for high power applications such as solid state batteries<sup>9</sup> and supercapacitors<sup>10</sup>. Monolayer WS<sub>2</sub> also possesses strong spin–orbit induced electronic band splitting and spin–valley coupling<sup>11,12</sup>, in addition to a band structure that can be tuned by variation in strain<sup>13</sup>. Until recently, most of the attention has been focused on the semiconducting phase of WS<sub>2</sub> and MoS<sub>2</sub>, both of which possess a prismatic coordination for the metal atom (2H). This 2H-structure is generally produced by physical exfoliation or vapour growth methods. In contrast, the metallic 1T-phase, with an octahedral coordination for the metal atom, has been synthesised primarily through lithium intercalation. Both 1T-WS<sub>2</sub> and 1T-MoS<sub>2</sub> layers have recently been shown to be very efficient hydrogen evolution electrocatalysts<sup>14–16</sup>.

Unlike related systems such as MoS<sub>2</sub>, techniques for effective preparation of 2D WS<sub>2</sub> remains largely underdeveloped. There are currently two main methods of preparing ultrathin WS<sub>2</sub>: (1) top-down exfoliation and (2) bottom-up substrate growth. Exfoliation can be mechanical<sup>17–18</sup>, chemically assisted<sup>19</sup> (e.g. sonication in a good solvent) or purely chemical<sup>14,20</sup> (intercalation with e.g. lithium), but the end results remain similar: high quality flakes can be prepared but there is very little control over their size, shape or nature of their edges. On the other hand, the substrate growth technique, commonly some form chemical vapour deposition<sup>21–23</sup> (CVD), is capable of producing large area monolayers with high crystallinity and good control over flake shape. However, many of these substrate growth techniques utilize different solid precursors heated to high temperatures in the range of 750–1000 °C and require long growth times. Although pulsed laser deposition (PLD) is also a type of bottom-up substrate growth technique, it is unique in that it is a purely physical method. Among its advantages is the ability to grow high quality films, to ablate any material and to obtain a stoichiometric transfer of target material onto the substrate, which is especially useful in the case of composite materials such as WS<sub>2</sub>. While PLD has found success in the fabrication of 2D materials such as few-layer graphene<sup>24</sup> and MoS<sub>2</sub><sup>25</sup>, it remains underdeveloped compared to exfoliation and CVD methods.

In this work, we describe the synthesis of highly crystalline few-layer WS<sub>2</sub> on Ag substrates using PLD at a relatively low temperature of 450 °C. Ag metal was selected as substrate for the ease of the ease of producing

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**Figure 1.** (a) Raman spectra of as-grown samples fabricated on Ag and sapphire. The  $\text{WS}_2$  film on sapphire was fabricated at 200 mJ and 10 s. The left and right dashed lines indicate the positions of the  $2\text{LA(M)}$  and  $\text{A}_{1g}^{\text{I}}$  phonon modes in bulk  $\text{WS}_2$  respectively. (b) Multi-peak Lorentzian fitting of Raman bands in the Ag sample fabricated at 100 mJ and 10 s, with bulk  $\text{WS}_2$  included for comparison. (c) Raman spectrum of as-deposited  $\text{WS}_2$  on quartz with varying Ag buffer thickness. All samples were synthesized using the parameters of 100 mJ and 10 s. (d) Optical transmittance of a quartz substrate with Ag buffer layer of 8 nm, before and after  $\text{WS}_2$  film growth. Inset: Photograph of the as-grown  $\text{WS}_2$  sample and bare quartz.

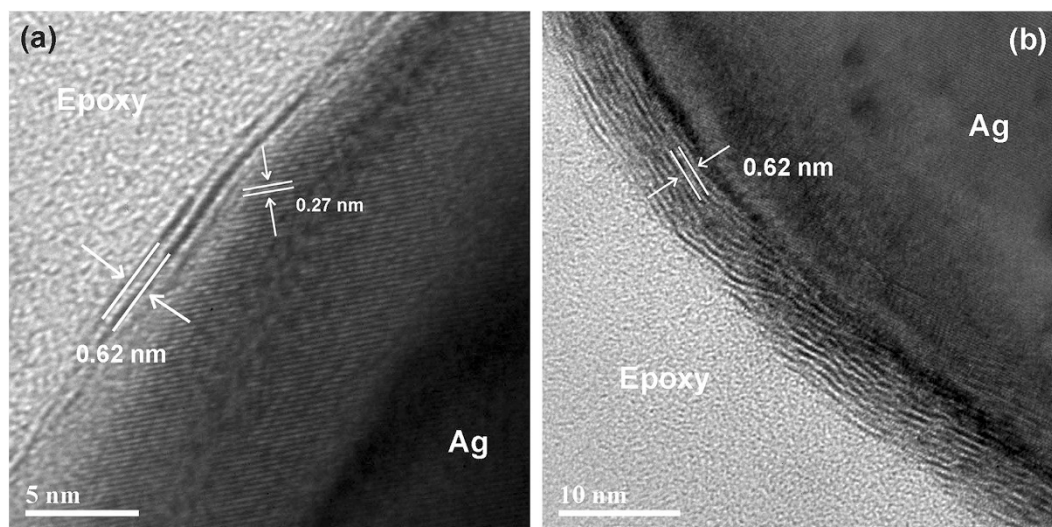
Laser Energy (mJ)	Deposition Time	$\text{A}_{1g}^{\text{I}}(\Gamma)$ ( $\text{cm}^{-1}$ )	$\text{E}_{2g}^{\text{I}}(\Gamma)$ ( $\text{cm}^{-1}$ )	$2\text{LA(M)}$ ( $\text{cm}^{-1}$ )	$\frac{I_{2\text{LA}}}{I_{\text{A}_{1g}^{\text{I}}}}$	Estimated Number of Layers
50	10	418	351	355	0.94	~2
100	10	418.7	350.6	355	0.75	~3
200	10	419.1	350.1	354.9	0.58	4–5
50	30	419.5	350.2	355	0.51	>5
Bulk	–	420	350.4	355.2	0.46	Bulk

**Table 1.** Peak position for the Raman modes  $\text{A}_{1g}^{\text{I}}(\Gamma)$ ,  $2\text{LA(M)}$  and  $\text{E}_{2g}^{\text{I}}(\Gamma)$ , as well as the intensity ratio of  $\text{A}_{1g}^{\text{I}}(\Gamma)$  to  $2\text{LA(M)}$  for the as-deposited samples on Ag.

metal-semiconductor contacts without film transfer. Additionally, the introduction of a metal support can substantially alter the H binding energy of 2D TMDCs<sup>26</sup>, enabled by charge transfer from the substrate to the overlying film and through strong interactions at the interface. Metal supported  $\text{WS}_2$  could thus potentially be used as novel catalysts for hydrogen production. Surprisingly, using PLD to grow  $\text{WS}_2$  does not produce the expected semiconducting prismatic 2H structure. Instead, it creates metallic  $\text{WS}_2$  films with the distorted octahedral 1T- $\text{WS}_2$  structure. The structure and properties of these films were explored through a combination of Raman spectroscopy, photoluminescence (PL) measurement, x-ray photoelectron spectroscopy (XPS), high-resolution transmission electron microscopy (HRTEM), and powder x-ray diffraction (PXRD).

## Results and Discussion

Figure 1a shows the Raman spectra of the as-grown  $\text{WS}_2$  film on Ag excited under ambient conditions, with the spectrum of bulk  $\text{WS}_2$  and a  $\text{WS}_2$  sample grown on insulating sapphire substrates included for comparison. The spectrum of the sapphire sample was normalized to remove the peaks from the substrate. At this excitation wavelength of 514.5 nm, the spectrum reveals many second-order Raman peaks, in addition to the first-order phonon modes. The strongest peak at  $\sim 352 \text{ cm}^{-1}$  can be resolved by multi-peak Lorentzian fitting into three individual contributions at 343, 351, and  $355 \text{ cm}^{-1}$  as shown in Fig. 1b for the sample fabricated at 100 mJ and 10 seconds. These modes are assigned to the in-plane vibrational  $\text{E}_{2g}^{\text{I}}(\text{M})$  mode, the second-order mode of longitudinal acoustic phonon  $2\text{LA(M)}$ , and the in-plane vibrational  $\text{E}_{2g}^{\text{I}}(\Gamma)$  mode respectively<sup>12</sup>. The other peaks at 418 and  $343 \text{ cm}^{-1}$  are attributed to the out-of-plane  $\text{A}_{1g}^{\text{I}}$  mode and the in-plane vibrational  $\text{E}_{2g}^{\text{I}}(\text{M})$  mode. Studies have previously shown that Raman characterization can provide unambiguous and nondestructive identification of the thickness of  $\text{WS}_2$ ; the  $\text{A}_{1g}^{\text{I}}(\Gamma)$  mode softens while both the  $2\text{LA(M)}$  and  $\text{E}_{2g}^{\text{I}}(\Gamma)$  modes present a subtle red-shift with a decreasing number of layers<sup>5,27</sup>. In particular, at 514.5 nm laser excitation, the  $\text{WS}_2$  spectrum reveals a striking increase in the intensity ratio of the  $2\text{LA(M)}$  to  $\text{A}_{1g}^{\text{I}}(\Gamma)$  phonon modes due to a double resonance process<sup>27</sup>. Table 1 summarizes the frequency for the three main Raman modes  $\text{A}_{1g}^{\text{I}}(\Gamma)$ ,  $2\text{LA(M)}$  and  $\text{E}_{2g}^{\text{I}}(\Gamma)$ , as well as the intensity ratio for the two most intense peaks in our pulsed laser fabricated samples on Ag. Based on these values, it can be concluded that atomically thin  $\text{WS}_2$  films ( $\leq 5$  layers) can be formed at any laser energy within the range of 50–200 mJ as long as the deposition time was capped at a maximum of 10 seconds. Compared to chemically derived  $\text{WS}_2$ <sup>21,23</sup>, our samples exhibited Raman peaks with broader full width at half maximums (FWHM), such that it becomes difficult to distinguish the phonon modes in the  $260\text{--}330 \text{ cm}^{-1}$  range without multi-peak Lorentzian fitting. This



**Figure 2.** Cross-section TEM of Ag samples fabricated at 50 mJ laser energy and a deposition time of (a) 10 seconds and (b) 30 seconds.

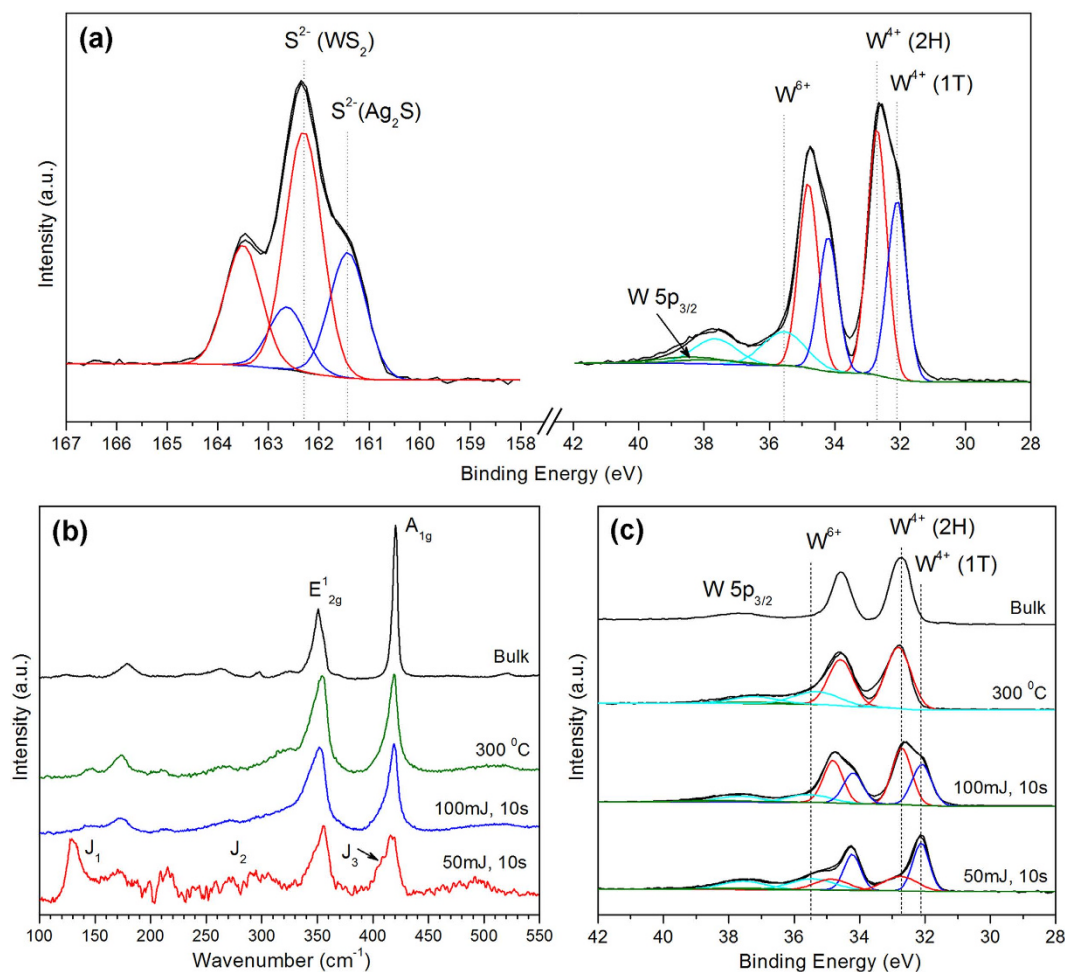
indicates the presence of defects arising from disorder in the atomic arrangement of the WS<sub>2</sub> film. In addition, the samples fabricated at a laser energy of 50 mJ demonstrate much lower signal intensity, possibly indicating only partial crystallization or a large concentration of defects. As such, even though the thinnest films are formed at the lowest energy, for an optimal balance between film thickness and crystalline quality, a laser energy of 100 mJ would be most ideal to grow highly crystalline few-layer WS<sub>2</sub>.

Although insulating substrates such as sapphire are preferred substrates for growing 2D materials for characterization of electronic properties, it was found that such substrates do not produce good quality WS<sub>2</sub> films in pulsed laser synthesis. Figure 1a shows that despite the use of higher laser energies of 200 mJ to fabricate the sapphire sample, its Raman spectrum resembles the partially crystalline and highly defective spectrum of the Ag sample fabricated at 50 mJ and 10 s. At lower energies of 100 and 50 mJ, the characteristic peaks of 2H-WS<sub>2</sub> are completely absent. This obstacle can be however, overcome by sputtering a very thin layer of the Ag buffer on the substrate. Our investigations reveal that the minimum thickness that the Ag layer can have before it begins to significantly affect the crystalline quality of the overlying WS<sub>2</sub> film is ~8 nm (Fig. 1c). These WS<sub>2</sub> films were grown on quartz, which not only permits the direct synthesis of crystalline few-layer WS<sub>2</sub> on insulating substrates without the need for transfer techniques, but also allows the fabrication of samples that are capable of transmitting visible light. Figure 1d shows that the triple-layered WS<sub>2</sub> film on quartz has an average optical transmittance of 71% in the wavelength range between 400 and 800 nm. Reducing the buffer layer to such minimal thickness also results in *in-situ* consumption of most, if not all, of the pure Ag metal to form Ag<sub>2</sub>S (a phase which was detected in the TEM, XPS and XRD analysis). This was deduced from the increase in optical transmission in the sample after deposition of WS<sub>2</sub>, also shown in Fig. 1d. The high optical transmittance of our WS<sub>2</sub> films on quartz means that they are suitable for solar energy applications. The presence of the Ag<sub>2</sub>S phase is also beneficial as the material appears to be a promising solar absorbing material with its narrow band gap of ~0.9 eV<sup>28</sup> and its unique combination of properties such as light absorbance in the near-infrared spectral regions<sup>29</sup>.

High resolution transmission electron microscope (TEM) images of our pulsed laser fabricated samples are shown in Fig. 2, and reveal the stacking of WS<sub>2</sub> (002) layers with an interplanar spacing of 0.62 nm on top of silver. Figure 3a provides direct evidence of the successful formation of double-layered WS<sub>2</sub> in the sample fabricated at 50 mJ and 10 seconds, with the first WS<sub>2</sub> layer forming covalent bonds to the previously mentioned Ag<sub>2</sub>S phase that develops *in-situ* on the surface of the Ag buffer. This bonding between the two layers manifests as an indistinct boundary at the interface contrary to the sharp interfaces of 2D layered materials grown by van der Waals epitaxy. The formation of the Ag<sub>2</sub>S phase is supported by the presence of lattice fringes with interplanar spacings of 0.27 nm located underneath the WS<sub>2</sub> film that can be ascribed to the (120) plane of monoclinic acanthite Ag<sub>2</sub>S [JCPDS #14-0072]. This Ag<sub>2</sub>S phase was also previously found to similarly promote the growth of ultrathin MoS<sub>2</sub> by PLD through lattice matching and conventional epitaxy<sup>25</sup>. Figure 3b shows the change in the WS<sub>2</sub> film when the deposition time is increased to 30 seconds. The number of layers have increased to >5, verifying the results of the Raman measurements.

To characterize the chemical nature and bonding state of WS<sub>2</sub> on Ag metal, X-ray photoelectron spectroscopy (XPS) was employed. Figure 3a depicts the W 4f and S 2p core level XPS scans for the WS<sub>2</sub> film on silver. Two doublets are present in the S 2p spectra, one occurring at 162.3 and 163.5 eV, consistent with the S<sup>2-</sup> species of WS<sub>2</sub>, while the second pair located at 161.1 and 162.4 eV can be assigned to the S 2p<sub>3/2</sub> and S 2p<sub>1/2</sub> peaks of Ag<sub>2</sub>S. The presence of both doublets confirms the successful growth of WS<sub>2</sub> as well as the *in-situ* formation of Ag<sub>2</sub>S in the samples. For the W 4f spectra, there is a small shoulder at 35.6 eV corresponding to the W<sup>6+</sup> state that shows the formation of tungsten oxide in the as-deposited films. This oxide can arise in part from surface oxidation of the WS<sub>2</sub> target, and in part from unreacted W atoms remaining on the substrate surface after completion of the WS<sub>2</sub> film. These unreacted W atoms are always produced because some S atoms are initially consumed in the making

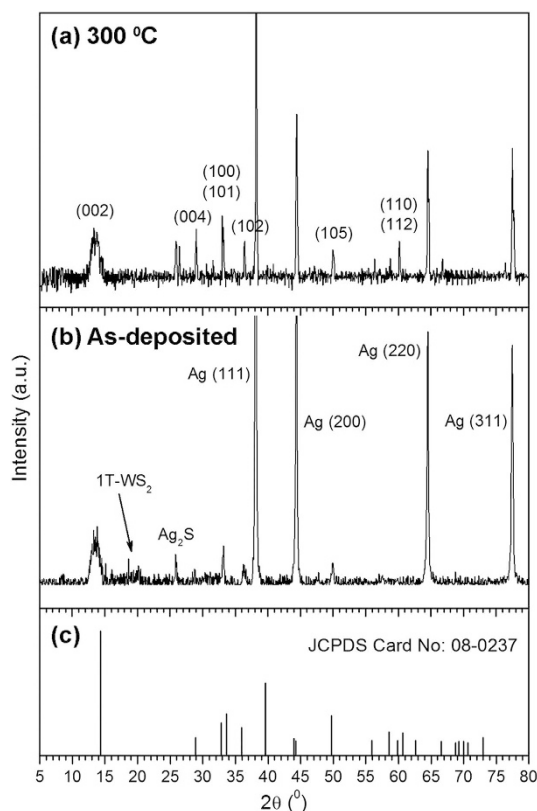




**Figure 3.** (a) XPS spectra showing W 4f and S 2p core level peak regions for as-deposited samples fabricated using 100 mJ laser energy and 10 s deposition time. (b) Raman spectra of  $WS_2$  films deposited on Ag. The  $J_1$ ,  $J_2$ , and  $J_3$  peaks are only active in the as-deposited sample fabricated at 50 mJ and 10 s. The spectrum for  $WS_2$  films annealed at 300 °C more closely resembles that of bulk 2H- $WS_2$ . (c) XPS spectra showing the W 4f core level peak regions of as-deposited, annealed and bulk samples. W 4f peaks were deconvoluted with 2H (red) and 1T (blue) components. As-deposited  $WS_2$  sheets at 50 and 100 mJ have 1T-to-2H phase ratio of 1.67 and 0.73 respectively, whereas after annealing at 300 °C the material is converted to purely 2H- $WS_2$ .

of the  $Ag_2S$  phase. The oxide is believed to form sort of capping layer over the  $WS_2$  film when it is first exposed to air, and does not negatively affect the quality the film itself as the calculated stoichiometric ratio of S : W is very close to the ideal value of 2. It was also observed that the main tungsten doublet peak can be deconvoluted into two separate pairs. The second doublet (red curve) occurs at binding energies of 32.7 and 34.8 eV, which corresponds well with the  $W^{4+}$  species of highly crystalline 2H- $WS_2$ <sup>30</sup>. The doublet at 32.1 and 34.2 eV (blue curve) on the other hand lies between the binding energies of metallic tungsten ( $W^{0+}$ ) and  $W^{4+}$ , and appears to be due to a partially sulfided, intermediate  $W^{x+}$  state. However, calculation of the stoichiometric ratio of S : W atoms reveal that both W tungsten species (red and blue curves) have a S : W ratio of ~2, indicating that the W 4f doublet located at 32.1 and 34.2 eV is also due to a  $W^{4+}$  state. The negative shift of peak binding energies by 0.6 eV is thus believed to be attributed to the formation of the 1T-phase of  $WS_2$ , consistent with its metallic nature, and is comparable to with previous studies on 1T-MX<sub>2</sub> materials<sup>14,31</sup>.

The formation of the 1T-phase can also be established by closer inspection of the Raman spectra, as well as by powder X-ray diffraction (XRD). As noted by Voiry *et al*<sup>14</sup>, the Raman spectra of 1T- $WS_2$  display the additional modes of  $J_1$ ,  $J_2$  and  $J_3$  that are attributed to the superlattice structure of the distorted 1T-phase. However, these Raman modes are only readily observable when the concentration of 1T- $WS_2$  dominates over the 2H-phase. Figure 3b shows the Raman results from two of our as-deposited samples with widely differing concentrations of 1T- $WS_2$ . It is apparent that the  $J_1$ ,  $J_2$  and  $J_3$  modes are noticeable only in the sample fabricated at 50 mJ and 10 s due to the much higher 1T-to-2H ratio of 1.67. In contrast, the sample fabricated at 100 mJ and 10 s only has a 1T-to-2H ratio of 0.73, and correspondingly the typical Raman peaks of 1T- $WS_2$  have almost completely vanished. As-deposited samples were also annealed at varying temperatures to test the stability of the 1T-phase. Included for comparison in Fig. 3b is a sample annealed at 300 °C for 30 minutes, wherein a complete phase transformation from 1T to 2H- $WS_2$  was achieved as revealed by XPS (Fig. 3c). The Raman bands of 2H- $WS_2$  in this sample appear more

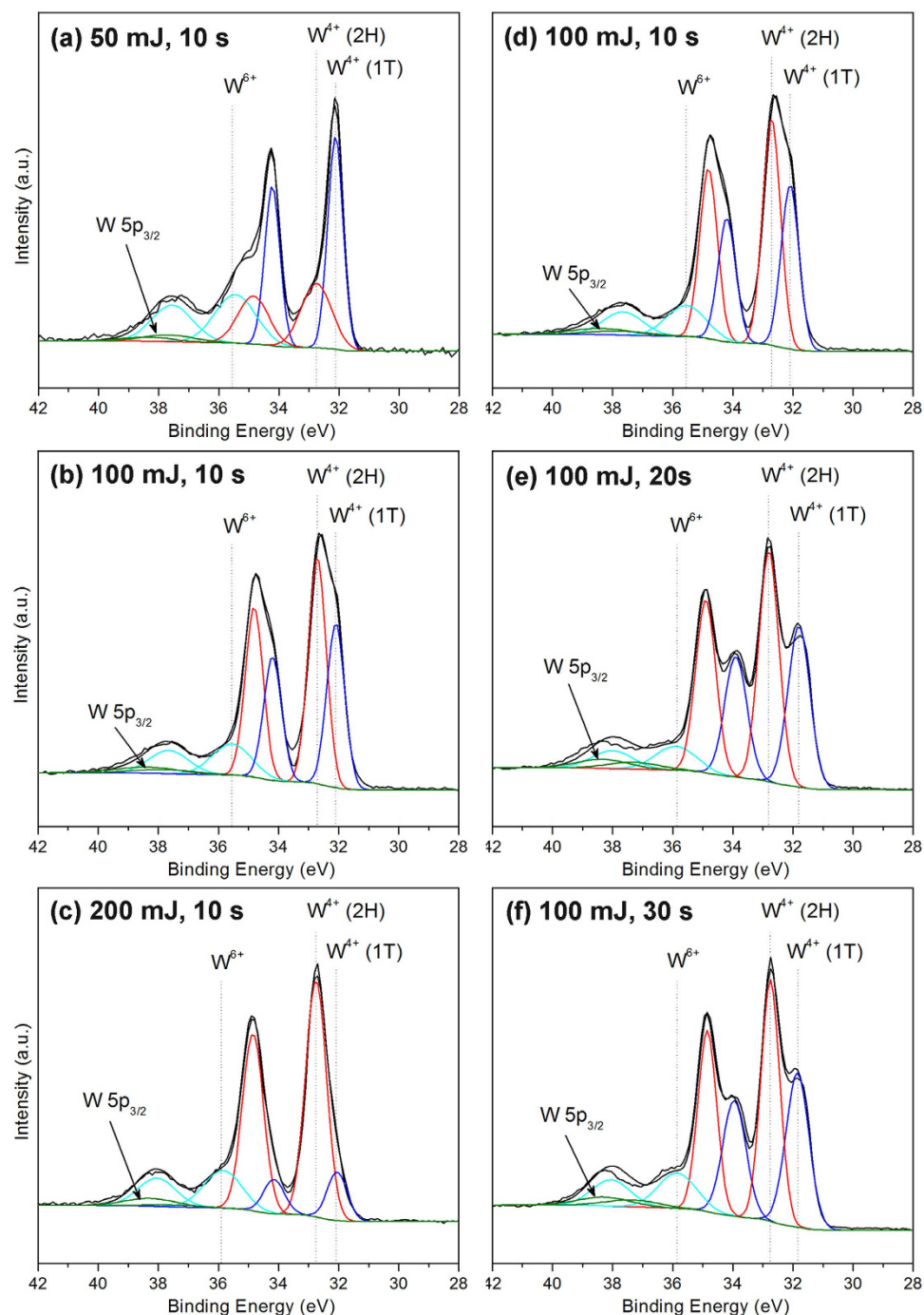


**Figure 4.** Powder X-ray diffractograms of WS<sub>2</sub> nanosheets (a) annealed at 300 °C and (b) as-deposited at 50 mJ and 10 s. (c) Bulk diffraction peaks of 2H-WS<sub>2</sub>.

distinct; in particular the combination mode of  $2\text{LA-E}_{2g}^2$  at  $\sim 320\text{ cm}^{-1}$  becomes visibly noticeable. Nonetheless, the peaks still lack the sharpness of CVD grown WS<sub>2</sub> sheets, and there is only a slight increase in the peak intensities while the FWHMs remained largely unchanged. This indicates that there is some degree of atomic disorder present even in the pure 2H-phase, and the crystalline quality of PLD grown WS<sub>2</sub> on Ag is not quite comparable to CVD grown WS<sub>2</sub> layers. The fact that the 1T to 2H phase transition can be induced with thermal annealing indicates that the 1T-WS<sub>2</sub> produced in this work is metastable. Indeed, it was observed that partial 1T to 2H transition begins to occur even at temperatures as low as 100 °C, which would be detrimental if the 1T-phase was intended for use at high temperatures. The ability to prepare metastable 1T-WS<sub>2</sub> layers at temperatures higher than the 1T to 2H transition is rather surprising. However, this is likely to be simply due to the short deposition times and fast cooling rate, such that the 1T-structure does not have sufficient time to relax to the more stable 2H-phase.

Unlike with 2H-WS<sub>2</sub>, characterization of 1T-WS<sub>2</sub> by powder XRD is rather challenging because the 1T-phase does not exhibit a well-defined crystalline structure and well-resolved diffraction peaks<sup>31</sup>. This problem is exacerbated by the fact that our samples are all mixtures of 1T and 2H-WS<sub>2</sub>. Nonetheless, we were still able to observe differences in the XRD diffractograms obtained before and after annealing. As shown in Fig. 4, both the as-deposited and annealed samples have peaks corresponding to the known 2H-WS<sub>2</sub> pattern. In the as-deposited sample, the presence of a broad peak at  $14.32^\circ$  with the highest intensity reveals the preferential growth of WS<sub>2</sub> sheets along the (002) direction. In contrast, the XRD pattern of the annealed sample shows a more intense (100) peak, which suggests the growth of protrusion edges along the (100) direction. Compared with 2H-WS<sub>2</sub>, the (002) peak of both the as-deposited and annealed sample is shifted to lower  $2\theta$  values, indicating a lattice expansion of (002) layers, i.e. 0.62 nm compared to 0.616 nm. The annealed sample also presents additional peaks corresponding to 2H-WS<sub>2</sub>, suggesting an improvement in crystallinity of the WS<sub>2</sub> film that may be attributed to a 1T to 2H-phase transformation. There is furthermore a very minor peak around  $19\text{--}20^\circ$  that vanishes after thermal annealing. As this peak cannot be indexed to either 2H-WS<sub>2</sub>, Ag, Ag<sub>2</sub>S or even WO<sub>3</sub>, and has previously been observed to be present in 1T-WS<sub>2</sub><sup>31</sup>, we believe that it can be attributed to the presence of the 1T-phase.

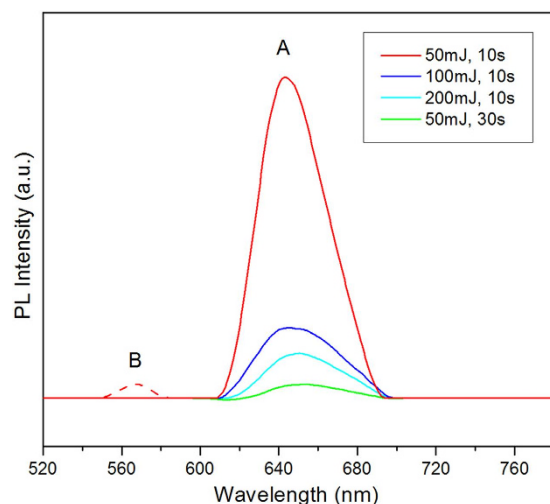
It was observed that the concentration of the 1T-WS<sub>2</sub> phase appears to decrease with longer deposition times and higher laser energies (Fig. 5). In addition, the peak position of the W 4f doublet for 1T-WS<sub>2</sub> continues to shift to lower binding energies with increasing deposition times, an indication of increasingly metallic character. This phenomenon could possibly be due to bombardment of the growing film by energetic species in the pulsed laser deposition process, which may lead to the breaking of W-S bonds in 1T-WS<sub>2</sub>. It is likely that with longer deposition times, the 1T-WS<sub>2</sub> phase would convert completely into metallic tungsten. Similarly, the formation of a metallic layer was previously observed during prolonged sputtering of WS<sub>2</sub><sup>32</sup>, another ion-assisted deposition technique. The breaking of the W-S bond appears to affect only 1T-WS<sub>2</sub> and not 2H-WS<sub>2</sub>, as seen by the negligible negative shift in the W 4f peak position for the 2H polymorph. This is deduced to be a consequence of the higher



**Figure 5.** XPS spectra showing W 4f core level peak regions for Ag samples fabricated at different laser energies and deposition times.

thermodynamic stability of 2H-WS<sub>2</sub>, rendering it less likely to be affected by the bombardment of energetic species during film growth.

The presence of the 1T-phase also affects the photoluminescence (PL) measurements of our samples. Weak PL was observed in as-deposited WS<sub>2</sub> nanosheets, as expected from their partial metallic character. Enlarging the emission spectra allows us to see that the sample fabricated at 50 mJ and 10 s presents a major peak at ~645 nm and a shoulder peak at ~570 nm (Fig. 6), corresponding to A and B excitonic transitions of the K point of the Brillouin zone. Curiously, the expected intensity and position dependence of the major PL peak with layer thickness can be observed even among the as-deposited samples despite the varying concentrations of 1T-WS<sub>2</sub>. As shown in Fig. 6, the major PL peak at ~640 nm quickly diminishes in intensity and gradually red-shifts with increasing layer thickness, a trend that has been observed by other groups<sup>12,22,33</sup>. This is unusual as the thinnest sample also has the highest concentration of 1T-WS<sub>2</sub>, which is expected to inhibit PL, and yet this sample gives the highest emission intensity. Because PL originates near the surface of a material, and is sensitive enough to be affected by surface



**Figure 6.** Photoluminescence spectra of as-grown samples fabricated on Ag substrates at  $\lambda_{\text{ex}} = 488$  nm.

adsorbates, the fact that the PL spectra of as-deposited samples appears to be largely unaffected by the presence of 1T-WS<sub>2</sub> suggests that the 1T-phase is located away from the surface of the film. While the origin of 1T-WS<sub>2</sub> in our samples is as yet unclear, we believe that it can be attributed to the capacity of Ag atoms to act as an electron donor for W (owing to Ag having more valence electrons than W). The stabilization mechanism is therefore similar to chemically exfoliated MoS<sub>2</sub> and WS<sub>2</sub>, wherein the 1T-phase is stabilized by substitutional doping of an electron donating atom. In such a case, the 1T-phase in our samples would be concentrated at the interface while 2H-WS<sub>2</sub> would dominate at the surface of the film, which could explain their unexpected PL results. Further experimental and theoretical studies would be required to verify these hypotheses.

From the above results, it appears that Ag is indeed capable of stabilizing the 1T-phase of WS<sub>2</sub>, which is noteworthy as recent studies have shown that the 2H and 1T phase have matching lattices and coherent interfaces between domains of the two phases can be formed<sup>34</sup>. As 2H-WS<sub>2</sub> is a semiconductor and 1T-WS<sub>2</sub> a metal, localized phase stabilization of the 1T phase to form a 2H-1T hybrid structure presents a viable route to achieving unique electronic heterojunctions across a chemically homogeneous layer, which is advantageous for the fabrication of molecular electronic devices. Furthermore, studies have revealed that the 1T-WS<sub>2</sub> phase possesses higher reactivity and catalytic ability compared to the 2H-phase<sup>14,22</sup>. With Ag also being catalytically active and considered a low-cost alternative to Pt<sup>35</sup>, our findings open up the possibility of Ag stabilized 1T-WS<sub>2</sub> nanostructures as superior catalysts for the hydrogen evolution reaction. Even so, some challenges remain, not least of which is the impact of the Ag<sub>2</sub>S phase on the intrinsic property of the WS<sub>2</sub> film. Due to the covalent bonds holding WS<sub>2</sub> to the buffer layer, it is unlikely that the Ag<sub>2</sub>S phase can be removed without damaging the crystal structure of the WS<sub>2</sub> sheets. Thus any effort to incorporate pulsed laser deposited WS<sub>2</sub> into electronic devices must take into account the presence of this Ag<sub>2</sub>S phase, and further fundamental studies on the exact electronic properties of the WS<sub>2</sub>/Ag<sub>2</sub>S material are necessary.

In summary, we have successfully fabricated crystalline few-layer WS<sub>2</sub> on Ag by pulsed laser deposition. The growth process involves the *in-situ* formation of a lattice matching Ag<sub>2</sub>S phase and that promotes the formation of crystalline WS<sub>2</sub> through conventional epitaxy rather than expected Van der Waals epitaxy. Samples fabricated at a laser energy of 100 mJ and a deposition time of 10 seconds demonstrated the optimum balance between crystalline quality and film thickness. It was also found that crystalline WS<sub>2</sub> is able to form on Ag layers as thin as 8 nm, producing samples with high optical transmittance in the visible range, which is advantageous for applications in optoelectronic devices and solar cells. In addition, the Ag substrate was observed to stabilize the 1T-polymorph of WS<sub>2</sub>, resulting in a hybrid 2H-1T structure. Further work would however be required to fully understand the factors behind the stabilization of the 1T phase on Ag before such controlled fabrication of such hybrid or even purely 1T-WS<sub>2</sub> structures can be achieved.

## Methods

**Synthesis.** Ag films of approximately 500 nm thickness were sputtered onto normal doped (n+) Si substrates and loaded in a KrF ( $\lambda = 248$  nm) Lambda Physik excimer PLD system. During deposition, a WS<sub>2</sub> target, 99.9% (Inlab Supplies) was ablated at a laser frequency of 10 Hz for a pulse duration of 25 ns in a vacuum environment of  $2 \times 10^{-6}$  Torr. The target was rotated at a speed of 6 rpm, with the laser (spot size 1 mm<sup>2</sup>) ablating a circular outline of 2 cm in radius in order obtain a uniform film. The laser energies used for ablation was kept within the range of 50–200 mJ, and the deposition time between 10–30 seconds. Substrate temperature was kept at 450 °C during deposition, and subsequently decreased at a controlled rate of 20 °C/min until the temperature reached 300 °C, thereupon natural cooling processes took over. The fabrication temperature of 450 °C was selected as it is the optimum for growth of ultrathin WS<sub>2</sub> films by PLD; any lower and the crystalline quality decreases because the atoms do not have sufficient thermal energy to rearrange themselves into a periodic configuration, higher temperatures only provide minimal increase in crystallinity, whereas temperatures >700 °C promotes the reaction of tungsten with oxygen.



**Characterization.** WS<sub>2</sub> films were characterized by Raman spectroscopy performed in a Horiba MicroRaman HR Evolution System using an Argon laser beam with an excitation wavelength of 514.5 nm.

The phonon mode from silicon at 520.6 cm<sup>-1</sup> was used for calibration. The Raman spectrum for bulk WS<sub>2</sub> was collected from the WS<sub>2</sub> target used. PL measurements were taken in a Perkin-Elmer fluorescence spectrometer LS 55 with excitation wavelength of 488 nm. A JEOL JEM-2010F high resolution transmission electron microscopy (TEM) operated at 200 kV was used to obtain bright field cross-section images of the samples. Surface composition was analyzed by x-ray photoelectron spectroscopy (XPS) using a Kratos Analytical Axis Ultra<sup>DLD</sup> UHV spectrometer with a monochromatized Al K $\alpha$  x-ray source (1486.6 eV) scanning a spot size of 700  $\mu$ m by 300  $\mu$ m. Core-level XPS spectra were obtained by photoelectrons at a take-off angle of 90°, measured with respect to the sample surface at a vacuum of  $5 \times 10^{-9}$  Torr. The pass energy for the narrow scans were set at 10 eV.

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### Author Contributions

T.A.J.L., D.H.C.C. and A.T.S.W. designed the experiments. The synthesis and characterization of samples was carried by T.A.J.L. All authors contributed to the manuscript preparation and discussion of results.

### Additional Information

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