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SUBJECT AREAS:
PHOTOCATALYSIS
SOLAR CELLS
QUANTUM DOTS

Towards Visible Light Hydrogen Generation: Quantum Dot-Sensitization via Efficient Light Harvesting of Hybrid-TiO₂

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Received
23 September 2013

Accepted
8 November 2013

Published
25 November 2013

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We report pronounced enhancement of photoelectrochemical hydrogen generation of a quantum dot-sensitized hybrid-TiO₂ (QD/H-TiO₂) electrode that is composed of a mesoporous TiO₂ layer sandwiched by a double sided energy harvesting layer consisting of a surface-textured TiO₂ inverse opals layer on the bottom and a patterned mesoporous TiO₂ layer on the top. CdSe/H-TiO₂ exhibits a maximum photocurrent density of ~ 16.2 mA/cm², which is 35% higher than that of the optimized control sample (CdSe/P25), achieved by matching of the bandgap of quantum dot-sensitization with the wavelength where light harvesting of H-TiO₂ is observed. Furthermore, CdSe/H-TiO₂ under filtered exposure conditions recorded current density of ~ 14.2 mA/cm², the greatest value in the visible range. The excellent performance of the quantum dot-sensitized H-TiO₂ suggests that alteration of the photoelectrodes to suitable nanostructures with excellent light absorption may offer optimal strategies for attaining maximum efficiency in a variety of photoconversion systems.

Photoelectrochemical (PEC) cells are useful devices designed for the chemical process of generating hydrogen via electrolysis of water, namely water splitting, under irradiation of a broad spectrum of sunlight^{1–4}. In general, PEC devices are composed of photoelectrodes, where oxidation/reduction of water occurs with response to sunlight, and appropriate aqueous electrolytes containing a redox couple^{5–7}. TiO₂ has proven to be a promising photoelectrode in photoconversion systems due to its flat band potential, fast electron transfer mechanism, and photoelectrochemical stability in most aqueous solutions^{8–10}. However, the use of TiO₂ as a single electrode material is fundamentally hampered by its large bandgap (3.2 eV), thus resulting in truncated and inclusively low efficiency in the visible region, which accounts for almost half of the entire sunlight spectrum. Accordingly, extending the light absorption window of the photoelectrode to visible light is considered critical to increase the efficiency of photoconversion^{11–13}. Notable efforts to lower the bandgap of TiO₂ include doping of heteroatoms^{14–16} or coating of metal or semiconductor particles utilizing excitation of surface plasmons^{17–19} or quantum dot-sensitization^{3,20,21}, permitting excitation of electrons in the photoelectrodes at energies much below TiO₂'s original energy bandgap^{22–24}. Among these approaches, quantum dot-sensitization has drawn considerable attention since it has shown an enhanced interfacial charge transfer mechanism via multiple excitations by single photon absorption.

Light harvesting refers to enhancement of absorbed light via either photon trapping^{25–27} or anti-reflection (AR)^{28–30}, leading to an increase in the fraction of available light^{31,32}. It has been reported that nanostructures can exhibit excellent absorption properties in the visible region via light harvesting, when their dimensions are comparable to the wavelength of light³³. By enhancing absorption in the visible range via light harvesting, it becomes possible to more effectively sensitize quantum dots with a bandgap in the range of visible light, leading to maximal power conversion efficiency of quantum dot-based PEC cells.

Mesoporous TiO₂ nanostructures fabricated by P25 with a mean diameter of about 25 nm have been widely used for fabrication of the working electrode in dye-sensitized solar systems, a representative photoconversion system, due to their advantage of large surface area for adsorption of dye molecules. However, the small particles are not able to efficiently harvest light in the red and near-infrared regions of the light spectrum and they also reflect a large amount of light due to the high refractive index of TiO₂. Proper modification of the P25

photoelectrode is thus needed and careful design of quantum dot-sensitized devices based on a light harvesting mechanism is crucial to maximize the efficiency in photoconversion systems.

Here, we report an optimized quantum dot-sensitized hybrid-TiO₂ (QD/H-TiO₂) electrode sandwiched with light harvesting layers. The H-TiO₂ electrode is constructed using a 4.5 μm -thick mesoporous TiO₂ film made of crystalline nanoparticles (25 nm) and additionally a 500 nm-thick patterned mesoporous TiO₂ layer on the top and 1 μm -thick surface-textured TIO (st-TIO) layer on the bottom, both with diameters on the order of hundreds of nanometers, with the capability of AR and photon trapping, respectively, contributing to maximized light harvesting. The configuration of the light harvesting layer sitting on both the bottom and top of the mesoporous TiO₂ layer provides devices with similar mechanism to mirrors in a laser, where entrapped light is enhanced by repeated reflection at both ends and thus maximizes the intensity of the absorbed light at a wider solar spectrum. Furthermore, CdSe/H-TiO₂ greatly broadens the range of photoconversion below the band gap of H-TiO₂ by enhanced quantum dot-sensitized electron transport substantially supported by the light harvesting mechanism of H-TiO₂. Under AM 1.5 G simulated sunlight illumination, the maximum photocurrent density of CdSe/H-TiO₂ reaches $\sim 16.2 \text{ mA/cm}^2$, which is 35% higher than that of the optimized control sample (CdSe/P25), at an applied bias of 0.5 V versus Ag/AgCl. More surprisingly, the photocurrent from hydrogen generation of CdSe/H-TiO₂ solely by the visible spectrum reaches $\sim 14.2 \text{ mA/cm}^2$, which is the highest record value in the visible range and almost 88% of that obtained under UV-visible irradiation. The greatly improved performance of CdSe/H-TiO₂, which is attributable to the ability of light harvesting of H-TiO₂, suggests the strong need for engineering light in photoconversion systems.

Results

H-TiO₂ was fabricated by the following four steps, as shown in Figure 1a. First, surface-activated polystyrene beads (PS) obtained by treatment with Triton X-100 were self-assembled using Colvin's method³⁴ into a few layers of a hexagonal structure on a FTO substrate with a thickness of 1 μm .

PS structures were infiltrated with a TiCl₄ aqueous solution and annealed in an air condition, inducing gelation of a TiO₂ sol and calcination of PS simultaneously, and thus forming the TIO structure (I). Second, a surface textured titanium inverse opal (st-TIO, II) coated with a diameter of 7 nm TiO₂ mesoporous structures was created by selective removal of triblock copolymer domains from the coated layer containing TiO₂ precursors and block copolymers on TIO. In principle, the st-TIO has been designed to minimize the loss of surface area for reaction sites while achieving a photon trapping effect. Third, P25 mesoporous structures (III) were deposited onto the st-TIO with a thickness of 5 μm . Fourth, a top-patterned mesoporous P25 nanostructure (p-P25, IV) was created by imprinting using a 500 nm-thick PDMS mold with 500 nm diameter and 600 nm pitch square patterns followed by the sintering process. The targeted structure is H-TiO₂, composed of a 4.5 μm -thick mesoporous TiO₂ layer with a double sided energy harvesting layer comprising a 1 μm -thick st-TIO layer of a few hundred nanometer scale on the bottom and a 500 nm thick mesoporous TiO₂ square pattern on the top, giving a total thickness of 6 μm . Figure 1b is a scanning electron microscopy (SEM) image of the st-TIO structure (II) with air holes with a diameter of 290 nm. The dimension of the TIO was determined so as to efficiently scatter light for maximized light trapping as optimized in our previous reports^{29,33,35}. Figure 1c is a SEM image of patterned mesoporous TiO₂ structures obtained by stamping a P25 layer using the PDMS molds followed by final sintering. The dimensions were fixed for the maximized AR effects, based on the previous reports^{28,29}. The volume fraction of TiO₂ in the AR layer is found to be 0.46 based on SEM imaging and the effective refractive

index of the AR layer is 1.62 as calculated in the Supporting Information. Figure 1d is a cross-sectional image of the H-TiO₂ photoanode (V) sandwiched by a thin upper-layer of AR and a photon trapping under-layer. The surface area of the 6 μm -thick H-TiO₂ obtained by BET was $119 \text{ m}^2/\text{g}$, which is comparable to that of a compact mesoporous TiO₂-nanocrystalline film ($117 \text{ m}^2/\text{g}$), as shown in Figure S1. Figure 1e illustrates the engineering mechanism of absorbed light to maximize the intensity of the light. When compared to a conventional P25 electrode, wherein a large portion of light is not only reflected but also limited to having a single pass through the device, H-TiO₂ offers a much improved mechanism for light harvesting. Here, the double sided-large dimension structures function as a mirror at both ends: maximally absorbed photons by the help of the AR layer on the top are scattered further by the st-TIO frame on the bottom, permitting effective light harvesting via multi-passes of entrapped photons through the entire device.

The enhanced AR property and photon trapping effects due to coupling of properly designed double-sided light harvesting layers can be quantified by the (specular angle) reflection spectra and diffuse reflection spectra, respectively. The reflection spectra were obtained by shining light from a 5 μm -thick pure P25 and 5 μm top-surface patterned mesoporous P25 (p-P25) composed of 4.5 μm -thick P25 and an upper layer of patterns with a diameter of 500 nm and 500 nm thickness. P25 was chosen as a control sample since it has been shown to have excellent features as a working electrode in photovoltaic devices. In this configuration, the light passes through air/(TiO₂ patterns)/nanocrystalline TiO₂ interfaces from air to the TiO₂ nanocrystalline layer, and it is assumed that the majority of reflectance originates from the interface of air ($n = 1$) and TiO₂ ($n = 2.49$) due to the notable difference of the refractive index. Figure 2a shows that p-P25 gives $\sim 20\%$ lower reflectance as compared to the pure TiO₂ structure over the entire spectral range. This implies that much more light is able to participate in the photoconversion process in p-P25. This is attributable to the gradual change in the refractive index between the air and TiO₂ nanocrystalline layers in p-P25 ($1/1.63/2.49$, $n_{\text{AR}} = 1.63$ as calculated in the Supporting Information), since the porosity of the patterned TiO₂ layer can modulate the effective refractive index of the diffraction grating layer^{29,36,37}. In order to investigate the intensity of scattered light under beam irradiation on the device, the diffuse reflectance spectra of 6 μm -thick TiO₂ films made in the form of P25, P25/st-TIO, and H-TiO₂ were measured as shown in Figure 2b. A schematic illustration of the samples is shown in Figure 2e. Interestingly, compared to the pristine P25 film, both films with a 1 μm -thick st-TIO layer (P25/st-TIO and H-TiO₂) showed significantly higher diffuse reflection spectra in the visible spectrum of light (380–800 nm), indicating that the absorbed light was efficiently scattered within the nanostructure film. The effect of introduction of the light harvesting layers into the P25 film was further exploited by the transmittance spectra by shining light from top to bottom (Figure 2c). As expected from the low absorption coefficient of P25 in the visible region, a significant fraction of the light passes through 6 μm -thick P25 at the longer wavelength range, which clearly shows light harvesting did not occur (black curve). However, when a 1 μm -thick st-TIO layer was introduced into 5 μm -thick P25 on the bottom, creating 6 μm -thick hybrid samples (blue and red curves), the transmission of light is greatly reduced, indicating that most of the light is re-absorbed by the structure in the devices. This is ascribed to the introduction of a few hundred nanometer scale structure as the photon trapping layer. Finally, UV-visible spectroscopic measurements were undertaken in order to explore the ultimate optical response of nanostructured TiO₂ with and without light harvesting layers. Figure 2d shows the UV-visible absorption spectra of an all 6 μm -thick plain P25, p-P25, P25/st-TIO, and H-TiO₂ illuminated at normal incidence from the air side. The absorption peaks spanning from 350 to 400 nm were assigned to the absorption of the conventional TiO₂ nanoparticles.

Figure 1 | Schematic illustrations and electron microscopy images of TiO₂ structures. (a) Schematic illustration of the fabrication process for the hybrid TiO₂ structure (H-TiO₂). (b) Top view SEM image of st-TiO for photon trapping. The inset shows transmittance electron microscopy (TEM) image of the mesoporous surface textured-TiO (st-TiO). (c) Top view SEM image of the patterned mesoporous TiO₂ layer for anti-reflection. The inset shows the close-up image of the P25 particles. (d) Cross-sectional SEM image of H-TiO₂. (e) Schematic illustration of pathway of incoming light engineered by H-TiO₂ compared with the case of P25.

Figure 2 | Optical properties of P25, p-P25, P25/st-TiO₂, and H-TiO₂ (p-P25/st-TiO₂). (a) Reflectance spectra recorded, within the UV-visible light wavelength range, on the pure P25 and top-surface patterned P25 (p-P25). (b) Diffuse reflectance spectra of P25, P25/st-TiO₂, and H-TiO₂, (c) Transmittance spectra, (d) UV-visible absorption spectra, and (e) Schematic illustration of P25, p-P25, P25/st-TiO₂, and H-TiO₂.

An intense and wider spectrum with a broad absorption band ranging from 400 to 800 nm compared to pure P25 is achieved when either a top or bottom light harvesting layer is adopted. Notably, st-TiO₂ contributed more to increment of the absorption band than anti-reflective p-P25 layer, implying the greater influence of photon trapping effects than AR effects on light harvesting. Conclusively, when the double sided layer is applied, greatly enhanced UV spectra in the visible range are observed, even though the density of TiO₂ particles is lower than that of P25. The enhanced UV absorption of H-TiO₂ is attributed to the synergetic combination of AR and photon trapping effects obtained by the larger dimension nanostructures, which allow for longer effective light path lengths and thus create a clear and strong absorption band in the longer wavelength region ranging from 400 nm to 800 nm.

In order to further enhance the working capability of H-TiO₂ in the visible wavelength region, CdSe quantum dots (QDs) were

deposited onto the H-TiO₂ structure. The deposition of CdSe on TiO₂ was carried out via 20 cycles of successive ionic layer adsorption and reaction (SILAR) method. Figure 3a presents TEM images of the CdSe QDs deposited onto H-TiO₂. Slightly randomly deposited spherical QDs with an average particle size of 4.5 nm with lattice spacings of 0.35 and 0.20 nm, which are consistent with the known values for the (002) and (103) planes of CdSe are observed. The enhanced optical properties due to the introduction of QDs on the hybrid samples are compared with the properties observed for CdSe/P25 and pure H-TiO₂ in Figure 3b. An apparent increase and broadening in the optical density near the band gap of CdSe covering a range of 400–800 nm confirms the creation of CdSe nanoparticles and the ability of QD-sensitization. It should be noted that the overlap of the excitation wavelength of QDs and the regions of the wavelength with great absorption of H-TiO₂ is crucial to maximize the efficient sensitization of QDs. CdSe/P25 also presents enhanced

Figure 3 | Characterization of quantum dots. (a) TEM images of CdSe/H-TiO₂. The image in the right panel is a close-up image of the rectangular area of the left image. (b) UV absorption spectra of P25 and H-TiO₂ before and after the deposition of CdSe. (c) Photoluminescence spectra of CdSe and CdSe/H-TiO₂.

absorption after the deposition of CdSe, yet with a much lower absorption coefficient than that of CdSe/H-TiO₂, likely due to low absorption of pristine P25 in the visible region where CdSe is excited. Figure 3c presents photoluminescence (PL) spectra of the CdSe QDs and CdSe/H-TiO₂. Compared to the PL spectrum of CdSe, which shows a typical near-band-edge emission at 630 nm, that of CdSe/H-TiO₂ is remarkably quenched. This clearly indicates a facile electron transfer mechanism at the interface between CdSe and TiO₂ nanoparticles, and thus low likelihood of electron-hole recombination within CdSe.

Discussion

In order to investigate the hydrogen generating efficiency of the samples, photo-electrochemical measurements were carried out on a P25 control sample, p-P25, P25/st-TiO₂, H-TiO₂, and CdSe/H-TiO₂ using Ag/AgCl, a Pt mesh, and as-prepared samples as a reference electrode, counter electrode, and working electrode, respectively, in a 0.24 M Na₂S and 0.35 M Na₂SO₃ electrolyte solution. This solution is commonly used as an electrolyte for quantum dot-based PEC systems, serving as the sacrificial agent. As shown in Figure 4a, which illustrates the electron transfer mechanism of CdSe/H-TiO₂, after the excited electrons are transferred from quantum dots to the conduction band of TiO₂, the electrons are withdrawn to the counter electrode where the hydrogen generation reaction occurs (An image

capturing a large amount of H₂ gas generation is presented in Figure S6 along with a movie clip)^{38,39}. Meanwhile, the holes diffused to the surface of TiO₂ or CdSe oxidize S²⁻. Figure 4b presents linear sweep voltammograms of the pristine P25, p-P25, P25/st-TiO₂, and H-TiO₂, which is described by the direct photoexcitation of electrons from the valence band of TiO₂, under AM 1.5 G simulated sunlight illumination (100 mW/cm²). All samples with a light harvesting layer show higher photocurrent density than the P25 control sample. Surprisingly, the photocurrent density of an optimized H-TiO₂ reaches ~1.65 mA/cm² at an applied bias of 0.5 V⁴⁰, which is approximately 2.5 times greater than that of the P25 control sample. This is the highest reported value of current density achieved using a single TiO₂ electrode material, implying that our approach to maximize light harvesting via engineering structural control is successful and also suggesting massive potential of pristine H-TiO₂ in photo-conversion systems. Figure 4c compares linear sweep voltammograms of CdSe/H-TiO₂ and CdSe/P25. When the system is further supported by QD sensitization, the photocurrent density of CdSe/H-TiO₂ reaches ~16.2 mA/cm², 35% higher than that obtained with CdSe/P25, under illumination of 100 mW/cm². More surprisingly, as shown in Figure 4c, the photocurrent density of CdSe/H-TiO₂ created by only visible light illumination (blue curve, AM 1.5, a cutoff filter (>420 nm)) reaches ~14.2 mA/cm², ~88% of that of UV-visible irradiation (red curve) and comparable to the values under UV-visible irradiation in other detailed studies^{41–43}, clearly indicating

Figure 4 | Schematic diagram and photocurrent response of P25 control, CdSe/P25, H-TiO₂, and CdSe/H-TiO₂ under AM 1.5 illumination (100 mW/cm²). (a) Schematic diagram showing the hydrogen generating mechanism of CdSe/H-TiO₂. (b) Linear sweep voltammetry measurements of P25, p-P25, P25/st-TiO₂, and H-TiO₂ under UV-visible and dark conditions. (c) Linear sweep voltammetry measurements of CdSe/P25 and CdSe/H-TiO₂ under UV-visible, visible light illumination, and dark conditions. (d) J-t curves of CdSe/H-TiO₂ at a bias of 0.5 V under UV-visible and visible light illumination. (e) EIS response of P25 and H-TiO₂. The measurement was conducted at a frequency range from 100 kHz to 0.1 Hz under open circuit voltage conditions. (f) Measurement of the incident photon-to-electron conversion efficiency. All experiments were performed with a 0.24 M Na₂S and 0.35 M Na₂SO₃ electrolyte (pH = 12) and 100 mW/cm² xenon lamp coupled with a UV/IR filter as the light source.

solely visible working capability of CdSe/H-TiO₂. This surprising result is ascribed to the large enhancement of absorption of H-TiO₂ in the visible range, which matches well with the excitation wavelength of CdSe, via efficient light harvesting. Figure 4d shows the photoresponse over time (J-t curve) of CdSe/H-TiO₂ measured at 0.5 V with chopped illumination of AM 1.5 at a rate of 10 s exposure followed by 10 s non-illumination for UV-visible and visible light.

The sharp spike in the photocurrent for both cases during the on/off illumination cycles indicates fast transport of photogenerated electrons from QDs to the current collector via the interior TiO₂ networks in accordance with the result of the PL study (Figure 3c). The greater values in the photocurrent density of CdSe/H-TiO₂ than in that of CdSe/P25 can be confirmed by a lower internal resistance measured by the electrochemical impedance spectroscopy (EIS). The

measurement was conducted at a frequency range from 100 kHz to 0.1 Hz under open circuit voltage conditions, which gives Nyquist plots composed of two semicircles. A small (first) semicircle at a high frequency represents the charge transfer resistance at the interfaces of the electrolyte/counter electrode whereas a large (second) semicircle at an intermediate frequency conveys the transport resistance of electrons at the $\text{TiO}_2/\text{QD}/\text{electrolyte}$ interface⁴⁴. The smaller diameter of the second semicircle in $\text{CdSe}/\text{H-TiO}_2$ indicates a lower interfacial charge transfer resistance between the TiO_2 electrode and electrolytes and thus a more effective separation of photo-generated electron-hole pairs (a faster interfacial charge transfer to the electron donor/acceptor). The incident photon-to-electron conversion efficiency (IPCE) of $\text{CdSe}/\text{H-TiO}_2$ exhibits overall higher values than that of $\text{CdSe}/\text{P25}$ with a maximum IPCE value of 70.1% (43.7% for $\text{CdSe}/\text{P25}$) at 575 nm, as shown in Figure 4e. Importantly, much stronger and broader responses of $\text{CdSe}/\text{H-TiO}_2$ were achieved than those of $\text{CdSe}/\text{P25}$ from 400 to 630 nm, clearly indicating that $\text{CdSe}/\text{H-TiO}_2$ offers superior quantum efficiency in water splitting in the visible region via light harvesting.

In conclusion, we report a finely designed H-TiO_2 structure that maximizes light harvesting triggered by the configuration of a double sided layer with a few hundred nanometer scale while retaining the sufficient surface area. The observed optical properties confirm excellent light harvesting effects in the H-TiO_2 structure, which can facilitate the water splitting reaction of a TiO_2 semiconductor in the UV-visible range. The photocurrent of $\text{CdSe}/\text{H-TiO}_2$ ($\sim 16.2 \text{ mA}/\text{cm}^2$) was 35% greater than that of the representative $\text{CdSe}/\text{P25}$ photoelectrode, which is clearly recognized for outstanding performance. Importantly, $\text{CdSe}/\text{H-TiO}_2$ under filtered exposure ($\lambda > 420 \text{ nm}$) conditions recorded current density of $\sim 14.2 \text{ mA}/\text{cm}^2$, the highest value in the visible range, which is attributable to the overlap of the excitation wavelength of CdSe with the wavelength where light harvesting is observed. The excellent performance of $\text{CdSe}/\text{H-TiO}_2$ suggests that this strategy can provide a novel route to design photoactive materials via systematic engineering of light for maximum photoconversion efficiency.

Methods

Fabrication of p-P25 and H-TiO_2 . In order to fabricate p-P25, a $5 \text{ mm} \times 5 \text{ mm}$ size patterned PDMS with 500 nm thickness was prepared from the patterned SU 8 with air holes with diameter of 500 nm and 500 nm in thickness. Scotch tape was attached to the FTO substrate to define the sample size as $5 \text{ mm} \times 5 \text{ mm}$. P25 paste was placed and flattened via the doctor blade method so that paste had uniform height corresponding with that of the scotch tape. The P25 paste was imprinted by the patterned PDMS with proper pressure so that no air bubbles formed. After annealing the samples at 55°C for 45 min to remove the residual solvent, the patterned PDMS and the scotch tape were peeled off. The PDMS can be reused after washing with ethanol. The sample was heated at 500°C for 30 min for calcination, followed by TiCl_4 treatment with a concentration of 40 mM at 70°C for 30 min and final heat treatment at 500°C for 30 min. When the st-TiO/FTO substrate was replaced by a FTO substrate, a H-TiO_2 structure was created. The fabrication process for the st-TiO substrate is described in the Supporting Information.

Fabrication of $\text{CdSe}/\text{H-TiO}_2$. The successive ionic layer adsorption and reaction (SILAR) process was used to deposit CdSe onto a $\text{H-TiO}_2/\text{FTO}$ substrate. A sodium selenosulfate (Na_2SeSO_3) solution was prepared by dissolving 0.01 mol Se powder in a 0.1 M Na_2SO_3 solution at 90°C for 9 hours. To assemble CdSe onto H-TiO_2 , the $\text{H-TiO}_2/\text{FTO}$ substrate was dipped into a 50 mM $\text{Cd}(\text{Ac})_2 \cdot 2\text{H}_2\text{O}$ ethanol solution for 1 min, washed with ethanol for 30 s, and then dipped into a Na_2SeSO_3 solution for 1 min and rinsed with D.I. water for 30 s. Finally, the $\text{CdSe}/\text{H-TiO}_2$ substrate was dried by blowing nitrogen gas. The two-step dipping procedure is defined as one SILAR cycle. This procedure was repeated twenty times to obtain the desired performance of CdSe and carried out in a glove box filled with nitrogen gas.

Characterization. The morphology of the structures was examined by SEM (NOVA NANOSEM 230 FESEM, 15 kV), TEM (JEM-2100, 200 kV). The powder diffraction data were obtained using a Rigaku Co. High Power X-Ray Diffractometer D/MAZX 2500V/PC from 20° to 80° . Optical properties measurement (Absorbance, Reflection, Transmission) of structures were investigated by UV-visible spectroscopy (VARIAN, Cary 100). The diffuse reflection was measured using a Cary 5000 UV/Vis/NIR. Photoluminescence (PL) spectrum of the $\text{CdSe}/\text{H-TiO}_2$ was measured by spectrofluorometer (JASCO, FP-8500). Photoconversion efficiency was confirmed using the IPCE measurement system (QEX10, Pv measurements).

PEC measurements. The PEC performance of the TiO_2 electrodes was exploited in a three-cell electrode system under front-side illumination of AM 1.5 G (Newport solar simulator) using a potentiostat (Princeton Applied Research VersaSTAT3, AMETEK). An Ag/AgCl electrode and a Pt mesh were used as reference and counter electrodes, respectively. A solution of 0.24 M Na_2S and 0.35 M Na_2SO_3 was used as an electrolyte and sacrificial solution. An exposed area of the working electrode was 0.25 cm^2 . The photoresponse was evaluated under chopped illumination from a Newport solar simulator equipped with an AM 1.5 G filter, calibrated with a standard Si solar cell simulating AM 1.5 G. The power of the solar simulator was measured to be $100 \text{ mW}/\text{cm}^2$. Photocurrent stability tests were carried out by measuring the photocurrent produced under chopped light irradiation (light/dark cycles of 10 s) at a bias of 0.5 V versus Ag/AgCl . The PEC performances in the visible range were acquired by the solar simulator coupled with a UV cutoff filter ($\lambda > 420 \text{ nm}$) in Figure S5. EIS measurement were performed under the same condition of gas evolutions under illumination of AM 1.5 G.

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Acknowledgments

This work is supported by NRF with the contract no. NRF-2010-0019408 and NRF-R1A1A2043076 and by the development program of local science park funded by the ulsan metropolitan city and the MSIP.

Author contributions

The experiments were designed by K.K. and M.K. with suggestions from J.J.; K.K. prepared H-TiO₂ samples. K.K. and M.K. performed PEC and optical property measurements. S.K. performed BET analysis and EIS measurements. J.J. analyzed the data and wrote the manuscript. All authors discussed the results and commented on the manuscript

Additional information

Supplementary information accompanies this paper at <http://www.nature.com/scientificreports>

Competing financial interests: The authors declare no competing financial interests.

How to cite this article: Kim, K., Kim, M.-J., Kim, S.-I. & Jang, J.-H. Towards Visible Light Hydrogen Generation: Quantum Dot-Sensitization via Efficient Light Harvesting of Hybrid-TiO₂. *Sci. Rep.* **3**, 3330; DOI:10.1038/srep03330 (2013).

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