

## Dirac lines and loop at the Fermi level in the time-reversal symmetry breaking superconductor LaNiGa<sub>2</sub>

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Unconventional superconductors have Cooper pairs with lower symmetries than in conventional superconductors. In most unconventional superconductors, the additional symmetry breaking occurs in relation to typical ingredients such as strongly correlated Fermi liquid phases, magnetic fluctuations, or strong spin-orbit coupling in noncentrosymmetric structures. In this article, we show that the time-reversal symmetry breaking in the superconductor LaNiGa<sub>2</sub> is enabled by its previously unknown topological electronic band structure, with Dirac lines and a Dirac loop at the Fermi level. Two symmetry related Dirac points even remain degenerate under spin-orbit coupling. These unique topological features enable an unconventional superconducting gap in which time-reversal symmetry can be broken in the absence of other typical ingredients. Our findings provide a route to identify a new type of unconventional superconductors based on nonsymmorphic symmetries and will enable future discoveries of topological crystalline superconductors.

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The combination of superconductivity with topology is expected to exhibit new types of quasiparticles such as non-Abelian Majorana zero modes<sup>1,2</sup>, or fractional charge and spin currents<sup>3</sup>. Experimental realization of topological superconductivity can also provide new platforms for fault-tolerant quantum computation technologies<sup>4</sup>. Topological superconductivity can be artificially engineered in hybrid structures<sup>5–9</sup> or it can exist intrinsically in certain unconventional superconductors<sup>10–15</sup>. In most intrinsic topological superconductors, the unconventional nature of superconductivity originates from the proximity to magnetic instabilities or strong electronic correlations<sup>15</sup>. Regardless of how the topological superconducting state is produced, it is well established that there is a vital connection between a material's symmetries, broken and preserved, and the resulting topology<sup>16–20</sup>.

We report that the time-reversal symmetry-breaking superconductor LaNiGa<sub>2</sub> derives its unconventional superconducting pairing from the previously unknown existence of Dirac lines and Dirac loop in the normal state. These features are pinned at the Fermi energy where they impact low energy properties including superconductivity. The rich topology of the electronic structure originates from the nonsymmorphic symmetry that guarantees band degeneracies, which in turn, enable interband and/or complex superconducting order parameters that can break time-reversal symmetry. Our results illustrate a novel method towards realizing intrinsic (single-material) topological superconductivity wherein the underlying space group symmetry intertwines the topology with the unconventional superconductivity.

The centrosymmetric superconductor LaNiGa<sub>2</sub> was previously known to break time-reversal symmetry when it becomes superconducting below  $T_{sc} = 2\text{ K}$ <sup>21</sup>. Subsequent penetration depth, specific heat, and upper critical field studies showed nodeless multigap behavior<sup>22</sup>, in contradiction with possible single-band spin-triplet pairing<sup>22,23</sup>. All previous experimental investigations were limited to polycrystalline samples and theoretical considerations were based on the previously reported

symmorphic *Cmmm* (No. 65) space group<sup>24</sup>. We reveal that single-crystal X-ray diffraction (SCXRD) analysis improves upon previous powder X-ray diffraction (PXRD) work and properly assigns LaNiGa<sub>2</sub> to a nonsymmorphic *Cmcm* (No. 63) unit cell. Difficulty discerning the difference between *Cmmm* and *Cmcm* from PXRD data has historical precedent<sup>25,26</sup>.

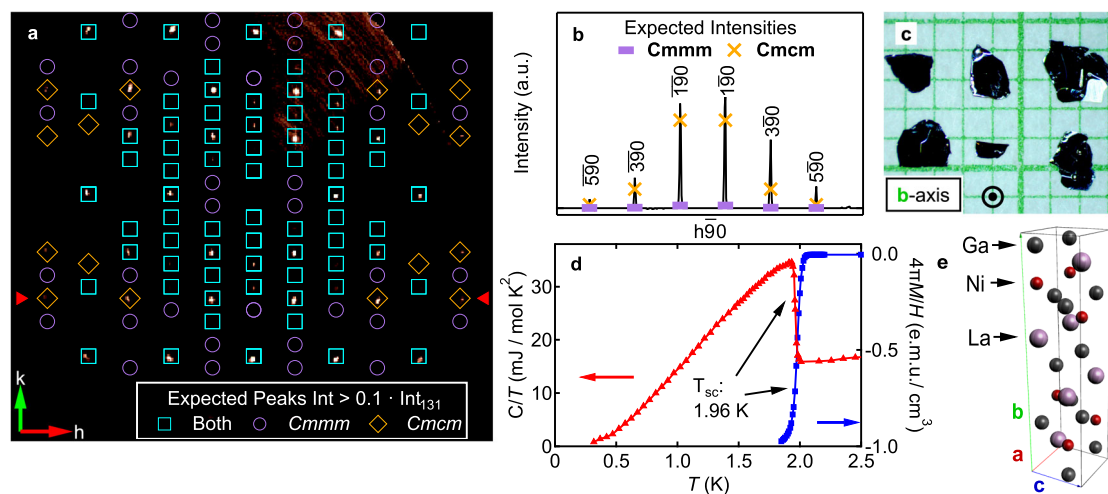
The nonsymmorphic symmetries of this new unit cell transform the  $k_z = \pi/c$  plane of the Brillouin zone (BZ) into a nodal surface, which hosts fourfold degenerate bands<sup>27,28</sup>. Here, the band degeneracies form two distinct Dirac crossings between two sets of Fermi surfaces (FSs) precisely at the Fermi level ( $E_F$ ), independent of chemical potential position. There are fluted lines closed by BZ periodicity and a closed loop. Of special note is that the Dirac loop contains two points which are protected against splitting from spin-orbit coupling (SOC). These “touchings” are shown from our band structure calculations, along with angle-resolved photoemission spectroscopy (ARPES) data.

We note that, among non-magnetic materials and outside of intercalated Bi<sub>2</sub>Se<sub>3</sub>, no other time-reversal symmetry breaking superconductor has been shown to exhibit a topological band structure (see Supplementary Note 12 and Table S4). Thus making LaNiGa<sub>2</sub> unique amongst this small set of bulk superconductors. Lastly, we discuss the impact of the topology of LaNiGa<sub>2</sub> as a natural platform for interband pairing and/or complex superconducting order parameter that can break time-reversal symmetry.

## Results and discussion

**Structural characterization.** Single-crystalline samples of LaNiGa<sub>2</sub> were successfully grown with a Ga deficient self-flux technique. Details about the growth are contained within the “Methods” section below. Highly reflective, plate-like crystals were produced as shown in Fig. 1c.

SCXRD data were collected on several samples and each data set resolved to a LaNiGa<sub>2</sub> unit cell with a *Cmcm* space group (see



**Fig. 1 Structural and superconducting properties of LaNiGa<sub>2</sub> single crystals.** **a** Compiled precession image of the  $hk0$  plane from a single-crystal X-ray diffraction data set of LaNiGa<sub>2</sub>. Overlaid are the predicted diffraction spots with a normalized intensity above 0.1. The diffraction spots are normalized to the 131 diffraction peak. The diffraction spots which are expected for both the original *Cmmm*<sup>24</sup> and our newly proposed *Cmcm* structures are denoted by the teal squares. The spots which are only expected for the *Cmmm* and *Cmcm* structures are shown by the purple circles and orange diamonds, respectively. **b** A normalized linear cut of the  $hk0$  precession image along the  $h90$ , shown by the red triangles in **a**. The expected intensities for the *Cmmm* and the *Cmcm* structures are shown as the purple dashes and the orange crosses, respectively. All intensity data, raw and theoretical, from **a** and **b** are normalized to the 131 peak. **c** Picture of representative single-crystal samples of LaNiGa<sub>2</sub>. The plate-like samples have the  $b$ -axis normal to the surface of the crystals. **d** A complete superconducting transition is observed with zero-field specific heat capacity ( $C/T$ ) and zero-field cooled magnetic susceptibility ( $M/H$ ) data, shown by the red and blue curves, respectively. The susceptibility was collected with an external magnetic field of 1 mT. **e** The *Cmcm* unit cell for LaNiGa<sub>2</sub>.

Supplementary Notes 1 and 2, Fig. S2, and Tables S1 and S2). Given the inherent similarities between the previously reported  $Cmmm$  structure<sup>24</sup> in real space, nearly all diffraction spots within the reciprocal space are predicted by both structures (see Fig. S1 for PXRD LeBail fittings using each structure and Fig. S3 for a real space comparison between the two structures). This is especially true for the most intense, low-angle diffraction spots. There are, however, a few observable differences amongst the weakly-diffracting high-angle spots that are sufficient to differentiate the two structures, as shown from the compiled  $hk0$  precession image in Fig. 1a.

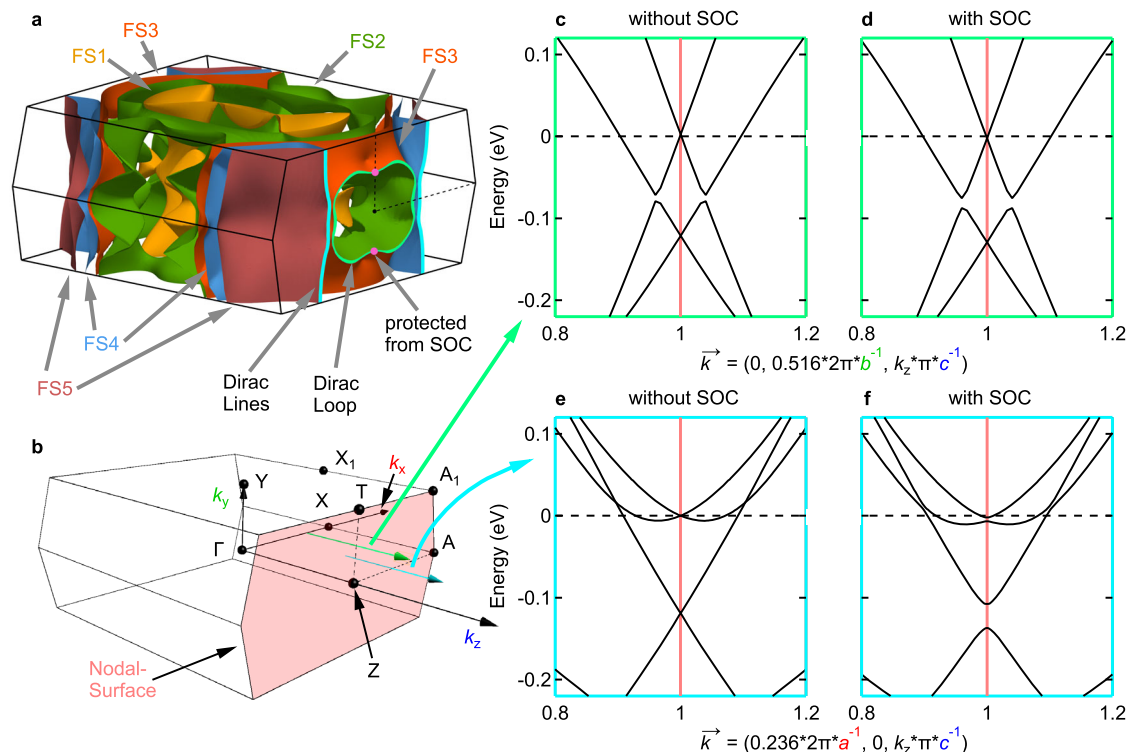
These discrepancies are highlighted by the differing expected intensities along the normalized  $h\bar{9}0$  line, red arrows in Fig. 1a, between the two structures (Fig. 1b). It should be noted that all peaks along this line from both models are displayed, regardless of intensity. The insufficient intensities for the  $Cmmm$  model at these high-angle spots reveal that the previously reported structure inadequately matches the observed diffraction data for this material. The data indicates that the more accurate structure for this material is that of the  $Cmcm$  space group.

While this new space group remains centrosymmetric, it adds nonsymmorphic symmetries: a  $c$  glide plane perpendicular to the  $b$ -axis and a  $2_1$  screw axis,  $\mathcal{S}_2^z$ , along the  $c$ -axis. First we will discuss the influence of the new symmetry operations on the electronic structure and later we will examine the ramifications on the superconducting pairing state.

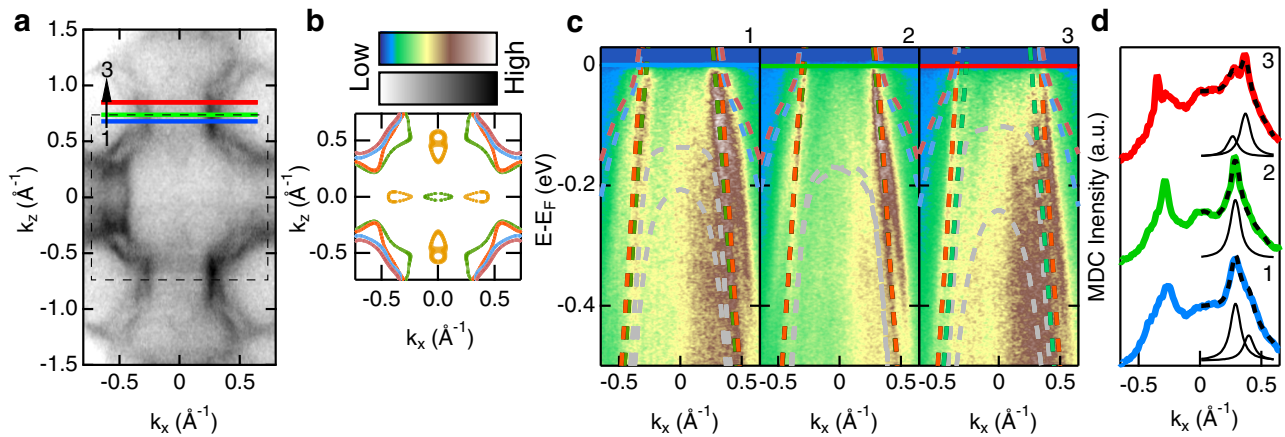
**Electronic structure and ARPES data.** Despite the new structure, Fig. 2a shows that there remain five FSs<sup>23,29</sup>. Highlighted with the previous space group, there are several regions within the BZ where pairs of FSs are parallel and quasi-degenerate<sup>22,23</sup>. A

crucial difference now is that the nonsymmorphic symmetry operations force the previously quasi-degenerate FS pairs to “touch” on the  $k_z = \pi/c$  plane (red plane) in the absence of SOC, turning this plane into a nodal surface<sup>27</sup> (see Supplementary Note 11 and the band structures along  $T-Z-A$  in Fig. S11 to see bands become degenerate in the  $Cmcm$  structure). The combination of  $\mathcal{S}_2^z$ , parity operation, and time-reversal symmetry force bands on the nodal surface to be fourfold degenerate (for the derivation and for the differences in the  $Cmmm$  and  $Cmcm$  band structures see Supplementary Note 8). This symmetry-enforced degeneracy results in two disjoint sets of Dirac crossings directly at the Fermi level. Both are between bulk bands crossing the nodal surface: fluted lines across the BZ face between FS4/5 and a closed loop between FS2/3 (highlighted lines top panel Fig. 2a). That these crossings occur at the Fermi level make  $\text{LaNiGa}_2$  uncommon compared to other superconductors with topologically non-trivial band structures<sup>10,14,30–37</sup> (see Table S4).

The Dirac crossings can be observed in the linear band dispersion plots without SOC along  $\vec{k} = (0, 0.516\frac{2\pi}{b}, k_z)$  (green arrow) for the Dirac loop (Fig. 2c) and  $\vec{k} = (0.236\frac{2\pi}{a}, 0, k_z)$  (blue arrow) for the Dirac lines (Fig. 2e). We note that small shifts of the Fermi energy will shift the  $k$ -space location of the Dirac lines and loop. However, since this is a matter of two surfaces intersecting (FS and nodal surface) and nonsymmorphic symmetry enforcement, the nodal lines will persist at the Fermi level as long as the FSs cross the nodal surface. When accounting for SOC, most band crossings become gapped (by a few to 40 meV), as pictured in Fig. 2f. Remarkably, the Dirac points between FS2/3 survive along the  $Z-T$  symmetry line under SOC, as seen in Fig. 2d, creating two true-Dirac points at the Fermi



**Fig. 2 Fermi surfaces and Dirac band crossings.** **a** The five Fermi surfaces (FSs), labeled as FS1–5, within the Brillouin zone (BZ). The Dirac lines (blue lines) and Dirac loop (green line) are highlighted on the nodal surface. The crossings along  $Z-T$  (magenta dots) are protected from spin-orbit coupling (SOC). **b** The BZ showing several high symmetry points and highlights the nodal surface (red plane). The green arrow,  $k_y = 0.516 \cdot 2\pi/b$ , shows where FSs 2 and 3 become degenerate on the nodal surface. The Dirac crossing is shown to remain both without **c** and with SOC **d**, where the SOC contribution to anticrossing is seen to be very small. The blue arrow shows the dispersion along  $k_x = 0.236 \cdot 2\pi/a$  without SOC **e** shows the Dirac lines between FSs 4 and 5. Once SOC is added **f**, the crossing becomes gapped at the nodal surface.



**Fig. 3 ARPES measurements revealing the band touching on the nodal surface.** Angle-resolved photoemission spectroscopy (ARPES) characterization of  $\text{LaNiGa}_2$ . **a** Constant energy map with an integration window of  $\pm 10$  meV around the Fermi energy,  $E_F$ . The black dotted line indicates the boundary of the Brillouin zone in the  $k_y = 0$  plane. The blue, green, and red solid horizontal lines indicate cuts (1), (2), and (3) in panel **c**, respectively. **b** Calculated Fermi surfaces (FSs) on the  $k_y = 0$  plane with the colors corresponding to FSs in Fig. 2a. See Fig. S10c for an overlay of the ARPES and the calculated FSs on the  $k_y = 0$  plane. **c** Energy vs momentum spectra along cuts indicated in **a**. The dotted lines are the overlay of electronic structure calculations and the colors show which FSs are associated with each band. The gray bands are low energy bands which do not cross  $E_F$ . **d** Momentum distribution curves (MDCs) at  $E_F$  from the cuts in panel **c**. Spectra are fit to one (two) Lorentzian peaks (dotted black curve) for cuts 2 (1,3), with a Gaussian background centered on  $k_x = 0$ . The black curves below the experimental data are the individual Lorentzian peaks marking where the bands cross  $E_F$ . The peak width for the Lorentzian fit on 2 is a free fitting parameter and fixed to that width for the peaks of cuts 1 and 3.

level. This protected feature results from the presence of the mirror reflection,  $\mathcal{M}_x$ , along the  $Z - T$  line, therefore, remaining fourfold degenerate even when accounting for SOC (see Supplementary Note 8 for derivation), illustrating that this degeneracy lies precisely at  $E_F$ , and is robust rather than accidental.

Single crystals of  $\text{LaNiGa}_2$  do not naturally cleave perpendicular to the crystallographic  $c$ -axis, making a direct observation of the Dirac dispersion by ARPES measurements challenging. However, with a photon energy of 144 eV we can probe the  $k_y = 0$  plane and confirm the presence of the band touchings (see Supplementary Note 10 and Fig. S10). Figure 3a shows the constant energy map centered at  $E_F$  and reveals the most prominent features of the spectra: the ruffled cylindrical bands centered on the BZ corners. Given that, near the corner of the BZ, the calculated FSs are very close to each other (see Fig. 3b), it is difficult to discern which bands are observed in the ARPES measurements from just this plane. Overlaid on Fig. 3c are the respective DFT band calculations (dashed lines) which reveal that the most prominent bands in the ARPES data originate from the bands associated with FS2 and FS3. The three parallel horizontal cuts on and near the nodal surface show the band dispersion plots at and below  $E_F$  (Fig. 3c). The green line, spectrum 2, represents the cut exactly on the nodal surface, while the blue, spectrum 1, and red, spectrum 3, lines are parallel cuts in the first and second BZs, respectively.

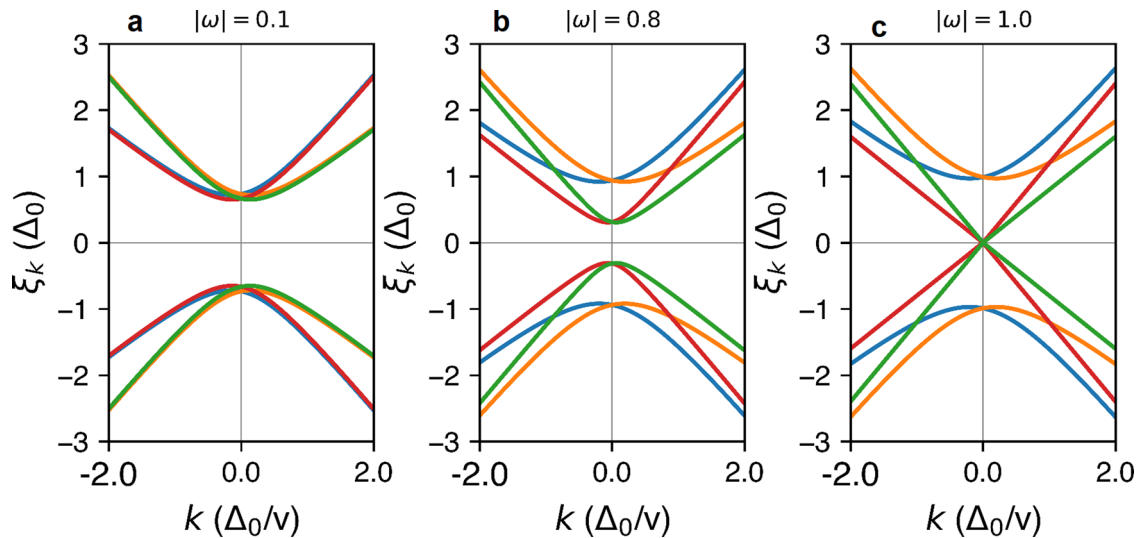
Each of these cuts was integrated within 50 meV of  $E_F$  to produce momentum distribution curves (MDCs) shown in Fig. 3d. On the nodal surface, spectrum 2 shows a single clear peak representing the degeneracy of FS2/3. Off the nodal surface, the MDCs for spectra 1 and 3 show that FS2 and FS3 separate and are no longer degenerate. Thus providing direct evidence for the band degeneracy between FS2/3 on the nodal surface. As mentioned above, we expect SOC to split the FS2/3 crossing on the  $k_y = 0$  plane of the nodal surface. We note, however, that the SOC gap cannot be resolved because the peaks have a smaller calculated momentum separation than the fitted experimental widths. This result is further evidence for the minimal impact of SOC on the electronic structure of  $\text{LaNiGa}_2$  in the normal state.

In the normal state and in the absence of SOC, the  $Cmcm$  space group makes  $\text{LaNiGa}_2$  a topological nodal line metal. The line (or loop) is topological<sup>38</sup>. Nodal lines (lines or loops of degeneracies) in band structures have been found to be rather common<sup>39,40</sup>. However,  $\text{LaNiGa}_2$  is so far unique in having the nodal lines lie precisely at the Fermi level. However, this confluence of bands will occur in any nonsymmorphic metal with FSs crossing the nodal surface where bands are guaranteed to be orbitally degenerate.

**Pairing model and quasiparticles.** Now we examine the consequences of the  $Cmcm$  space group assignment for the superconducting state.  $\text{LaNiGa}_2$  has low symmetry and previous symmetry analysis based on the  $D_{2h}$  point group revealed only 4 possible gap functions that break time-reversal symmetry<sup>21,41</sup>. All of them have nodes inside the BZ, which is incompatible with thermodynamic measurements on polycrystals<sup>22</sup>, as well as our heat capacity measurements on single crystals which indicate nodeless fully gapped superconductivity (Fig. 1d). The presence of nonsymmorphic symmetries modifies the nodal behavior on the  $k_z = \pi/c$  plane with or without SOC (see our classification in Supplementary Note 9), but does not provide a scenario for the absence of nodes inside the BZ. The five FSs in Fig. 2 indicate that the full FS is large and pervasive throughout the zone, thus any superconducting gap nodes in a direction  $\hat{k}$  would produce a gap node on the FS and thus be detectable in thermodynamic measurements. This observation limits the possible superconducting states to  $A_{1g}$  with or without SOC (see Supplementary Note 9), but these states do not break time-reversal symmetry. The superconducting properties of  $\text{LaNiGa}_2$  cannot be understood without involving interband pairing<sup>22,23</sup>. The topological properties of the normal state now provide a natural platform for such unconventional superconductivity.

As mentioned earlier, the Dirac lines and Dirac loop are gapped by SOC, except for the true-Dirac points surviving on the  $Z - T$  line where SOC vanishes. A feature of more interest for the superconducting phase is that, unlike the case for conventional FSs, in Dirac (or Weyl) metals interband transitions persist all the





**Fig. 4 Degeneracies in the Bogoliubov-de Gennes quasiparticles.** Bogoliubov-de Gennes (BdG) quasiparticle bands near the Dirac point for three values of the  $|\vec{\omega}|$  = **a** 0.1, **b** 0.8, and **c** 1.0 resulting in gapped, weakly nonunitary, and gapless fully nonunitary states, respectively. Energy units correspond to  $\Delta_0 = 0.7$  meV,  $v = 500$  and  $\beta = 100$ , see Eq. (1).

way to zero energy. Any single-band model breaks down, and a two-band model is a minimal model<sup>42</sup>. LaNiGa<sub>2</sub> thereby becomes an intrinsically two, degenerate and topological, band superconductor.

If the interband pairs are symmetric in the band index, then the Cooper-pair wave function will have the same symmetry as the intraband pairs do, *s*-wave will be spin-singlet and *p*-wave will be spin-triplet. But if the interband pairs are antisymmetric in the band index, we can have *s*-wave spin-triplet, or *p*-wave spin-singlet pairing while still satisfying the overall fermionic nature required for a superconducting order parameter<sup>43</sup>. If both symmetric and antisymmetric pairing exists simultaneously on the nodal surface (weak SOC) or on the true-Dirac points on the *Z* – *T* line (strong SOC), time-reversal symmetry could be broken in two ways: the band (orbital) channel or the spin channel.

In the band-orbital channel, two-gap functions (for example *s*-wave spin singlet  $A_{1g}^s$  and *s'*-wave spin-triplet  $A_{1g}^t$ ) could form a complex combination similar to  $s + is'$  to break time-reversal symmetry. Another possibility is to combine several triplet components. For example, the  $B_{1g}^t$  triplet will be split by SOC into  $A_{2g}$ ,  $B_{2g}$ , and  $B_{3g}$ . The complex combination could also break time-reversal symmetry. However, a first-order transition or multiple transitions are expected in these cases<sup>22</sup>, but there is no such evidence in our heat-capacity measurements (Fig. 1d and Fig. S7). Time-reversal symmetry breaking in the band-orbital channel is therefore unlikely.

Because of the possibility of *s*-wave spin-triplet pairing on the Dirac lines, loop, or points, time-reversal symmetry could also be broken in the spin channel from an internally antisymmetric nonunitary triplet pairing (INT) state<sup>23</sup>. The power of symmetry analysis is remarkable in the sense that, even with the incorrect space group, the correct point group ( $D_{2h}$ ) already led to the conclusion that the INT pairing is the only reasonable solution<sup>22,23</sup>. However, the necessary degeneracy was not identified because of the wrong space group. An orbital-singlet equal spin pairing has also been proposed for doped Dirac semimetals<sup>44</sup>. The INT model has also been proposed to explain time-reversal symmetry breaking and fully gapped superconductivity in LaNiC<sub>2</sub><sup>23,45,46</sup>. While LaNiGa<sub>2</sub> and LaNiC<sub>2</sub> are compositionally related, our results on LaNiGa<sub>2</sub> highlight new significant differences between the two compounds. LaNiC<sub>2</sub> has a

symmorphic and noncentrosymmetric space group (*Amm*2, No. 38), and thus far no topological band crossing has been reported<sup>45,47–52</sup>. In addition, electrical resistivity measurements under pressure showed the proximity to a different state characterized by a high-energy scale<sup>53</sup>, and magnetic penetration-depth measurements under pressure suggested the proximity of a quantum critical point in LaNiC<sub>2</sub><sup>54</sup>. Further studies remain needed on both materials to confirm the validity of the INT model, and the mechanism of time-reversal symmetry breaking. Our discovery of symmetry-imposed band crossing, even under SOC, in LaNiGa<sub>2</sub>, reinforce the relevance of the INT model for this compound, as well as for other nonsymmorphic superconductors.

Breaking of time-reversal symmetry requires a nonunitary triplet pairing potential  $\hat{\Delta} = i(\Delta_0 \vec{\eta} \cdot \vec{\sigma}) \sigma_y \otimes i\tau_y$ , where the tensor products include the first (spin,  $\sigma$  matrices) channel  $\sigma = \uparrow, \downarrow$  and the second (band,  $\tau$  matrices) channel, with the bands labeled by  $m = \pm$  being degenerate along the Dirac loops, taken to be at  $k_{\perp} = 0$ .

Note that the spontaneous vector field  $\Delta_0 \vec{\eta}$  couples to spin like a magnetic moment. The pairing matrix describes triplet pairing but is antisymmetric in the band channel ( $i\tau_y$ ) to ensure the fermionic antisymmetry of the pair wave function<sup>22</sup>. The experimentally observed time-reversal symmetry breaking is ensured by the non-unitarity, which is characterized by a nonvanishing real vector  $\vec{\omega} = i(\vec{\eta} \times \vec{\eta}^*)$  which satisfies  $|\vec{\omega}| \leq |\vec{\eta}|^2 = 1$ . A noteworthy difference with previous INT proposals<sup>22,23</sup> is that the true two-band situation in LaNiGa<sub>2</sub> is enforced by symmetry, and persists in the Bogoliubov - de Gennes (BdG) quasiparticle bands. Accounting for the linear band coupling  $\beta k_{\perp}$  away from the nodal surface, the dispersion of the eight BdG quasiparticles (two bands, two spins, electrons and holes) becomes

$$\epsilon_k = \pm [\sqrt{\{vk_{\perp} - \mu\}^2 + |\Delta_0|^2(1 \pm |\vec{\omega}|)} \pm \beta k_{\perp}] \quad (1)$$

with degenerate eigenvalues on the nodal surface of  $|\Delta_0| \sqrt{1 \pm |\vec{\omega}|}$ .

The spectrum, shown in Fig. 4, illustrates the 8-band behavior versus the strength of non-unitarity. The nonsymmorphic operations guarantee that pairs of BdG quasiparticle bands persist in

“sticking together” on both sides of the gap at  $k_{\perp} = 0$ , thereby retaining topological character. However, now massive points of degeneracy arise in the BdG band structure, unlike the bands of Ghosh et al.<sup>23</sup> which retain no degeneracies and hence no topological character. The linear band mixing results in the gap edges lying slightly away from the plane  $k_{\perp} = 0$ . Unit values of  $|\vec{\omega}|$  lead to gaplessness, with unusual Weyl dispersion of the BdG quasiparticles. The measured magnetic moment of  $0.012 \mu_B$ , if from spin, corresponds to a small conventional band exchange splitting  $\Delta_{ex} = m/N(E_F) = 1.8$  meV. This splitting is comparable to (larger than) the superconducting gap  $2\Delta_0 \sim 3.5 - 4k_B T_c \sim 0.7$  meV, suggesting it may be central to the exotic pairing mechanism.

In contrast to the INT state, conventional  $p$ -wave spin-triplet superconductivity is expected to lead to high upper critical fields, because equal spin pairing is not subject to Pauli limiting, and because most proposed  $p$ -wave superconductors are heavy fermion systems with high orbital limit<sup>55–60</sup>. LaNiGa<sub>2</sub>, however, is not a heavy-fermion material ( $\gamma_n = 14.1$  mJ mol<sup>-1</sup> K<sup>-2</sup>) and interband pairing is suppressed by the application of a strong magnetic field. Thus the upper critical field in LaNiGa<sub>2</sub> remains low (see Fig. S9), even though time-reversal symmetry breaking superconductivity was observed at zero field in  $\mu$ SR experiments.

Our findings reveal that LaNiGa<sub>2</sub> is a topological nonsymmorphic crystalline superconductor<sup>61–64</sup>. The normal state electronic structure features Dirac lines and Dirac loop at the Fermi energy enforced by nonsymmorphic symmetries, as well as true-Dirac points that retain their degeneracy under SOC. These findings are expected to be common to a large number of materials with similar crystalline symmetries. In general, when topological materials become superconducting, the superconducting state is unconventional. LaNiGa<sub>2</sub> was previously reported as a time-reversal symmetry-breaking superconductor with evidence for a fully gapped superconducting state, but the topological properties were unknown. The topological character now provides a natural platform for the INT state to exist. Further experiments and theoretical proposals are necessary to further elucidate the time-reversal symmetry breaking mechanism.

Because of the possibility of a fully gapped behavior mimicking conventional superconductivity, many other materials could have been overlooked. Our results on LaNiGa<sub>2</sub> motivate the need to characterize the time-reversal symmetry, with zero-field  $\mu$ SR experiments or Kerr effect, of other crystalline topological metals<sup>18–20</sup> that become superconducting. While LaNiGa<sub>2</sub> is the only intrinsic topological material with nodal features at the Fermi level which has been experimentally shown to break time-reversal symmetry in the superconducting state without any overlapping magnetic fluctuations or ordering, other materials could soon be discovered based on our findings.

## Methods

**Sample preparation.** Single-crystalline samples of LaNiGa<sub>2</sub> were grown with a Ga deficient self-flux technique. Ga (99.99999%) atomic composition ranged from 32–36% Ga and the remaining percentage was equally split between La (99.9996%) and Ni (99.9999%). Precursor ingots were first synthesized by arc melting all the elements in an argon environment. The ingots were subsequently loaded into an alumina Canfield crucible set<sup>65</sup> and sealed in an evacuated quartz ampule. The reaction was heated up to 1150 °C and held at temperature for five hours. The reaction was slowly cooled down to 800 °C over 100 h and then quickly centrifuged. Overall, high-quality single crystals were synthesized and characterized (see Supplementary Notes 3–7, Figs. S4–9, and Table S3).

It was noted that different starting Ga percentages did not produce a noticeable difference in crystal quality, as evaluated by the residual resistivity ratio. However, larger single crystals (up to 7 mm) were obtained in the more Ga deficient syntheses. Additionally, it was also discovered that the superconducting properties were highly sensitive to oxidation throughout the reaction. Lastly, in more Ga deficient growths, below 32% Ga, no crystals were obtained when the reactions were centrifuged at 800 °C.

**Crystal Structure Determination.** Each synthesis was checked to produce the desired phase by PXRD performed on a Rigaku Miniflex with a Cu X-ray source. LeBail fitting was performed using both the  $Cmmm$  and  $Cmcm$  space groups in GSAS-II<sup>66</sup>. Selected samples were aligned using a Laue X-ray diffractometer to distinguish the  $a$ - and  $c$ -axes directions. SCXRD data were collected on several samples of LaNiGa<sub>2</sub> at 100 K using a sealed-tube Mo X-ray source on a Bruker Photon 100 CMOS X-ray diffractometer (Bruker AXS). Across several crystals, obvious twin domains were observed within diffraction space; although not all samples exhibited this. Regardless of the presence of multiple domains, initial unit cell parameters for each sample suggested a  $C$ -orthorhombic unit cell that matches well with previous reports:  $a = 4.2808$  Å,  $b = 17.466$  Å, and  $c = 4.25778$  Å (ICSD Nos. 634496 and 634508)<sup>24</sup>. The collected frames were integrated using SAINT within APEX3 version 2017.3. For every crystal that was diffracted, XPREP suggested the centrosymmetric space group  $Cmcm$  (No. 63) and the structure was refined down to a R value of 0.0288 using SHELXL-2018/3<sup>67</sup>. See Tables S1 and S2 for further crystallographic information and Supplementary Data 1 for the CIF file. This  $Cmcm$  structure is of the BaCuSn<sub>2</sub> structure type (ICSD No. 58648). The precession image was compiled within APEX3. Structure factors for the precession image models were calculated from Visualization for Electronic and Structural Analysis (VESTA) Ver.3.4.7<sup>68</sup>.

**Physical property measurements.** Low-frequency AC resistivity measurements were measured using a four-probe technique on a Quantum Design Physical Property Measurement System (PPMS) from 300–1.8 K. The PPMS was also used to obtain heat capacity data for selected samples using a relaxation technique down to 1.8 K. A <sup>3</sup>He insert for the PPMS allowed for measurements of AC resistivity and heat capacity down to 0.4 K. Magnetization measurements were collected in a Quantum Design DC Magnetic Property Measurement System down to 1.85 K.

**Electronic structure methods.** Density functional-based electronic structures were produced by the precise linearized augmented planewave code WIEN2k<sup>69,70</sup> using the generalized gradient functional for exchange and correlation. The atomic coordinates used for these calculations were the same as listed above. The sphere sizes were, in Bohr: La, 2.50; Ni, 2.40; Ga, 2.12. The plane-wave cutoff  $K_{max}$  was determined by  $RK_{max} = 7$ , and the  $k$ -point mesh for self-consistency was  $14 \times 14 \times 14$ . Exchange and correlation contributions to the energy and potential were included by using the generalized gradient approximation functional<sup>71</sup>. Effects of spin-orbit coupling were included by using second variation method as implemented in WIEN2k.

**ARPES measurements.** ARPES measurements were performed at Stanford Synchrotron Radiation Lightsource National Laboratory beamline 5–2 using a Scientia DA30 electron spectrometer. Samples were cleaved in-situ at 20 K and with a pressure better than  $5 \times 10^{-11}$  Torr. The nominal resolution for the ARPES measurements was  $\Delta E = 20$  meV, the beam spot size was approximately  $21 \mu\text{m} \times 36 \mu\text{m}$ .

## Data availability

The data that support the findings of this study are available from the authors upon reasonable request.

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### Author contributions

J.R.B. and V.T. conceived the project. Single-crystal synthesis work was initiated by K.N. and V.T. J.R.B. completed the synthesis and performed all X-ray crystallography, electrical resistivity, magnetic susceptibility, and heat capacity measurements with assistance from J.C.F., K.P.D., and S.M.K. for single-crystal X-ray diffraction, and P.K., D.A., V.T. for other physical property measurements. Y.Q. and W.E.P. completed the electronic structure calculations and its discussion. M.C.S., A.R., D.S.S., and I.M.V. performed the ARPES measurements and the subsequent analysis. J.R.B., M.C.S., I.M.V., W.E.P., and V.T. wrote the manuscript with input from all authors.

### Competing interests

The authors declare no competing interests.

### Additional information

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