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# Super-Nernstian ion sensitive field-effect transistor exploiting charge screening in WSe<sub>2</sub>/MoS<sub>2</sub> heterostructure

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lon-sensitive field-effect transistors (ISFETs) have gained a lot of attention in recent times as compact, low-cost biosensors with fast response time and label-free detection. Dual gate ISFETs have been shown to enhance detection sensitivity beyond the Nernst limit of 59 mV pH<sup>-1</sup> when the back gate dielectric is much thicker than the top dielectric. However, the thicker back-dielectric limits its application for ultrascaled point-of-care devices. In this work, we introduce and demonstrate a pH sensor, with WSe<sub>2</sub>(top)/ MoS<sub>2</sub>(bottom) heterostructure based double gated ISFET. The proposed device is capable of surpassing the Nernst detection limit and uses thin high-k hafnium oxide as the gate oxide. The 2D atomic layered structure, combined with nanometer-thick top and bottom oxides, offers excellent scalability and linear response with a maximum sensitivity of 362 mV pH<sup>-1</sup>. We have also used technology computer-aided (TCAD) simulations to elucidate the underlying physics, namely back gate electric field screening through channel and interface charges due to the heterointerface. The proposed mechanism is independent of the dielectric thickness that makes miniaturization of these devices easier. We also demonstrate super-Nernstian behavior with the flipped MoS<sub>2</sub>(top)/WSe<sub>2</sub>(bottom) heterostructure ISFET. The results open up a new pathway of 2D heterostructure engineering as an excellent option for enhancing ISFET sensitivity beyond the Nernst limit, for the next-generation of label-free biosensors for single-molecular detection and point-of-care diagnostics.

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#### **INTRODUCTION**

Measuring the ionic content or pH of an analyte is critical for a wide variety of applications, ranging from medical diagnostics, early disease detection, and genome sequencing to industrial process control, homeland security, and environmental monitoring<sup>1,2</sup>. Solid-state sensors, such as the ion-sensitive field-effect transistors (ISFETs), offer considerable benefits over state-of-theart electrochemiluminescence, electrochemical, and optical sensing platforms for point-of-care applications. ISFETs enable rapid, real-time, low-power, low-cost, and label-free electronic detection in a compact setup that is highly compatible with the standard complementary metal-oxide-semiconductor (CMOS) process<sup>3,4</sup>. In addition, ISFETs allow on-chip integration of sensor and measurement circuitry on the same substrate using monolithic integration<sup>5</sup>. Unlike metal-oxide-semiconductor (MOS) field-effect transistors (FETs), the gate metal in ISFETs is removed and replaced by a fluid gate which consists of a reference electrode capacitively coupled to the oxide layer through an electrolyte solution<sup>6</sup>. The pH sensing is based on the modulation of surface potential and channel conductance, when hydrogen ions in the electrolyte combine with hydroxyl groups on the gate dielectric, causing a shift in the threshold voltage<sup>7</sup>. This shift in threshold voltage for a unit change in pH is defined as the voltage sensitivity of the pH sensor<sup>8</sup>. For classical ISFET devices, the maximum achievable voltage sensitivity for pH sensing is 59 mV pH<sup>-1</sup> at 25 °C, known as the Nernst limit<sup>9</sup>. Enhancement of sensitivity, beyond the Nernst limit, is necessary to improve the resolution of the pH sensor, especially for detecting small pH changes in biological samples with ultralow concentrations. Several efforts have been made to achieve super-Nernstian sensitivity (>59 mV pH<sup>-1</sup>) in ISFETs. Recent studies have exploited the steep-subthreshold behavior in negative capacitance FETs<sup>10</sup>

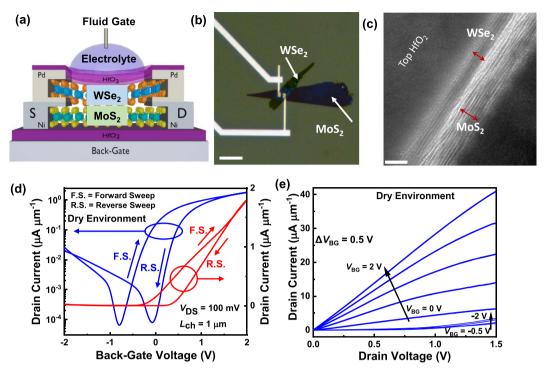
and tunnel FETs<sup>11</sup> to demonstrate ISFETs with sensitivities beyond the Nernst limit. A more common approach to enhance sensitivity uses capacitive amplification through asymmetric dual gate configurations<sup>9,12-16</sup>. Compared with classical single gate pH sensors, the combination of a fluid gate and back gate offers significant signal amplification, overcoming the charge screening effect caused by the counter ions in the electrolyte<sup>17</sup>. The threshold voltage shift is modulated by the top-  $(C_{top})$  and bottom- (C<sub>bottom</sub>) gate capacitances, allowing the sensitivity to be increased by simply adjusting the ratio,  $C_{\text{top}}/C_{\text{bottom}}^{18}$ . A simple approach to achieve super-Nernst sensitivity is by increasing the thickness of bottom dielectric or changing the dielectric constant (making  $C_{\text{bottom}} < C_{\text{top}}$ ). However, thicker dielectric limits miniaturization and restricts the use of these sensor devices in point-ofcare applications. Considering the ultimate aim of scaling down the device to nanoscale dimensions, an alternate method of enhancing sensitivity needs to be explored, which is independent of the thickness of the gate dielectric. This can be achieved by exploiting the interface traps and defects in order to modulate the capacitance, without (necessarily) increasing its device dimensions.

Several studies have reported ultrascaled FET-based biosensors using transition-metal dichalcogenide (TMDC)-based two-dimensional (2D) materials, like molybdenum disulfide (MoS<sub>2</sub>) or tungsten diselenide (WSe<sub>2</sub>), as the channel material<sup>19–23</sup>. Compared with silicon, the atomically thin 2D materials offer desired scalability, good electrostatics, and reduced flicker noise, leading to highly sensitive nano biosensors with excellent signal-to-noise ratio (SNR)<sup>24</sup>. Recently, FETs with Van der Waals WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) vertical heterojunction have been intensively investigated for their unique electrical properties such as band-to-band tunneling, negative differential resistance, and ambipolar

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**Fig. 1 Device structure and dry electrical characteristics. a** Schematic of the WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) ISFET device, **b** optical image, showing the WSe<sub>2</sub> on MoS<sub>2</sub> along with the metal contacts (scale bar 10  $\mu$ m), **c** HR-TEM of the channel region of HfO<sub>2</sub>/WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom)/HfO<sub>2</sub> stack (scale bar 5 nm), and **d** transfer and **e** output characteristics of the ISFET measured in the dry environment.

transfer characteristics, arising from the atomically sharp interface and strong interlayer coupling between the MoS<sub>2</sub> and WSe<sub>2</sub> films<sup>25-29</sup>. However, the use of TMDC 2D heterostructures has been limited to novel electronics, tunneling FETs, memory devices, photodetector, etc. In most of these applications, the nonidealities such as interface traps, defects, etc., deteriorate device performance. Alternatively, these nonideal effects can also be exploited and even engineered for sensing applications, as reported in recent literatures for super-Nernstian graphene-ISFETs<sup>30–32</sup>. In these devices, sensitivity above the theoretical Nernst limit is achieved through the process of charge-transfer doping, which directly modulates the carrier density of graphene with changes in pH level. However, this requires the fabrication of defectengineered nanocrystalline graphene grain boundaries, which may add to processing complexities. Here, we propose a WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) pH sensor, which uses the interface effects of Van der Waals heterostructure to surpass the Nernst sensitivity limit of 59 mV pH<sup>-1</sup> at room temperature. Unlike the defect-engineered graphene sensors, these still utilize an oxide to sense the pH by protonation or deprotonation. The sensing and signal transduction are decoupled using a double-gate FET architecture (fluid gate for sensing and back gate for gate voltage), while the charge screening due to the heterostructure interface is used to provide a gain over the Nernst limit.

In this work, we present experimental demonstration of a WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) vertical heterostructure ISFET for pH sensing beyond the Nernst sensitivity limit. The double-gated FET structure uses a thin high-k hafnium oxide (HfO<sub>2</sub>) as the top and bottom dielectric. The heterostructure enables the modulation of back gate transconductance by introducing a screening layer of interface and channel charges. This enables super-Nernst sensitivity by lowering effective back gate capacitance without increasing the device dimension and compromising the overall device performance. The atomically thin 2D heterostructure channel along with use of thin high-k dielectric allows extreme

device scalability that is crucial for next generation nanobiosensors, targeted for point-of-care applications. In addition, we have also used technology computer-aided design (TCAD) simulations to elucidate the working principle of our device and also to investigate pH sensors with the flipped  $MoS_2(top)/WSe_2(bottom)$  heterostructure as the channel material.

# **RESULTS AND DISCUSSIONS**

# **Device characterization**

Figure 1a shows the schematic of the ISFET device with the WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) vertical heterostructure. Few-layer flakes of exfoliated MoS<sub>2</sub> (bottom) and WSe<sub>2</sub> (top) were stacked together and sandwiched between two ~30 nm thick, HfO<sub>2</sub> layers, serving as the top and the bottom gate dielectrics. The bottom HfO<sub>2</sub> is deposited on a highly doped p-type silicon (Si) substrate, which also acts as a global back gate. Previous studies have reported pH sensors with HfO2 as the sensing surface due to the high pH sensitivity, low drift voltage, low hysteresis, and low body effect of HfO<sub>2</sub><sup>33–35</sup>. Moreover, the medium permittivity and high bandgap of HfO<sub>2</sub> provide excellent chemical stability when in contact with Si and silicon dioxide<sup>35</sup>. The MoS<sub>2</sub> layer in the proposed device was contacted by nickel (Ni) electrodes and showed n-channel operation<sup>36</sup>. Palladium (Pd) electrodes were used to make the contacts to the top WSe2 layer and shows p-channel (hole) transport<sup>37</sup>. The individual contacts on the flakes were also connected to form a single pair of source/drain contact. Figure 1b shows the optical image of the WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) heterostructure, with the source and drain metal contacts. Highresolution transmission electron microscope (HR-TEM) imaging was used to analyze the cross section of the fabricated device. The HR-TEM images shown in Fig. 1c show excellent interface quality with ~1.2 nm Van der Waals gap between HfO2-MoS2 and MoS2-WSe2 interfaces. This gap is expected due to the dry transfer of the exfoliated flakes. There is minimal gap in

WSe $_2$ -HfO $_2$  interface, as the top HfO $_2$  is deposited directly on WSe $_2$ . Both the top and bottom HfO $_2$  dielectric are also of excellent quality without any visible cracks or nonuniformities. The same was also confirmed by atomic force microscopy (AFM) on the fabricated device (on the channel after top HfO $_2$ ), which also showed good uniformity of deposition without any island formation (see Supplementary Fig. 1). The measured physical characteristics of the top and bottom HfO $_2$  films are provided in Supplementary Table 1. The top HfO $_2$  was slightly thicker and showed a slight increase in roughness due to process variations. A more detailed physical and surface analysis of the same HfO $_2$  films has been reported in the recent study by Ganapathi et al.  $^{38}$ 

The device was initially characterized in a dry environment by measuring the transfer and output characteristics as a function of back gate ( $V_{BG}$ ) and drain voltage ( $V_{DS}$ ), respectively. The transfer curve, in Fig. 1d, shows ambipolar characteristics, with p ( $V_{BG} \downarrow$  $I_{\rm D}\uparrow$ ) and n regions ( $V_{\rm BG}\uparrow I_{\rm D}\uparrow$ ) labeled. The transfer characteristics of the reverse sweep are considered to focus on the hole transport region. For  $V_{BG} < 0 \text{ V}$ , the p-type WSe<sub>2</sub> turns ON and starts conducting current through holes (p-channel operation), whereas for  $V_{BG} > 0 \text{ V}$ , n-type MoS<sub>2</sub> begins to conduct current by electrons (n-channel operation). The ON regions of these two materials are mutually exclusive, i.e., when one is ON and conducting, the other is fully depleted or even inverted. At large  $V_{\rm BG}$ , minority inversion carriers can form in one channel while the other is in accumulation. However, their contribution toward the overall device current is very small. The space charge regions at the p-n junctions of the hetero-interface, as well as the corresponding band offsets, pose a significant potential barrier for the minority carriers to overcome. Furthermore, measurements on individual back-gated MoS<sub>2</sub> and WSe<sub>2</sub> FETs show minimal currents in inversion mode due to the large Schottky barrier at their respective contacts (see Supplementary Fig. 2). Also, the lack of any negative differential resistance (NDR) eliminates the possibility of band-to-band tunneling in our devices<sup>28</sup>. Our findings are consistent with the results reported by Lee et al. for a similar WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) heterostructure<sup>39</sup>. The curves in Fig. 1e display the output characteristics of the proposed device at different  $V_{\rm BG}$ . The drain current increases with increasing  $V_{\rm BG}$  during n-channel operation (from  $V_{\rm BG}=0$  to 2 V), whereas it increases with decreasing  $V_{\rm BG}$  for p-channel operation (from  $V_{\rm BG}=-0.5$  to -2 V). The linearity of the curves indicates good (~ ohmic) contacts with a low Schottky barrier.

The current in the n-channel region is ~2 orders higher than the p-channel due to better back gate coupling for MoS<sub>2</sub>. The backgated transfer characteristics of the standalone WSe<sub>2</sub> FET (Supplementary Fig. 2) show much higher currents, highlighting weaker back gate control in the WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) heterostructure. The characteristics of the n-channel region are similar to that of a back-gated MoS<sub>2</sub> FET, showing minimal effects of heterostructure on the bottom layer. The reverse is valid for topgate coupling, which allows the WSe<sub>2</sub> layer to be better sensitive to pH changes during wet measurements. The bidirectional sweep of the  $V_{BG}$  reveals hysteresis in the device. The hysteresis direction confirms the lack of any ferroelectric polarization in the HfO<sub>2</sub> films. The clockwise direction of the hysteresis arises from interface traps in the HfO<sub>2</sub>-MoS<sub>2</sub> and MoS<sub>2</sub>-WSe<sub>2</sub> interfaces<sup>40</sup>. These interface traps play a key role in reducing the back gate transconductance without the need for a thicker bottom gate dielectric. The density of these interface traps  $(D_{it})$  is crucial for understanding the operation of the device. MOS capacitors are a popular way to extract interface characteristics. However, the small area of the exfoliated flakes (compared with the pads) prevents such measurements. The subthreshold slope (SS) can provide an approximation of  $D_{it}$  (higher  $D_{it}$  causes larger SS). But the complexity of the device structure and the transfer

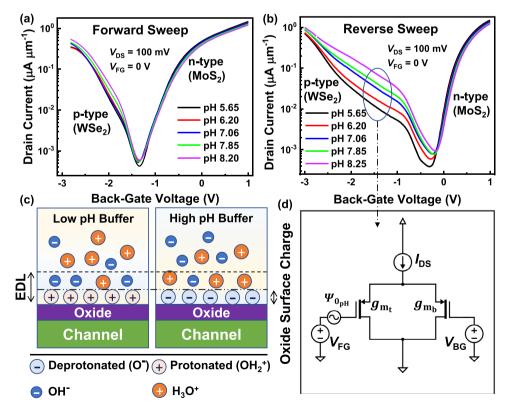


Fig. 2 Wet electrical characteristics and pH sensing. Transfer characteristics of the WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) ISFET at different pH values during **a** forward and **b** reverse sweep of the back gate voltage ( $V_{BG}$ ), **c** schematic of the pH sensing mechanism showing the formation of surface charge on oxide and electric double layer (EDL) in the electrolyte, and **d** simplified equivalent model showing signal transduction in a double-gated device.

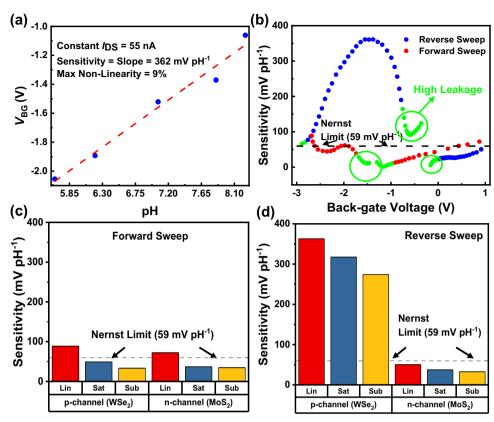


Fig. 3 Sensitivity. a Back-gate voltage ( $V_{BG}$ ) against pH at constant drain current, **b** Voltage sensitivity ( $S_V$ ) with respect to the  $V_{BG}$  of the transfer curve at pH = 7. Peak  $S_V$  at linear, saturation, and subthreshold region of operation during **c** forward sweep and **d** reverse sweep for the WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) ISFET.

characteristics make such approximations difficult. Hence, an analysis based on TCAD simulations is provided, which numerically estimates a  $D_{\rm it}$  of more than ~ $10^{12}\,{\rm cm}^{-2}$  at the MoS<sub>2</sub>–WSe<sub>2</sub> interface<sup>41</sup>.

# pH sensing with $WSe_2(top)/MoS_2(bottom)$ heterostructure ISFET

For the wet characterization and pH sensing, phosphate buffered saline (PBS) solutions of different pH values were prepared. These solutions were drop-casted onto the top  $HfO_2$  layer of the ISFET. A platinum wire, biased at 0 V, was used as the reference electrode ( $V_{FG}$ ). This reference electrode, in contact with the electrolyte solution, also serves as the fluid gate of the ISFET. The back gate transfer characteristics were obtained at multiple pH values, ranging from pH = 5.65 to 8.25, in both forward- and reverse-sweep directions, as shown in Figs. 2a and 2b, respectively.

The effective surface potential of the top dielectric is determined by the combination of fluid gate voltage, applied through the reference electrode, and the surface charge, modulated by the pH value of the electrolyte. A lower pH value results in protonation (OH + H<sup>+</sup> = OH<sub>2</sub><sup>+</sup>), giving rise to positive surface charges, while at higher pH values, the top HfO<sub>2</sub> surface is deprotonated (OH–H<sup>+</sup> = O<sup>-</sup>), yielding negative charges. These surface charges are partially screened by the counter-ions in the electrolyte, which move toward the charged surface and forms an electric double layer. The resultant surface charge or the change in potential at the oxide surface ( $\Psi_{0_{pH}}$ ) causes a corresponding shift in the device threshold voltage, which, in turn, determines the sensitivity of the pH sensor. A schematic of this sensing mechanism is shown in Fig. 2c.

Figure 2d represents a simplified model of the signal transduction once a surface charge and potential have developed at the oxide surface. A dual gated configuration can be represented as two transistors in parallel with a common source and drain. The top gate with a transconductance  $g_{\rm m_t}$  acts as the sensor for the oxide surface potential  $\Psi_{\rm 0_{pH}}$  (represented as an ac source) due to the pH. The fluid gate bias  $V_{\rm FG}$  is kept at a reference voltage of 0 V. The device current is controlled by the back gate voltage ( $V_{\rm BG}$ ) and has a transconductance of  $g_{\rm m_b}$ . When operated at constant current mode, the voltage sensitivity is given by Eq. 1. Hence, super-Nernstian sensitivity can be obtained if  $g_{\rm m_b} < g_{\rm m_t}$  providing a gain in the sensitivity  $^{12,42,43}$ .

$$S_{V} = \frac{\Delta V_{BG}}{\Delta pH} = \left(\frac{g_{m_{t}}}{g_{m_{b}}}\right) \left(\frac{\Delta \Psi_{0_{pH}}}{\Delta pH}\right) \tag{1}$$

To measure the sensitivity of our device, we plot  $V_{BG}$  against the corresponding pH of the electrolyte solution at a fixed drain current ( $I_{DS} = 55 \text{ nA}$ ), as shown in Fig. 3a. The voltage sensitivity (S<sub>v</sub>) is derived directly from the slope and amounts to 362 mV pH<sup>-1</sup>, clearly exceeding the Nernst limit by more than six times. The response, with a nonlinearity of 9%, indicates excellent sensing performance. The  $S_V$  can be tuned by adjusting the  $V_{BGV}$  as shown in Fig. 3b. Here,  $S_v$  is determined at a constant drain current and plotted with respect to the  $V_{BG}$  of the transfer curve at pH = 7 for both forward and reverse sweeps. The peak  $S_v$  in the subthreshold, linear, and saturation region of the WSe2(top)/ MoS<sub>2</sub>(bottom) ISFET is displayed in Figs. 3c and 3d for the forward and reverse sweeps, respectively. In the forward sweep,  $S_v$  is high only in the extreme linear regions of both p and n sides. On the other hand, the p-channel region in the reverse sweep shows high S<sub>v</sub> throughout subthreshold, linear, and saturation regions. In fact, the highest  $S_v$  (362 mV pH<sup>-1</sup>) is obtained in the reverse sweep when the back gate is biased in the linear region of hole transport. Although super-Nernstian sensitivity is visible in both the forward

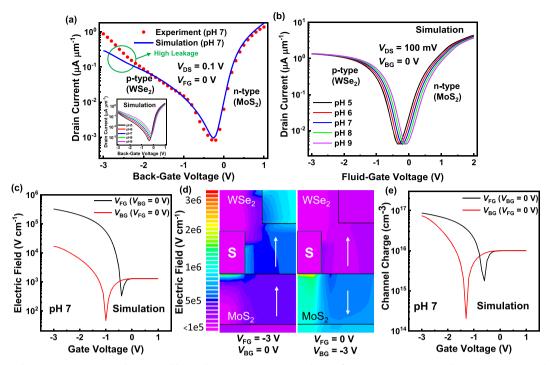


Fig. 4 TCAD simulations. a TCAD simulations calibrated against experimental transfer curve at pH = 7 in the reverse-sweep direction for the  $WSe_2(top)/MoS_2(bottom)$  ISFET. The inset shows simulation results for pH values ranging between 5 and 9. **b** Drain current vs fluid-gate voltage characteristics at different pH values, extracted using TCAD simulations. **c** Simulated magnitude of the electric field across the  $WSe_2$  channel v/s gate voltage (for  $V_{FG}$  and  $V_{BG}$  when the other is set to 0 V). **d** Simulated electric field contours with applied  $V_{FG}$  and  $V_{BG}$  (arrows indicate the field direction). **e** Simulated density of channel (WSe<sub>2</sub>) charge v/s gate voltage (for  $V_{FG}$  and  $V_{BG}$  when the other is set to 0 V).

and the reverse sweeps, it is more prominent in the reverse sweep direction. The enhanced sensitivity, beyond the theoretical Nernst limit of 59 mV pH $^{-1}$ , originates from the transconductance ( $g_{\rm m}$ ) modulation between the fluid-gate and the back gate, since  $g_{\rm mb} < g_{\rm mt}$ .

In the present work, the  $g_{\rm m}$  mismatch arises from the hetero-interface properties, whereas it originates from capacitance mismatch in devices with a thicker back gate dielectric<sup>44</sup>. The donor traps in the MoS<sub>2</sub>-WSe<sub>2</sub> interface capture holes from WSe<sub>2</sub> during the reverse sweep and reduces  $g_{\rm m_b}$ , leading to enhanced  $S_{\rm v}$ . In addition, the combination of screening effect<sup>45</sup> and mobility degradation<sup>46</sup> increases  $S_{\rm v}$  in the linear region of operation. In the former case, the inversion charges in the bottom MoS<sub>2</sub> channel lower the  $g_{\rm m_b}$  when the device is biased in strong linear region of the p-channel side. In addition, high vertical electric field degrades the mobility of carriers through the channel through surface scattering, reducing  $g_{\rm m_b}$  and improving  $S_{\rm v}$ . A similar effect is observed in the strong linear region of the n-channel region during the forward sweep. In all cases, the lowest  $S_{\rm v}$  is observed in the subthreshold region, primarily due to the high gate leakage.

#### Charge-screening effect for high sensitivity

To further understand the role of screening and interface traps in enhancing  $S_{v,}$  we investigate the electrical characteristics of  $WSe_2(top)/MoS_2(bottom)$  heterostructure ISFET using TCAD. All simulations were carried out with Silvaco ATLAS device simulation tool, which self-consistently solves Poisson's equations along with current continuity equations to obtain the drain current as a function of gate voltage. Silvaco ATLAS TCAD device simulator tool cannot directly model devices based on 2D materials. As a workaround, we have modeled  $MoS_2$  and  $WSe_2$  using the material parameters available in the literature and considering the 3-D equivalent of their 2-D density of states. Using a few layers (~6 nm) of exfoliated flakes means that the quantum effects

dominant in monolayer or bilayer materials can be avoided. This approach is consistent with previous studies which used TCAD tools to simulate 2-D material based FETs, with reasonable accuracy<sup>41,47–49</sup>. The mobile ions in the electrolyte solution are modeled as charge carriers in a semiconductor material such that pH values correspond to the charge density at the electrolyte/ dielectric interface<sup>50</sup>. The use of HfO<sub>2</sub> as the gate dielectric and sensing layer has been well explored in the literature and usually shows near-ideal sensitivity even in single gate configurations<sup>15,33,35</sup>. Hence, the modeling of surface protonation or deprotonation on the HfO2 surface is not included in the simulations. The TCAD model is calibrated against the drain current vs back gate voltage curve of pH = 7 in the reverse-sweep direction. Figure 4a shows a good agreement between experiment and simulation. A slight deviation is observed at high  $V_{BG}$ due to gate-leakage in the practical device, which is not captured in the TCAD simulations. The inset of Fig. 4a shows the simulation results for a wide range of pH values.

Figure 4b shows the variation in the drain current with respect to the change in the fluid-gate voltage ( $V_{\rm FG}$ ) at different pH values, when  $V_{\rm BG}=0\,\rm V$ , as obtained from the TCAD simulations. The sensitivity with change in  $V_{\rm FG}$  is found to be constant across all operating regions and is ~59.5 mV pH<sup>-1</sup>, very close to the Nernst limit. This is expected as there is no gain mechanism with a single gated device. However, the device can still provide near-ideal sensitivity even when operated as a single gated ISFET.

The TCAD simulations confirm the presence of interface traps ( $\sim 10^{12} \, \text{cm}^{-2}$ ) at the MoS<sub>2</sub>–WSe<sub>2</sub> interface as well as inversion carriers in the bottom MoS<sub>2</sub> channel (Supplementary Fig. 3). The higher density of traps in WSe<sub>2</sub>–MoS<sub>2</sub> interface and their distribution of energy levels have a profound impact on the device operation. The much lower currents and larger subthreshold slope of the p-channel region are due to these donor-interface traps that capture holes from the WSe<sub>2</sub> channel as back gate bias is applied. On the contrary, the traps in the HfO<sub>2</sub>–MoS<sub>2</sub>

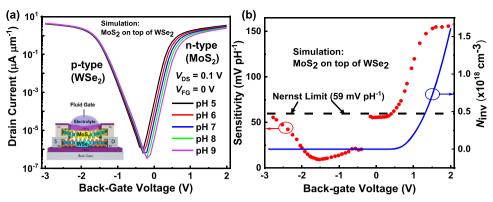


Fig. 5 Simulated MoS<sub>2</sub>(top)/WSe<sub>2</sub>(bottom) ISFET. a Drain current vs back-gate voltage characteristics at different pH values, when  $V_{FG} = 0 \text{ V}$ . b Voltage sensitivity and density of inversion carriers ( $N_{inv}$ ) in the bottom channel against back-gate voltage.

interface are shallower and of lower density, with minimal effects on the n-channel region. The donor traps in WSe<sub>2</sub>-MoS<sub>2</sub> interface effectively weaken or screen the controlling of the WSe<sub>2</sub> channel charges by the applied back gate voltage. In addition, the inversion charges that form in the bottom MoS<sub>2</sub> layer do not contribute significantly to device current (due to band offset barriers) but can again weaken the back gate control by screening.

To illustrate the effects of screening, parameters such as electric field and charge are plotted against the back gate and fluid gate voltages when the other is set to 0 V (Fig. 4(c-e)). Figure 4c shows the magnitude of the electric field across the WSe<sub>2</sub> channel, which is an order of magnitude lower with the back gate bias than with the applied fluid-gate voltage. The distribution of the electric field is illustrated in Fig. 4d for  $V_{BG} = -3 \text{ V}$  ( $V_{FG} = 0 \text{ V}$ ) and  $V_{FG} = -3 \text{ V}$  $(V_{BG} = 0 \text{ V})$ . With the back-gate voltage, the field lines are mainly confined within MoS<sub>2</sub>, whereas with the applied fluid-gate voltage, there is much higher field penetration. This is evident by considering the electric field direction in the respective channels. Even at  $V_{BG} = -3 \text{ V}$  ( $V_{FG} = 0 \text{ V}$ ), the field in the WSe<sub>2</sub> channel is directed away from the back gate, i.e., there is a certain amount of band pinning at the WSe2-MoS2 interface due to the interface traps. However, this does not occur when  $V_{\rm FG} = -3 \, {\rm V}$  $(V_{BG} = 0 \text{ V})$ . These effects are also visible in the charge induced in the WSe<sub>2</sub> channel by the respective gates, which is shown in Fig. 4e. The channel charge in WSe<sub>2</sub> induced by the back-gate bias is much lower than the channel charge formed by the fluid-gate for a considerable range of voltages. These voltages correspond to the range in  $V_{BG}$  where the super-Nernstian response is visible. These factors emphasize the role of screening in reducing the back-gate coupling and the necessity of interface and inversion charges to cause it. This reduced back-gate coupling and hence reduced  $g_{m_b}$  directly translates to the high sensitivity in our heterostructure-based ISFET.

#### MoS<sub>2</sub>(top)/WSe<sub>2</sub>(bottom) pH sensor

Next, we simulate the flipped  $MoS_2(top)/WSe_2(bottom)$  heterostructure device, using the calibrated TCAD model, with  $MoS_2$  (n-type) stacked on top of  $WSe_2$  (p-type). To be consistent with the simulation setup for  $WSe_2(top)/MoS_2(bottom)$  ISFET, the same donor trap density has been used, while the acceptor traps are being ignored. Although, in a real device, the presence of acceptor traps will reduce the current in the n-region, having an identical simulation setting provides a fair comparison between the  $MoS_2(top)/WSe_2(bottom)$  and the  $WSe_2(top)/MoS_2(bottom)$  ISFETs. Figure 5a shows the simulated drain current of the flipped  $MoS_2(top)/WSe_2(bottom)$  ISFET, plotted against  $V_{BG}$  for different pH values at  $V_{FG} = 0$  V. Compared with the  $WSe_2(top)/MoS_2(bottom)$  ISFET, the drain current obtained is higher in the  $MoS_2(top)/WSe_2(bottom)$  ISFET. This is because of the lower Schottky barrier

height in  $MoS_2$  and the absence of acceptor traps in the interface in the simulation setup. The  $S_{\rm v}$  of the simulated  $MoS_2$ (top)/ $WSe_2$ (bottom) device, shown in Fig. 5b, displays super-Nernstian sensitivity when the n-layer is ON. On the contrary, the sensitivity remains within the Nernst limit in the p-region. The density of inversion electrons ( $N_{\rm inv}$ ) in the  $WSe_2$  layer, with change in  $V_{\rm BG}$ , is also displayed in Fig. 5b. A higher  $N_{\rm inv}$  corresponds to an enhanced  $S_{\rm V}$ , confirming the contribution of the screening effect in improving sensor performance. Mobility degradation due to higher vertical field at higher  $V_{\rm BG}$  may also play a role in enhancing  $S_{\rm V}$ .

In summary, we have demonstrated a double gated super-Nernstian ISFET for pH sensing, using a 2D heterostructure of WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) as the channel material. The proposed device shows a maximum sensitivity of 362 mV pH<sup>-1</sup> for pH sensing due to the reduced back-gate transconductance. The novelty of this work lies in the underlying sensing mechanism that is attributed to the traps and defects formed at the heterointerface. The bottom semiconducting channel acts both as a screening layer and a site for interface traps, contributing to the enhanced sensitivity. The proposed device and its flipped variant (MoS<sub>2</sub>(top)/WSe<sub>2</sub>(bottom)) both offer high sensitivity with great tunability, depending on the region of operation (subthreshold, linear, and saturation). Although the Nernst limit can be surpassed in other dual-gate ISFETs by tailoring the thickness and the material of the top and bottom gate dielectrics, the thick back gate oxide limits device scaling. The proposed device, on the other hand, with thin high-k gate oxides and atomically thin 2-D channel materials, offers extreme scalability, enabling ultrascaled, next-generation biosensing platforms for point-of-care applications.

# **METHODS**

## Device fabrication and material characterization

In all, 30-nm-thick HfO<sub>2</sub> was deposited (e-beam evaporation) on p++ Si substrate, which is used as a global back gate. The HfO2 thin film deposition parameters were optimized based on the reports by Ganapathi et al.<sup>38</sup>. MoS<sub>2</sub> flakes (source: Graphene supermarket) were mechanically exfoliated onto this substrate using scotch-tape and polydimethylsiloxane (PDMS) stamps. The samples were cleaned with acetone, isopropyl alcohol (IPA), and deionized water after each process. Flakes, ranging between 4 and 7 nm in thickness, were identified optically and later confirmed using AFM (Bruker) (Supplementary Fig. 4). Source/drain contacts with a channel length of 1 µm were patterned using e-beam lithography (Raith Pioneer), and nickel/palladium (nickel in contact with MoS<sub>2</sub>) of thickness 15 nm each were deposited using e-beam evaporation, followed by liftoff. WSe<sub>2</sub> flakes (source: 2D Semiconductors) of 4-7 nm thickness were transferred on top of the MoS<sub>2</sub> flake using PDMS stamps after aligning them using the stage of an optical lithography setup (MJB4 mask aligner). Source/drain contacts of palladium/gold, each with a thickness of 15 nm, were



connected to the WSe $_2$  layer, aligned with the contacts on MoS $_2$ . The structural and chemical integrity of the flakes were verified using Raman spectroscopy (Supplementary Fig. 4). HfO $_2$  (~30 nm) was deposited on top, forming the top dielectric, which acts as the pH sensing layer. This layer also isolates the electrodes from direct contact with the electrolyte. Oxide on the contact pads was etched using buffered hydrofluoric acid (BHF) solution. For the cross-section HR-TEM of the device, the sample was prepared using focused ion beam (FIB) Helios G4-UX and imaged using Titan Themis 300 kV from ThermoFisher.

#### **Device characterization**

PBS was used as the electrolyte. PBS solution with different pH values, ranging from pH = 5 to 8, was prepared and the pH was accurately determined using a standard laboratory pH meter (Eutech Instruments). The electrolyte solution is then drop-casted ( $\sim\!10~\mu$ l) onto the exposed region (top HfO $_2$  layer) of the ISFET for sensing. A fluid gate, made up of a platinum electrode, connects to the top HfO $_2$  layer through the electrolyte solution, and is used to maintain a constant reference potential during measurement. Finally, the electrical measurements were carried out using Keysight B1500 Semiconductor device analyzer at room temperature.

# **TCAD** simulations

The current–voltage characteristics of WSe<sub>2</sub>(top)/MoS<sub>2</sub>(bottom) and MoS<sub>2</sub>(top)/WSe<sub>2</sub>(bottom) ISFET heterostructure devices were simulated using Silvaco ATLAS simulator<sup>51</sup>. Material parameters of MoS<sub>2</sub> and WSe<sub>2</sub>, such as carrier mobility, were defined based on the reports by Mirabelle et al.<sup>47,48</sup> and Arora et al.<sup>52</sup>. Shockley–Read–Hall (SRH) and Auger Quantum effects have been ignored in the TCAD to simplify the simulation procedure. The electrolyte was modeled using an undoped semiconductor. The details of this modeling are given in Supplementary Notes.

The simulated ISFET was then calibrated against the experiment–transfer curve, corresponding to pH = 7. The device parameters such as mobility, doping, interface traps, and Schottky barrier heights were adjusted in the TCAD model to match the experimental results accurately. These final values after calibration of these parameters are provided in Supplementary Table 2, along with the material parameters used in the simulation. To better fit the subthreshold region of the transfer curve and its shape, interface traps and fixed–oxide charges were added at various interfaces. These are summarized in Table 3 of the Supplementary Information.

# **DATA AVAILABILITY**

The data that support the findings of this study are available from the corresponding author upon reasonable request.

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#### **AUTHOR CONTRIBUTIONS**

N. Bhat conceived the project. S. Sanjay and M. Hossain fabricated the devices, performed the characterizations, and conducted the experiments. A. Rao carried out TEM sample preparation, characterization, and analysis. S. Sanjay carried out the TCAD simulations with inputs from M. Hossain. S. Sanjay and M. Hossain analyzed the results and prepared the paper with N. Bhat. All authors approve the final version of the paper.

#### **COMPETING INTERESTS**

The authors declare no competing interests.

#### ADDITIONAL INFORMATION

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