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## OPEN Bi<sub>0.9</sub>Ho<sub>0.1</sub>FeO<sub>3</sub>/TiO<sub>2</sub> Composite Thin Films: Synthesis and Study of **Optical, Electrical and Magnetic Properties**

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A visible light active Bi<sub>0.9</sub>Ho<sub>0.1</sub>FeO<sub>3</sub> nanoparticles/TiO<sub>2</sub> composite thin films with different mol.% of Bio 9Hoo 1 FeO3 were successfully prepared via non-aqueous sol-gel method. The incorporation of 5, 10 and 20 mol.% Bio 9 Hoo 1 FeO3 nanoparticles in the precursor solution of TiO2 brings modifications in the functional properties of the composite thin films. XPS analysis indicates that interdiffusion of Fe<sup>3+</sup>, Ho<sup>3+</sup>, Bi<sup>3+</sup>/Ti<sup>4+</sup> ions through the interfaces between Bi<sub>0.9</sub>Ho<sub>0.1</sub>FeO<sub>3</sub> nanoparticles and TiO<sub>2</sub> matrix reduces the concentration of  $Ti^{3+}$  ions. X-ray diffraction analysis affirms that  $TiO_2$  and  $Bi_{0.9}HO_{0.1}FeO_3$ retain anatase and orthorhombic phase respectively in composite films. The composite thin film containing 20 mol.% Bio 9 Hoo 1 FeO 3 nanoparticles exhibits the most prominent absorption phenomenon in visible region and has significantly reduced indirect band gap of 2.46 eV compared to that of pure TiO<sub>2</sub> (3.4 eV). Hall effect measurements confirm that the resistivity of composite film increases by  $\sim$ 2.33 orders of magnitude and its carrier concentration decreases by 1.8 orders of magnitude at 5 mol.% Bi<sub>0.9</sub>Ho<sub>0.1</sub>FeO<sub>3</sub> nanoparticles addition compared to those of pure TiO<sub>2</sub> film. Moreover, the pure film exhibits diamagnetism, whereas the composite films have both large ferromagnetic and small diamagnetic components. The findings in this research justify that the composite film can be a potential candidate for making improved photocatalyst, resistors and spintronic devices.

TiO<sub>2</sub> as powder or thin film has been comprehensively investigated due to their excellent photo-chemical stability, low cost and non -toxicity<sup>1</sup>. However, pure TiO<sub>2</sub> responds only to ultraviolet (UV) light which comprises only 4% of the total sunlight. The absorption occurs in the UV light due to large band gap (3.2 eV) which eventually restricts its practical application in degradation of organic pollutants such as dyes, detergents and pesticides etc1-4. Because of diamagnetism and low resistivity, TiO<sub>2</sub> has become a concern in the field of spintronics, resistors, sensors and multistate memory devices<sup>5,6</sup>. These limitations have galvanized endeavors to enhance photocatalytic activity, induce magnetism and raise resistivity of TiO<sub>2</sub>.

One of the attempts is to substitute a portion of the titanium with different electropositive atoms. Transition metals (e.g. Fe, Mn and Cr etc) are the most commonly used constituents that induce visible light absorption<sup>2-4</sup>. Besides, transition metal (TM) doped  $TiO_2$  has exhibited improved ferromagnetism at room temperature<sup>7,8</sup>. Moreover, Mn and Fe dopants in TiO<sub>2</sub> have increased the resistivity significantly<sup>7,8</sup>. Another endeavor involves substitution of oxygen with electronegative atoms (N, C and S) which aids TiO<sub>2</sub> to absorb visible light efficiently<sup>9</sup>. However, the introduction of substitutional elements creates scattering centers which decrease the photochemical activity. It has been reported that substitutional TM serves as a center for charge carrier recombination<sup>2</sup>. Besides, it has been controversial as to whether Co or Fe doped TiO2 is a true magnetic semiconductor or if the effect is due to the clustering of a ferromagnetic second phase<sup>10,11</sup>. Indeed, according to these reports doped  $TiO_2$  has failed to exhibit magnetism to any greater extent.

Since, the physical quantities like electrical, magnetic, and optical properties come in the category of sum or product properties, many have tried to make composites of TiO<sub>2</sub> to utilize this criteria<sup>12,13</sup>. Magnetic properties

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have been improved by the fabrication of  $TiO_2/Fe_3O_4/SiO_2$  or  $Fe_xO_y$ - $TiO_2$  composites<sup>14,15</sup>. Similarly the nanostructured ZnO/TiO<sub>2</sub> has been found to improve the optical properties, even though, ZnO is a wide band gap semiconductor<sup>16</sup>. Nowadays, the visible light activity of  $TiO_2$  is increased by coupling it with other narrow bandgap semiconductors. The proposition is that coupled semiconductors form a hetero-junction structure which can transfer electrons from an excited small band gap semiconductor to the other attached one in the case of proper band potentials.  $In_2O_3^{17}$ ,  $BiFeO_3^{18}$  or CdSe co-sensitized  $TiO_2^{19}$  heterostructures have been reported for showing better visible light activity. Multiferroic  $BiFeO_3$  (BFO) deserves a special mention for its narrow band gap (2.2 eV) and chemical stability which have made it a suitable visible- light responsive photocatalyst<sup>1</sup>. The coupling of  $TiO_2$  with BFO has been accomplished through different synthesis processes. For instance, core-shell structured  $BFO/TiO_2$  by hydrolysis precipitation approach<sup>1</sup>, deposition of  $TiO_2$  on ferroelectric BFO substrate by pulsed laser deposition<sup>20</sup>, growth of  $TiO_2$  nanofibers on BFO nanoparticles by electrospinning<sup>18</sup>. These reports focus on enhancing visible light activity of  $TiO_2$ . Subsequently, several studies ponder on improving electrochemical energy storage capacity and solar energy conversion efficiency of titania by anchoring BFO nanoparticles on  $TiO_2^{21,22}$ .

The magnetic and electrical properties (e.g. resistivity, carrier concentration and mobility) still remain unreported with BFO nanoparticles due to its various limitations. Literature survey indicates that BFO nanoparticles are antiferromagnetic<sup>23</sup> and have higher leakage current due to formation of oxygen vacancies  $(OVs)^{24}$ . Recently, doping with tri-positive rare-earth ions  $(R^{3+})$  at the Bi site has been proposed to mitigate some of the issues mentioned above<sup>25</sup>. Previous reports suggest that 10 mol.% Holmium  $(Ho^{3+})$  doping at Bi site increases the magnetization by 6 times and exhibits higher resistivity compared to pure BFO<sup>25</sup>. Since Ho<sup>3+</sup> is a rare earth metal it is expected that  $Bi_{0.9}Ho_{0.1}FeO_3$  (BHFO) nanoparticles will have better absorbance in the visible region and lower band gap than BFO<sup>23</sup>. The aforementioned advantages make BHFO nano particles a good candidate for incorporation in TiO<sub>2</sub>.

 $TiO_2$  thin film is better candidate for spintronics, sensors and magnetic memory devices creating faster, smaller and more energy efficient devices compared to titania nanoparticles, nanowire, nanorod or nanofiber. Both undoped and doped  $TiO_2$  thin films have been of immense interest due to their enhanced photocatalytic activity, better electrical properties and room temperature ferromagnetism<sup>3,7,26,27</sup>. To the best of our knowledge the optical, magnetic and electrical properties based on incorporation of BHFO nanoparticles in  $TiO_2$  thin film have not been explored yet. Herein motivated by the above concerns, BHFO nanoparticle/ $TiO_2$  composite thin films were prepared by non-aqueous sol-gel method and their structural, magnetic, optical and electrical properties were investigated in detail. The current study revealed that composite thin films could be a more efficient visible light absorber than pure  $TiO_2$ . Furthermore room temperature magnetism and improved electrical properties of this composite thin films have been found.

#### **Experimental Section**

**Materials.** Bismuth nitrate pentahydrate (Bi(NO<sub>3</sub>)<sub>3</sub>.5H<sub>2</sub>O, Merck -India), iron nitrate nonahydrate (Fe(NO<sub>3</sub>)<sub>3</sub>.9H<sub>2</sub>O, Merck-India), holmium nitrate pentahydrate (Ho(NO<sub>3</sub>)<sub>3</sub>.5H<sub>2</sub>O, Sigma Adrich-USA), citric acid (C<sub>6</sub>H<sub>8</sub>O<sub>7</sub>, Merck -India) and ethylene glycol (C<sub>2</sub>H<sub>6</sub>O<sub>2</sub>, Merck -India) were used for the synthesis of BHFO nanoparticles. In this case, C<sub>6</sub>H<sub>8</sub>O<sub>7</sub> acted as the chelating agent to complex the metal cations and C<sub>2</sub>H<sub>6</sub>O<sub>2</sub> was added as polymerization agent<sup>23</sup>.

The solutions of pure TiO<sub>2</sub> thin film and BHFO nanoparticles/TiO<sub>2</sub> composite thin films were prepared using titanium(IV) n-butoxide (Ti (O-nBu)<sub>4</sub>, Aldrich-USA) as precursor and n-butyl alcohol (n-BuOH, Merck-India) as solvent. Besides, acetylacetone ( $C_5H_8O_2$ , Loba chemie) was added in the solution as chelating agent to decrease the reactivity of Ti(O-nBu)<sub>4</sub><sup>28</sup>. Glacial acetic acid (CH<sub>3</sub>COOH, Qualikems) was used further to begin hydrolysis via an esterification reaction.

**Synthesis of BHFO nanoparticles.** BHFO nanoparticles were synthesized using a modified sol-gel method described elsewhere<sup>23</sup>. Shortly, for a typical 1 g BHFO powder synthesis process stoichiometric proportion of Bi(NO<sub>3</sub>)<sub>3</sub>.5H<sub>2</sub>O (0.003 mol), Fe(NO<sub>3</sub>)<sub>3</sub>.9H<sub>2</sub>O (0.0033 mol), Ho(NO<sub>3</sub>)<sub>3</sub>.5H<sub>2</sub>O (0.00033 mol), C<sub>6</sub>H<sub>8</sub>O<sub>7</sub> (0.0067 mol) and C<sub>2</sub>H<sub>6</sub>O<sub>2</sub> (10 ml) were dissolved in 400 ml deionized water. Consecutively the solution was heated under continuous stirring at 75–85 °C for 4 h to obtain gel. The gel was dried at 100 °C for 24 h in a drier to obtain precursor xerogel. The ground precursor xerogel powders were annealed at 500 °C for two hours with a heating rate of 3 °C/min to obtain BHFO nanoparticles. BHFO pellets (thickness 0.11 cm and diameter = 1.3 cm) were prepared by mixing precursor xerogel powders with PVA binder followed by pressing (5 tons pressure) and annealing at 500 °C with a heating rate of 3 °C/min to measure electrical properties.

**Preparation of pure TiO**<sub>2</sub> thin films and BHFO nanoparticles/TiO<sub>2</sub> composite thin films. The detail synthesis process of pure TiO<sub>2</sub> thin film is delineated elsewhere<sup>28</sup>. Briefly, first n-BuOH (0.0884 mol) and  $C_5H_8O_2$  (0.0015 mol) were mixed and, then Ti(O-nBu)<sub>4</sub> (0.005 mol) was added to the solution. Subsequently this mixture was stirred for 30 min at room temperature. CH<sub>3</sub>COOH (0.001 mol) was slowly added into the alkoxide solution and stirred continuously for another 30 min. Concentration of the final solution was 0.5 M and its color was yellowish.

Thin films containing 5 mol.% BHFO (marked as  $T^{95}B^5$ ), 10 mol.% BHFO (marked as  $T^{90}B^{10}$ ) and 20 mol.% BHFO (marked as  $T^{80}B^{20}$ ) were prepared using non aqueous sol-gel method. To synthesize composite films, molar amount of n-BuOH,  $C_5H_8O_2$ , Ti(O-nBu)<sub>4</sub> and CH<sub>3</sub>COOH was kept similar as that of pure thin film. Firstly 0.0442 mol of n-BuOH was taken and calculated amount of BHFO nanoparticles was mixed into it for every composite film. The whole mixture was dispersed vigorously in an ultrasonic bath for 30 min to avoid the aggregation of BHFO nanoparticles. Then  $C_5H_8O_2$ , Ti(O-nBu)<sub>4</sub> and rest of n-BuOH (0.0442 mol) were added into this mixture and it was stirred for 30 min on a stirring plate at room temperature. This mixture was then taken into ultrasonic



**Figure 1.** (a) XPS survey spectra of BHFO nanoparticles,  $TiO_2$  film and  $T^{80}B^{20}$  composite thin film. (**b**,**c**) Ti 2p spectra of  $TiO_2$  and  $T^{80}B^{20}$ . (**d**-**f**) O 1s spectra of BHFO,  $TiO_2$  and  $T^{80}B^{20}$ .

bath in which CH<sub>3</sub>COOH was added dropwise. The final mixture was under intensive stirring for 2 h in an ultrasonic bath. The prepared mixture eventually turned into brownish color.

The brownish and yellowish solutions were spin-coated onto the  $2.5 \times 2 \text{ cm}^2$  glass substrate at 2000 rpm for 30 s and dried at 200 °C for 10 mins. During coating brownish solution was in stirring state continuously so that BHFO nanoparticles did not settle down. Micro pipette was used to ensure that specific amount of solution was taken every time. The spin-coating and drying processes were repeated three times. The pure and composite films were annealed in static air at 500 °C for 2 h with a heating rate of 3 °C/min.

**Characterization.** XPS measurements were carried out on the PHI Quantera II spectrometer and peak fitting was performed by Origin Pro. XRD (PANalytical Empyrean X-Ray Diffractometer system) was used for phase analysis of nanoparticles and thin films utilizing a Cu x-ray source (wavelength:  $K\alpha_1 = 1.540598$  Å and  $K\alpha_2 = 1.544426$  Å). Rietveld refinement was done using FULLPROF software<sup>29</sup>. Field emission scanning electron microscope (FESEM: JEOL, JSM, 7600F) was employed to monitor the morphologies of nanoparticles and thin films and to determine the thickness of these thin films. Energy dispersive X-ray spectroscopy (EDS) was used for mapping the elements present in the thin films. The optical properties were evaluated using UV/Vis/NIR spectrometric measurements (Perkin Elmer, Lambda 1050). To measure the magnetic property of the thin films, 5 mm × 5 mm size sample was cut from every thin film substrate. Electrical characterizations of nanoparticles were carried out by coating both sides of the pellet with silver paste and its conductivity was measured using Precision Materials Analyzer (Radiant Technologies, Inc.: P-PMF, PMF0215-377). Hall measurement was carried out using the van der Pauw configuration (HMS-3300). The four points contacts were made by indium tin alloy soldering which enabled us to obtain the carrier concentration, resistivity and mobility of thin films. The room temperature magnetic property was determined by vibrating sample magnetometer (VSM: EV-9 Microsense).

#### **Results and Discussion**

The chemical composition and electronic structure of the as synthesized BHFO,  $TiO_2$  and  $T^{80}B^{20}$  was determined by performing XPS analysis. Figure 1(a) depicts the XPS survey spectrum which clearly shows the presence of all the main constituent elements like Ti, Bi and O. Here, all the spectra were corrected for the C 1s peak appeared at 284.6 eV. However, Fe 2p peaks near 720 eV were not easily detectable by XPS nonetheless the EDX analysis (see Fig. S4) suggests the abundance of Fe in composite thin films. The low sensitivity of Fe to XPS analysis may lead to its absence in the survey spectra. However, the high resolution XPS scanning of BHFO nanoparticle depicts the peak corresponding to Fe having low intensity compared to other elements see Fig. S1. The high resolution elementary XPS peaks of Ti 2p and O 1s core levels observed for TiO<sub>2</sub>,  $T^{80}B^{20}$  and BHFO after curve- fitting are presented in Fig. 1(b–f). The binding energies of the Ti  $2p_{3/2}$  and Ti  $2p_{1/2}$  are found to be nearly 458.4 and 464.1 eV respectively, giving a spin orbit splitting of ~5.7 eV between the two core levels, which corresponds to the 4+ oxidation state of Ti in anatase phase<sup>30,31</sup>. Moreover, the appearance of extra peaks at around 457.8 eV and 463.3 eV



Figure 2. XRD patterns of the prepared pure film, composite films and nanoparticles.

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confirms the presence of  $Ti^{3+31}$ . The peak position of Ti 2p in  $T^{80}B^{20}$  has shifted negligibly. It is important to note here that the peak areas ratio for  $Ti^{3+}$  to  $Ti^{4+}$  given by  $(Ti^{3+} 2p_{3/2} + Ti^{3+} 2p_{1/2})_{Area}/(Ti 2p_{3/2} + Ti 2p_{1/2})_{Area}$  in  $TiO_2$  and  $T^{80}B^{20}$  are 0.348 and 0.231 respectively. The diminishment of the area ratio in  $T^{80}B^{20}$  implies that the concentration of  $Ti^{3+}$  has decreased. The interdiffusion of  $A^{3+}$  (here  $A^{3+}$  represents  $Bi^{3+}$ ,  $Fe^{3+}$  and/or  $HO^{3+}$ ) from BHFO to  $TiO_2$  across the interface could lead to a reduction in the concentration of  $Ti^{3+}$  ion. The substitution of  $Ti^{4+}$  in  $TiO_2$  by  $A^{3+}$  generates holes to maintain charge neutrality which oxidizes an existing  $Ti^{3+}$  to  $Ti ^{4+}$  according to the following reactions decreasing the density of electron hopping sites.

$$A^{3+} \xrightarrow{IIO_2} A'_{Ti} + h^{\bullet} \tag{1}$$

$$i^{3+} + h^{\bullet} = Ti^{4+}$$
 (2)

Combining Equations (1,2)

$$A^{3+} + Ti^{3+} \xrightarrow{TiO_2} A'_{Ti} + Ti^{4+}$$
(3)

Figure 1(d) displays the XPS spectra of O 1s for BHFO, which can be deconvoluted into two peaks at around 529.4 eV and 531 eV. The lower binding energy peak is affiliated with the intrinsic O 1s core spectra and the higher energy peak is attributed to the oxygen vacant sites in BHFO. The XPS spectrum of O 1s core level for TiO<sub>2</sub>, as shown in Fig. 1(e), can also be de-convoluted into two symmetric Gaussian peaks centered at 529.6 and 532.2 eV<sup>22</sup>. The low binding energy peak at 529.6 eV of O 1s can be assigned to the -2 oxidation state of oxygen<sup>31</sup>. The enhancement of peak intensity at around 532 eV can be ascribed to the presence of both the oxygen vacancy related defects (OV) and the hydroxyl group (OH<sup>-</sup>) together in T<sup>80</sup>B<sup>20</sup>2<sup>2,31</sup>. The appearance of OH<sup>-</sup> in TiO<sub>2</sub> is normal since it is potential photocatalyst<sup>2,4</sup>. It is noteworthy here that the ratio of the two peak areas (OV/OH<sup>-</sup>)<sub>Area</sub>/(O 1s)<sub>Area</sub> in TiO<sub>2</sub> and T<sup>80</sup>B<sup>20</sup> are 1.83 and 2.20 respectively. The increment in area ratio could be due to an increment in OH<sup>-</sup> group concentration in T<sup>80</sup>B<sup>20</sup>. This indicates that the composite film could be a better candidate for making improved photocatalyst than many other ceramic thin films.

Figure 2 shows room temperature XRD patterns of nanoparticles, pure thin film and composite thin films. The obtained pattern of TiO<sub>2</sub> as depicted in Fig. 2(a) shows satisfactory correspondence with anatase phase. The diffraction peaks at  $2\theta = 25.2^{\circ}$  (101), 37.8° (004), 48.0° (200), 53.9° (105) and 55.1° (211)) are matched well with the standard diffraction data of tetragonal crystal structure of TiO<sub>2</sub><sup>32</sup>. Thus, it is confirmed that the annealed TiO<sub>2</sub> film deposited on glass substrates has anatase phase with no trace of other polymorphs of TiO<sub>2</sub> (e.g., rutile and brookite).

On the other hand, XRD pattern of BHFO as shown in Fig. 2(e) is in line with previous report of Ho<sup>3+</sup> doped BFO ceramics<sup>33</sup>. In addition to the desired spectra one extra peak [marked by asterisk (\*)] becomes visible which indicates the presence of impurity phase. To uncover the crystal structure and the impurity phase, Rietveld refinement of BHFO was carried out using the FULLPROF software<sup>29</sup>. The refinement corroborates that BHFO crystallizes in a orthorhombically distorted perovskite structure having Pnma space group and the extra peak in the XRD pattern belongs to Bi<sub>2</sub>Fe<sub>4</sub>O<sub>9</sub> impurity phase. However, refinement indicates that Ho substituted BFO has different bond angles (Fe-O2-Fe = 151.33°, Fe-O1-Fe = 147.91°) and bond lengths (Fe-O1 = 1.49 Å, Fe-O2 = 1.82 Å) compared to that of undoped BFO<sup>23</sup>. The alternation of these parameters is expected to have profound influence on the magnetic and optical properties of BHFO. Weight fraction of the phases, lattice parameters and Rietveld agreement factors listed in Table S1 and fitted plot of XRD pattern as Fig. S2 are provided as Supplementary Information.

			Lattice parameters	
Compound	Position $(2 \theta)_{(101)}$	Position $(2 \theta)_{(200)}$	a=b (Å)	c (Å)
TiO <sub>2</sub>	25.2681	48.0128	3.7867	9.5092
T <sup>95</sup> B <sup>5</sup>	25.2907	48.0459	3.7843	9.4964
T <sup>90</sup> B <sup>10</sup>	25.2452	48.0097	3.7869	9.5136
$T^{80}B^{20}$	25.2485	47.9995	3.7877	9.5136

Table 1. Position (20) and lattice parameters values obtained from XRD for pure film and composite films.

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Looking at the XRD patterns for composite thin films (see Fig. 2(b-d)) it can be immediately appreciated that intensity of the peaks of BHFO increases with an increase mol.% of BHFO nanoparticles. In particular, no impurity peak can be discerned in any of the composite films. This observation solidifies that the anatase phase of  $TiO_2$  is retained in the films and addition of nanoparticles does not introduce any phase change. Another feature of the composite thin films is to be observed- the diffraction peaks of different planes of composite films have shifted and the lattice parameters have changed with respect to pure one (see Table 1). The positions of diffraction peaks (101) and (200) were determined using Pseudo -Voight function in high score plus software and lattice parameters were calculated using the following formulae<sup>34</sup>

$$\frac{1}{d_{hkl}^2} = \left[ h^2 + k^2 + l^2 \left( \frac{a^2}{c^2} \right) \right] \frac{1}{a^2}$$
(4)

where  $d_{hkl}$  spacing has been calculated using the Bragg's law.

The peak shifting of composite films may be attributed to the diffusion of Fe<sup>3+</sup>, Ho<sup>3+</sup> or Bi<sup>3+</sup> ions into TiO<sub>2</sub> across the interfaces of BHFO and TiO<sub>2</sub>. Table 1 shows that diffraction peaks (101) and (200) have shifted to higher angle for T<sup>95</sup>B<sup>5</sup> and this film has the lowest lattice parameters. According to a previous study, this phenomenon manifests the introduction of tensile stresses due to the interstitial incorporation of doping ions in the crystal lattice of TiO<sub>2</sub><sup>35</sup>. Some other previous studies also report that Fe ions can go either in octahedral interstitial sites or in the substitutional positions of TiO<sub>2</sub> lattice<sup>36,37</sup>. Since the ionic radii of Fe<sup>3+</sup> (0.785 Å) is lower than that of Bi<sup>3+</sup> (1.14 Å) and Ho<sup>3+</sup> (1.015 Å)<sup>38</sup>, Fe<sup>3+</sup> ions have the highest chance of interstitial incorporation in TiO<sub>2</sub>. On the other hand, the diffraction peaks of T<sup>90</sup>B<sup>10</sup> and T<sup>80</sup>B<sup>20</sup> has shifted to the lower angle and they have larger lattice parameters compared to the pure film (see Table 1). These can be ascribed to the generation of compressive stresses as a result of substitutional incorporation of Bi<sup>3+</sup> (1.14 Å), Ho<sup>3+</sup> (1.015 Å) or Fe<sup>3+</sup> (0.785 Å)<sup>35,38</sup>. The XPS analysis also supports this phenomenon. The comparatively larger ionic radii of Bi, Ho and Fe compared to Ti<sup>4+</sup> (0.745 Å) ion<sup>38</sup> possibly cause an expansion of the crystal lattice and concordant shift in the TiO<sub>2</sub> diffraction peaks to a smaller angle. The substituted Ti<sup>4+</sup> ions may have also diffused in crystal structure of BHFO.

Figure 3 exhibits the surface images of nanoparticles and films annealed at 500 °C. The observed nanoparticles are nearly spherical in shape with average diameter of 45 nm as depicted in Fig. 3(a). It can be clearly seen from Fig. 3(a) that nanoparticles are agglomerated and thus they were dispersed in an ultrasonic bath before composite thin film preparation. On the other side, the grains of pure TiO<sub>2</sub> film are polygonal and crack-free (see Fig. 3(b)) signifying the fact that annealing and densification has occurred to a satisfactory degree. The average grain size of this film is larger than that of nanoparticles and is found to be ~60 nm. The average diameter of nanoparticles and grain size of thin films were determined using Image J software.

The addition of only 5 mol.% BHFO in  $TiO_2$  has introduced crack as well as diminished the uniformity and polygonality of the grains (see Fig. 3(c)). Further additions of 10 and 20 mol.% BHFO nanoparticles have totally changed the morphologies of thin films (see Fig. 3(d,e)).  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films have irregular polygonal grains and spherical nanoparticles which have made the surface rough. It has become difficult to identify  $TiO_2$  and BHFO nanoparticles separately from Fig. 3(d,e). Indeed,  $T^{80}B^{20}$  with 20 mol.% BHFO has larger number of porosities, cracks and rougher surface compared to  $T^{90}B^{10}$ . The FESEM surface morphology of  $T^{80}B^{20}$  film is consistent with the peak intensities ( $2\theta = 22.54^\circ$ ,  $32.14^\circ$ ,  $39.48^\circ$ ,  $46.185^\circ$  and  $57.34^\circ$ ) of XRD pattern (as depicted in Fig. 2) implying high amount of BHFO nanoparticles. It can be seen from the insets of Fig. 3(b–e), the average thicknesses of TiO<sub>2</sub>,  $T^{95}B^5$ ,  $T^{90}B^{10}$  and  $T^{80}B^{10}$  are 200, 244, 477, and 635 nm respectively. The cross sections of the films are also indicating the well-adherence of films with the glass substrate. Considering all these aforementioned observations it can be reasonably stated that there is no appreciable segregation of BHFO nanoparticles in the TiO<sub>2</sub> matrix.

Figure 4 depicts the absorbance spectra for thin films and nanoparticles. The absorbance of BHFO nanoparticles shown in the inset of Fig. 4 was obtained from diffused reflectance data using Kubelka-Munk conversion function<sup>39</sup>. The diffused reflectance data was converted to Kubelka-Munk function given by the following equation:

$$F(R) = \frac{(1-R)^2}{2R}$$
(5)

where R is the diffused reflectance value. On the other hand the absorbance of the films was determined using Beer- Lambert law as follows<sup>40,41</sup>:



**Figure 3.** FESEM images of (a) BHFO nanoparticles (b)  $TiO_2$  film (c)  $T^{95}B^5$  film (d)  $T^{90}B^{10}$  film and (e)  $T^{80}B^{20}$  film along with the films' cross-section micrographs. The magnification for all the images is x50000.

$$A = log \frac{(1-R)^2}{T}$$
(6)

where A, T and R stand for the absorbance, transmittance and reflectance of the film.

Nanoparticles exhibit extensive visible light absorbance and absorb larger wavelength than the synthesized films. The steep shape of the spectrum indicates that the visible light absorption has happened as a result of electronic transition from the valence band to conduction band  $(O^{2-} 2p \rightarrow Fe^{3+} 3d)$  in BHFO lattice<sup>23</sup>. This phenomenon will be discussed in detail in the following section. In contrast, the pure TiO<sub>2</sub> film does not start absorbing substantially until the incident wavelength falls below 370 nm (see Fig. 4). This is typical behavior of TiO<sub>2</sub> film<sup>19</sup>. It is noticeable from Fig. 4 that the absorption edge of composite thin films shows progressive shift between the pure TiO<sub>2</sub> thin film and BHFO nanoparticles. Even though the absorption edge of  $T^{95}B^5$  film does not shift much and still absorb shorter wavelength (~380 nm) than two other composite films Fig. 4 clearly depicts that  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films show a redshift and absorb visible light. The absorption edges of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films show a redshift and absorb visible light. The absorption edges of TiO<sub>2</sub> gets affected by diffusion of foreign ions which in turn influences the optical properties<sup>42,43</sup>. Indeed, the decrement of Ti<sup>3+</sup> concentration in T<sup>80</sup>B<sup>20</sup> film obtained from XPS analysis (see Table 1) suggests the possibility of A<sup>3+</sup> ions diffusion in the crystal structure of TiO<sub>2</sub>. Therefore, the bathochromic shift of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films could be due to the interdiffusion of A<sup>3+</sup>/Ti<sup>4+</sup> ions across the interfaces between BHFO and TiO<sub>2</sub>. In fact, due to interdiffusion of ions the electrons of these two films could be excited from the partially filled d orbitals of Fe<sup>3+</sup> (rather than the lower-lying O p orbitals) to the Ti 3d orbitals with lower energy visible light<sup>1</sup>.



Figure 4. UV-vis absorbance spectra of films. The inset shows the absorbance spectra of BHFO nanoparticles.



**Figure 5.** (a)  $[E^*F(R)]^2$  vs energy, E plot to calculate band gap energy of nanoparticles (b) indirect and (c) direct Tauc's plots demonstrating the band gap energies of pure and composite thin films.

In order to have better ideas on the optical properties of the films, their band gaps were determined as shown in Fig. 5. The band gap of BHFO nanoparticles obtained from  $[E *F(R)]^2$  vs E plot depicted in Fig.  $5(a)^{44}$  is found to be 1.88 eV, which is notably smaller than the previously reported values for both undoped<sup>44</sup> and many doped BFO<sup>45</sup>. The reduced band gap of BHFO could be attributed to the following reasons: (a) Ho<sup>3+</sup> ions are likely to have minimal degree of hybridization for a stable electronic configuration ( $4f^{10} 5d^0 6s^0$ ) which in turn may lead to the formation of unique energy level in between Fe 3d and O 2p and thereby effective band gap of Ho doped BFO is diminished, (b) according to some previous investigations, changes in Fe- O bond length and Fe-O-Fe bond angle by cation doping play a pivotal role in modifying one-electron bandwidth (W) and thus band gap of BFO<sup>23,46</sup>. This claim has been substantiated by the empirical formula relating W with bond length and angle:

$$V \approx \frac{cos\omega}{d_{Fe-O}^{3.5}}.$$
(7)

where  $\omega$  stands for  $\frac{1}{2}[\pi$ - (Fe-O-Fe] and  $d_{Fe-O}$  is the Fe-O bond length<sup>47</sup>. The relationship between band gap and W can be given as:  $E_q = \Delta - W$  where  $\Delta$  is the charge transfer energy<sup>47</sup>. Generally, BFO crystallizes in a rhombohedral phase and bond length (Fe-O) of this phase is greater than that of orthorhombic phase<sup>23</sup>. Since BHFO nanoparticles have orthorhombic phase and have smaller bond length, they have appreciably larger bandwidth (W) value than that of BFO which could reduce the effective band gap energy of BHFO.



**Figure 6.** Schematic illustration for the calculated energy level diagram indicating the conduction and valence band potentials of TiO<sub>2</sub> film and BHFO nanoparticles.

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The optical band gap of the films were calculated using the Tauc's law<sup>48</sup> as given below:

$$(\alpha h v) = A' \left( E - E_o \right)^n \tag{8}$$

where A' is a constant, E is the photon energy,  $E_g$  is the band gap energy,  $\alpha$  is the absorption coefficient and n is the power factor of the optical transition mode in a semiconductor, i.e., direct transition (n = 1/2) or indirect transition (n = 2). The band gaps of the films were obtained by extrapolating the linear portion of of  $(\alpha h\nu)^2$  vs E plot for direct transition (see Fig. 5(c)) and  $(\alpha h\nu)^{1/2}$  vs E plot for indirect transition (see Fig. 5(b)). Since the indirect band gaps of the films are more related to their absorption edge values (see Fig. 4) they have been assigned to the respective composite thin films. Even both experimental results and theoretical modeling suggest that TiO<sub>2</sub> has a direct forbidden gap, which is almost degenerated with an indirect allowed transition. Therefore, the indirect allowed transition dominates in the optical absorption just above the absorption edge due to the weak strength of the direct forbidden transition<sup>49</sup>.

The Tauc's plots of composite thin films corresponding to indirect transition shown in Fig. 5 have only one linear region. However, the curve for pure  $\text{TiO}_2$  has more than one linear regions which indicates that there is more than one optical transition occurring in this film. The first transition of pure film occurred at 3.4 eV is its effective band gap which agrees reasonably well with previous studies<sup>35,50</sup>. The second indirect transition occurs at 2.1 eV and it is believed that this sort of transition is attributed to the presence of OVs in deposited TiO<sub>2</sub> film<sup>19</sup>.

The indirect band gaps of  $T^{95}B^5$ ,  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are 3.13, 2.87 and 2.46 respectively (see Fig. 5). Thus, the optical transition energies of composites thin films monotonically decrease with increasing mol.% of BHFO nanoparticles. A perceptive explanation for the reduced band gap is related to conduction band of two different semiconductors. The hypothesis is that if the conduction band edge of the sensitizing material is higher than the conduction band edge of TiO<sub>2</sub>, electrons can transfer from the smaller band gap material to the conduction band of TiO<sub>2</sub><sup>18,19</sup>. According to some previous reports BFO has relatively high conduction band edge<sup>1,51</sup> which makes electrons in BFO nanoparticles to be transferred easily into the lower lying CB of TiO<sub>2</sub> through the interface. Thus it is required to determine the conduction band (CB) and valence band (VB) positions of both BHFO and TiO<sub>2</sub> to elucidate optical transition phenomenon. The CB potential of BHFO and TiO<sub>2</sub> at the point of zero charge can be calculated successfully by the following empirical equation<sup>52,53</sup>.

$$E_{CB} = \chi_{\rm S} - E^e - 0.5E_{\rm g} \tag{9}$$

where  $E_{CB}$  is the CB edge potential,  $E^e$  is the energy of free electrons on the hydrogen scale (~4.5 eV),  $E_g$  is the band gap energy of the semiconductor,  $\chi_s$  is the electronegativity of the semiconductor, which is the geometric mean of the electronegativity of the constituent atoms<sup>53</sup>. The detail information on  $\chi_s$  ihas been provided in the Supplementary Information. Plugging  $E^e$ ,  $\chi_s$  values of TiO<sub>2</sub> (5.81 eV) and BHFO (5.86 eV), and their corresponding  $E_g$  values into the above equations,  $E_{CB}$  of TiO<sub>2</sub> and BHFO stands as -0.39 and 0.42 eV respectively. The valence band edge potential  $E_{VB}$  can be obtained by the equation:  $E_{VB} = E_{CB} + E_g$  and the calculated  $E_{VB}$  of TiO<sub>2</sub> and BHFO are 3.01 and 2.3 eV respectively. Figure 6 depicts the calculated energy levels of TiO<sub>2</sub> and BHFO.

It can be clearly seen from Fig. 6 that BHFO has relatively higher CB potential than that of  $TiO_2$ . As discussed earlier, the redshift shown by  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films is probably induced by Fe, Bi, Ho/Ti interdiffusion in the interfaces which could raise the CB edge potentials of BHFO and  $TiO_2$  and hence lower the optical transition energy at the interface between BHFO and  $TiO_2$ . Such reduced band-gaps of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  could be beneficial for the efficient utilization of visible light for photocatalysis. Moreover from XPS analysis a higher amount of OH<sup>-</sup> has been found in  $T^{80}B^{20}$  which allows the generation of large amount of highly reactive hydroxyl radicals (OH<sup>•</sup>) through oxidation of OH<sup>-</sup> by photo induced holes rendering the composite film highly reactive during photocatalysis.



**Figure 7.** (a) Resistivity, carrier concentration and mobility of films with different mol.% of BHFO (b) Conductivity ( $\Omega$ -cm) vs. Electric Field, E (kV/cm) plot of BHFO nanoparticles.

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From the Hall effect analysis, the Hall coefficient is found to be negative for all the thin films which indicates that they are n- type semiconductors. Generally,  $TiO_2$  is an intrinsic n-type semiconductor<sup>8</sup>. The formation of the OVs and titanium interstitials are responsible for this sort of conductivity<sup>54</sup>. TiO<sub>2</sub> is a sub-stoichiometric compound with excess titanium under standard conditions. This sub-stoichiometry is accommodated as OVs or titanium interstitials formation. The following reaction mechanisms can be considered for the formation of OVs in TiO<sub>2</sub> crystal<sup>55</sup>:

$$O_{O}^{x} \xrightarrow{TiO_{2}} \frac{1}{2}O_{2}(g) + V_{O}^{\bullet \bullet} + 2e^{'}$$

$$\tag{10}$$

$$2Ti_{Ti}^{x} + 2e' = 2Ti_{Ti}^{'} \tag{11}$$

Combining Equations (4 and 5)

$$2Ti_{Ti}^{x} + O_{O}^{x} = 2Ti_{Ti}^{'} + \frac{1}{2}O_{2}(g) + V_{O}^{\bullet\bullet}$$
(12)

As follows the n-type conductivity behavior found in TiO<sub>2</sub> implies Equation (6) to be the dominant step for providing two excess electrons. The carrier concentration (CC), resistivity and mobility of pure TiO<sub>2</sub> thin film deposited using non aqueous sol-gel method are found to be  $2.0 \times 10^{16} \text{ cm}^{-3}$ , 1550.8  $\Omega$ -cm and 208 cm<sup>2</sup>/V.s respectively. The CC of TiO<sub>2</sub> found in the literatures ranges from  $1 \times 10^{16} \text{ cm}^{-3}$  to  $1 \times 10^{20} \text{ cm}^{-356,57}$ . Thus the measured value of CC is within accepted range. The measured mobility of the film is also consistent with the previous study which reported almost similar CC<sup>8</sup>. The resistivity of conductor and semiconductor normally varies from  $10^{-3}$  to  $10^8 \Omega$ -cm<sup>58</sup>. Based on this range, pure TiO<sub>2</sub> film can be classified as semiconductor. However, the resistivity of this pure TiO<sub>2</sub> film is lower than prepared by many methods such as chemical bath deposition<sup>58</sup>, spray pyrolysis<sup>59</sup> and DC magnetron sputtering technique<sup>26</sup>.

Figure 7(a) depicts that 5 mol.% BHFO addition has significantly altered the CC, mobility and resistivity of the film. The resistivity of  $T^{95}B^5$  film is ~2.33 orders of magnitude higher than pure film while the CC and mobility of this film are 62.5 and 2.6 times lower than those of pure one. Moreover, the charge carriers of  $T^{95}B^5$  film are dominated by free electrons at 5 mol.% BHFO addition, even though, BFO is reported to be a p-type semiconductor where holes are the majority carriers. Here, substitution of Bi<sup>3+</sup> with Ho<sup>3+</sup> should retain BHFO as p-type semiconductor<sup>60,61</sup>. The determination of CC of BHFO was out of scope for the current study. Moreover, to the best of our knowledge there is no report on CC of BFO or on doped BFO<sup>60,61</sup>. Since TiO<sub>2</sub> is the major phase its carrier will ultimately control the carrier concentration. The concentration of electrons may also be affected subtly by the diffused Fe ions. Since Fe ions have more than one valence states such as Fe<sup>2+</sup> or Fe<sup>3+62</sup>, some of diffused Fe<sup>3+</sup> may have changed into Fe<sup>2+</sup> by accepting free electrons of TiO<sub>2</sub> matrix and reduces total CC of T<sup>95</sup>B<sup>5</sup> slightly. Thus the total CC of the composite film decreases with the addition of BHFO.

The additions of 10 and 20 mol.% BHFO have not changed the electrical properties of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films (see Fig. 7(a)) appreciably. The mobilities of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films have decreased by a factor of 1.23 and 1.6 respectively and the CCs of these two films have been reduced by a factor of ~4 .1 when compared to those of  $T^{95}B^5$  film. Similarly, the resistivities of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are 1.05 and 1.15 times higher than that of  $T^{95}B^5$  film. However, the charge carriers of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are still found to be dominated by free electrons. The aforestated variations in CC, mobility and resistivity point out that addition of higher mol.% of BHFO has little effect on electrical properties of composite films. The trivial reduction of CC in  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are  $T^{4+}$  from the lattice sites of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are still found to be dominated by a factor of BHFO has little effect on electrical properties of composite films. The trivial reduction of CC in  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are  $T^{4+}$  from the lattice sites of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are trivial reduction of CC in  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are substitute  $T^{4+}$  from the lattice sites of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are trivial reduction of CC in  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are substitute  $T^{4+}$  from the lattice sites of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films are trivial reduction of  $T^{90}B^{10}$  and



**Figure 8.** Room temperature (**a**) M-H hysteresis loop for BHFO nanoparticles with inset showing its magnified M-H loop (**b**) M-H hysteresis loops for pure and composite films with inset showing magnified M-H loops of  $T^{95}B^5$ ,  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films.

eventually reduces the films' electron/ $Ti^{3+}$  concentration. However, further reduction in CC (electron hoping site concentration) is probably hindered by intergranular cracks and pore formations in the films (see Fig. 3(d,e)) as well as the direct contact interface areas between BHFO and  $TiO_2$  are likely to become gradually saturated at higher BHFO nanoparticles content (10 and 20 mol.%) in the composite films. These phenomena could hamper further enhancement in migration of Ho<sup>3+</sup>, Fe<sup>3+</sup> from BHFO to  $TiO_2$  in  $T^{90}B^{10}$  and  $T^{80}B^{20}$  restricting any marked alteration in overall carrier concentration of the system.

The increased resistivity of composite thin films is directly related to low conductivity of BHFO (see Fig. 7(b)). The conductivity of BHFO varies from  $1 \times 10^{-12}$  to  $5 \times 10^{-12}$  ( $\Omega$ . cm)<sup>-1</sup> which is lower than that of previously reported doped BFO<sup>23</sup>. The CC and mobility are also related to resistivity via the following equation<sup>63</sup>:

ρ

$$\rho = \frac{1}{qn\mu} \tag{13}$$

where q is the electron charge, n is carrier concentration and  $\mu$  is the mobility. The reduced CC and mobility of composite films could be ascribed to increased resistivity of films. Moreover, the poor morphology of the films may reduce mobility of films. Yasuno *et al.*<sup>64</sup> and Trinh *et al.*<sup>65</sup> observed that electrical properties of the film depend on surface roughness and porosity of film. These surface roughness and porosity increase the scattering of electrons and holes which in turn reduce their mean free path. From FESEM images it is clear that the pure film has a smoother surface than the composite films. With increasing mol.% of BHFO, the surface roughness and porosity of the films decrease and charge carrier mean free path of the films decreases. Thus, the mobilities of the films decrease with reduced mean free path of electrons. The composite films with enhanced resistivity, and reduced CC and mobility will have minimum leakage current and reduced power losses in devices. Therefore, these films could be electrically reliable for resistors, sensors and memory devices.

Generally BFO nanoparticles exhibit G-type antiferromagnetism due to its cycloid spiral spin structure<sup>23</sup>. However, an enhancement in magnetization by Ho-doping at Bi-site is evidenced in this study (see Fig. 8(a)). The saturation magnetization (Ms), coercivity (Hc) and magnetic susceptibility ( $\chi_M$ ) of BHFO nanoparticles are 24.33 emu/cm<sup>3</sup>, 260 Oe and 0.247 respectively. The density of BHFO is 8.40 g/cm<sup>3</sup> which was calculated from unit cell parameters listed in Table S1. The magnetization unit of BHFO was converted from emu/gm to emu/cm<sup>3</sup> by multiplying with 8.40 g/cm<sup>3</sup>. The bond lengths and bond angles of BHFO listed in Table S1 are quite different from the reported bond lengths and bond angles of pure BFO<sup>23</sup>. The enhanced magnetization could be attributed to these changes in bond angle and bond length which modify the tilting angle of FeO<sub>6</sub> octahedron and thereby suppress spiral modulated spin stricture<sup>23</sup>. Moreover, the substitution of non-magnetic Bi<sup>3+</sup> ([Xe] 4f<sup>14</sup> 5d<sup>10</sup> 6 s<sup>2</sup> 6p<sup>0</sup>) with magnetic Ho<sup>3+</sup> ([Xe] 4f<sup>10</sup>) having high magnetic moment (10.6  $\mu_B$ ) could lead to improved magnetization.

Room temperature M-H curve of thin films were obtained after subtracting the contribution of glass substrate. TiO<sub>2</sub> film deposited on this substrate is diamagnetic (see Fig. 8(b)) and the diamagnetic susceptibility  $(\chi_D')$  is  $\sim -1.58 \times 10^{-03}$ . Donor type defects (e.g. OVs) have been proposed to be a vital element for ferromagnetism in TiO<sub>2</sub><sup>66,67</sup> indicating that synthesized TiO<sub>2</sub> film may not have sufficient amount of OVs to induce magnetism. The other three composite thin films show no such diamagnetic behavior yet they display no complete ferromagnetic loop either. The composite thin films have ferromagnetic behavior accompanied with a diamagnetic component. The discernibly open hysteretic loop shown in the inset of Fig. 8(b) is signifying the presence of a ferromagnetic fields and all the films have diamagnetic components at higher magnetic fields. Figure 8(b) clearly depicts that saturation magnetization of the films increases with increasing mol.% of BHFO. T<sup>80</sup>B<sup>20</sup> has higher magnetization and coercivity than two other composite films (see inset of Fig. 8(b)). Moreover, it has retained its ferromagnetism upto 11 kOe and has the lowest diamagnetic susceptibility (3.28 × 10<sup>-4</sup>). The diamagnetic susceptibilities of other two films are  $1.01 \times 10^{-03}$  (T<sup>95</sup>B<sup>5</sup>) and  $6.045 \times 10^{-4}$  (T<sup>90</sup>B<sup>10</sup>). Previous investigations reveal that diffusion of magnetic ions into an oxide ceramic through interfaces aids in emergence of room temperature ferromagnetism<sup>68,69</sup>. As conferred earlier, the interdiffusion of Fe<sup>3+</sup>, Ho <sup>3+</sup> and Ti<sup>4+</sup> in the interfaces brings about a decrement of Ti<sup>3+</sup> concentration in the crystal structure of TiO<sub>2</sub>. The ferromagnetic exchange between the diffused magnetic ions (Fe<sup>3+</sup> or Ho<sup>3+</sup>) and OVs could induce magnetism<sup>68,70</sup>. The lacking of sufficient diffused magnetic ions to couple with OVs possibly leads to diamagnetism at higher magnetic field for the films. Since T<sup>95</sup>B<sup>5</sup> and T<sup>90</sup>B<sup>10</sup> have lower mol.% of BHFO compared to T<sup>80</sup>B<sup>20</sup>, they are expected to have smaller amount of diffused Fe<sup>3+</sup> or  $Ho^{3+}$  ions in  $TiO_2$  phase than that of  $T^{80}B^{20}$ . Therefore, the ferromagnetic exchange of  $T^{95}B^5$  and  $T^{90}B^{10}$  is weak due to small amount of diffused magnetic ions and thus diamagnetic susceptibility of these films is higher than that of T<sup>80</sup>B<sup>20</sup>. On the other hand, T<sup>80</sup>B<sup>20</sup> film has higher amount of magnetic ions when compared with other two composite films. Hence, a stronger ferromagnetic exchange between magnetic ions and OVs is expected in T<sup>80</sup>B<sup>20</sup> film leading to a lower diamagnetic susceptibility compared to the other two composite films. Indeed, composites thin film with further addition of BHFO nanoparticles (>20 mol.%) will have diamagnetic component. It can be seen from Fig. S5 that the reduction in diamagnetic susceptibility from  $TiO_2$  film to  $T^{95}B^5$  film is prominent but further addition of BHFO nanoparticles does not bring such kind of change. In fact Fig. S5 depicts that diamagnetic susceptibility of T<sup>80</sup>B<sup>20</sup> film is approaching saturation phenomenon conforming our explanation provided in electrical section. The presence of pores and cracks in films (see Fig. 3(d,e)) as well as possible saturation of total amount of direct contact interface areas between BHFO and TiO<sub>2</sub> not only affect the electrical properties but also the magnetic properties of  $T^{90}B^{10}$  and  $T^{80}B^{20}$  films. Nevertheless, this finding suggests that incorporation of BHFO nanoparticles into TiO<sub>2</sub> matrix introduces magnetism in the final product which can be applied to thin films-based spintronics application.

### Conclusion

In conclusion, novel BHFO/TiO<sub>2</sub> composite thin films were prepared for the first time by non-aqueous sol-gel method. The incorporation of BHFO nanoparticles improves magnetism and electrical properties of composite thin films, and also reduces their band gaps. The composite thin films exhibit a redshift in optical absorption edge and absorb visible light. Such behavior has been attributed to interdiffusion of  $A^{3+}/Ti^{4+}$  ions through the interfaces between BHFO and TiO<sub>2</sub>. Besides, interfacial microstructural defects and the direct contact areas between BHFO and TiO<sub>2</sub> in composite films control the migration of these ions through their interfaces which eventually determines the amount of CC in the films. Moreover, the ferromagnetic exchange between the diffused magnetic ions (Fe<sup>3+</sup> or Ho<sup>3+</sup>) and OVs has been considered to influence the induction of magnetism in composite thin films.

#### References

- 1. Li, S., Lin, Y.-H., Zhang, B.-P., Li, J.-F. & Nan, C.-W. BiFeO3/TiO2 core-shell structured nanocomposites as visible-active photocatalysts and their optical response mechanism. *Journal of Applied Physics* **105**, 054310 (2009).
- Wang, X., Li, J.-G., Kamiyama, H., Moriyoshi, Y. & Ishigaki, T. Wavelength-sensitive photocatalytic degradation of methyl orange in aqueous suspension over iron (III)-doped TiO2 nanopowders under UV and visible light irradiation. *The Journal of Physical Chemistry B* 110, 6804–6809 (2006).
- Nkosi, S. et al. An instant photo-excited electrons relaxation on the photo-degradation properties of TiO2-x films. Journal of Photochemistry and Photobiology A: Chemistry 293, 72–80 (2014).
- 4. Serpone, N., Lawless, D., Disdier, J. & Herrmann, J.-M. Spectroscopic, photoconductivity, and photocatalytic studies of TiO2 colloids: naked and with the lattice doped with Cr3+, Fe3+, and V5+ cations. *Langmuir* **10**, 643–652 (1994).
- Weissmann, M. & Errico, L. A. The role of vacancies, impurities and crystal structure in the magnetic properties of TiO2. *Physica B: Condensed Matter* 398, 179–183 (2007).
- 6. Gale, E. TiO2-based memristors and ReRAM: materials, mechanisms and models (a review). *Semiconductor Science and Technology* 29, 104004 (2014).
- 7. Wang, Z. et al. Extraordinary Hall effect and ferromagnetism in Fe-doped reduced rutile. Applied Physics Letters 83, 518-520 (2003).
- Li, X. et al. Structures and magnetic properties of p-type Mn: TiO2 dilute magnetic semiconductor thin films. Journal of Applied Physics 106, 043913 (2009).
- Chen, X. & Burda, C. The electronic origin of the visible-light absorption properties of C-, N-and S-doped TiO2 nanomaterials. Journal of the American Chemical Society 130, 5018–5019 (2008).
- Kim, Y. J. et al. Growth and properties of molecular beam epitaxially grown ferromagnetic Fe-doped TiO2 rutile films on TiO2 (110). Applied physics letters 84, 3531–3533 (2004).
- 11. Kim, J.-Y. et al. Ferromagnetism Induced by Clustered Co in Co-Doped Anatase TiO2 Thin Films. Physical review letters **90**, 017401 (2003).
- 12. Gao, X., Liu, X., Zhu, Z., Wang, X. & Xie, Z. Enhanced photoelectrochemical and photocatalytic behaviors of MFe2O4 (M = Ni, Co, Zn and Sr) modified TiO2 nanorod arrays. *Scientific reports* 6, 30543 (2016).
- 13. Van Suchtelen, J. Product properties: a new application of composite materials. Philips Res. Rep 27, 28-37 (1972).
- 14. Watson, S., Beydoun, D. & Amal, R. Synthesis of a novel magnetic photocatalyst by direct deposition of nanosized TiO2 crystals onto
- a magnetic core. *Journal of Photochemistry and Photobiology A: Chemistry* 148, 303–313 (2002).
  15. Huang, W., Tang, X., Felner, I., Koltypin, Y. & Gedanken, A. Preparation and characterization of FexOy–TiO2 via sonochemical synthesis. *Materials research bulletin* 37, 1721–1735 (2002).
- 16. Cheng, C. et al. Enhanced photocatalytic performance of TiO2-ZnO hybrid nanostructures. Scientific reports 4, 4181 (2014).
- Mu, J. et al. Enhancement of the visible-light photocatalytic activity of In2O3–TiO2 nanofiber heteroarchitectures. ACS applied materials & interfaces 4, 424–430 (2011).
- Yang, Y. et al. Electrospun nanofibers of p-type BiFeO3/n-type TiO2 hetero-junctions with enhanced visible-light photocatalytic activity. RSC Advances 4, 31941–31947 (2014).
- Larsen, G. K., Fitzmorris, B. C., Longo, C., Zhang, J. Z. & Zhao, Y. Nanostructured homogenous CdSe–TiO2 composite visible light photoanodes fabricated by oblique angle codeposition. *Journal of Materials Chemistry* 22, 14205–14218 (2012).
- Zhang, Y., Schultz, A. M., Salvador, P. A. & Rohrer, G. S. Spatially selective visible light photocatalytic activity of TiO2/BiFeO3 heterostructures. *Journal of Materials Chemistry* 21, 4168–4174 (2011).
- Zhu, A., Zhao, Q., Li, X. & Shi, Y. BiFeO3/TiO2 nanotube arrays composite electrode: construction, characterization, and enhanced photoelectrochemical properties. ACS applied materials & interfaces 6, 671–679 (2013).

- Sarkar, A., Singh, A. K., Sarkar, D., Khan, G. G. & Mandal, K. Three-dimensional nanoarchitecture of BiFeO3 anchored TiO2 nanotube arrays for electrochemical energy storage and solar energy conversion. ACS Sustainable Chemistry & Engineering 3, 2254–2263 (2015).
- Hasan, M. et al. Saturation magnetization and band gap tuning in BiFeO3 nanoparticles via co-substitution of Gd and Mn. Journal of Alloys and Compounds 687, 701–706 (2016).
- 24. Sakar, M., Balakumar, S., Saravanan, P. & Jaisankar, S. Annealing temperature mediated physical properties of bismuth ferrite (BiFeO3) nanostructures synthesized by a novel wet chemical method. *Materials Research Bulletin* **48**, 2878–2885 (2013).
- Song, G. et al. Crystal structure refinement, ferroelectric and ferromagnetic properties of Ho3+ modified BiFeO3 multiferroic material. Journal of Alloys and Compounds 696, 503–509 (2017).
- Yildiz, A., Lisesivdin, S., Kasap, M. & Mardare, D. Electrical properties of TiO2 thin films. *Journal of Non-Crystalline Solids* 354, 4944–4947 (2008).
- Hong, N. H., Sakai, J., Poirot, N. & Brizé, V. Room-temperature ferromagnetism observed in undoped semiconducting and insulating oxide thin films. *Physical Review B* 73, 132404 (2006).
- Legrand-Buscema, C., Malibert, C. & Bach, S. Elaboration and characterization of thin films of TiO2 prepared by sol-gel process. *Thin solid films* 418, 79–84 (2002).
- Rodriguez-Carvajal, J. FULLPROF: a program for Rietveld refinement and pattern matching analysis. Satellite meeting on powder diffraction of the XV congress of the IUCr. 127 (1990).
- 30. Liu, H. *et al.* Synthesis and characterization of titania prepared by using a photoassisted sol- gel method. *Langmuir* **19**, 3001–3005 (2003).
- 31. Sanjines, R. et al. Electronic structure of anatase TiO2 oxide. Journal of Applied Physics 75, 2945–2951 (1994).
- 32. Hellenbrandt, M. The inorganic crystal structure database (ICSD) Present and Future. Crystallography Reviews 10, 17–22 (2004).
- 33. Qi, Y., Alima, B. & Shifeng, Z. Lattice distortion of holmium doped bismuth ferrite nanofilms. *Journal of Rare Earths* **32**, 884–889 (2014).
- 34. Ladd, M. F. C., Palmer, R. A. & Palmer, R. A. Structure determination by X-ray crystallography. (Springer, 1985).
- Munir, S., Shah, S. M. & Hussain, H. Effect of carrier concentration on the optical band gap of TiO2 nanoparticles. *Materials & Design* 92, 64–72 (2016).
- Shah, S., Li, W., Huang, C.-P., Jung, O. & Ni, C. Study of Nd3+, Pd2+, Pt4+, and Fe3+ dopant effect on photoreactivity of TiO2 nanoparticles. *Proceedings of the National Academy of Sciences* 99, 6482–6486 (2002).
- Nguyen, V. N., Nguyen, N. K. T. & Nguyen, P. H. Hydrothermal synthesis of Fe-doped TiO2 nanostructure photocatalyst. Advances in Natural Sciences: Nanoscience and Nanotechnology 2, 035014 (2011).
- Shannon, R. T. & Prewitt, C. T. Effective ionic radii in oxides and fluorides. Acta Crystallographica Section B: Structural Crystallography and Crystal. Chemistry 25, 925–946 (1969).
- 39. Von Kubelka, P. Ein Beitrag zur Optik der Farbanstriche. Zeitschrift fur technische Physik, 593-601 (1931).
- 40. Beer, A. Bestimmung der absorption des rothen lichts in farbigen flussigkeiten. Ann. Physik 162, 78-88 (1852).
- 41. Lambert, J. H. Photometria sive de mensura et gradibus luminis, colorum et umbrae. (Klett, 1760).
- Nagaveni, K., Hegde, M. & Madras, G. Structure and Photocatalytic Activity of Ti1-xMxO2± δ (M = W, V, Ce, Zr, Fe, and Cu) Synthesized by Solution Combustion Method. *The Journal of Physical Chemistry B* 108, 20204–20212 (2004).
- Moriguchi, I., Maeda, H., Teraoka, Y. & Kagawa, S. Preparation of TiO2 ultrathin film by newly developed two-dimensional sol-gel process. *Journal of the American Chemical Society* 117, 1139–1140 (1995).
- Mocherla, P. S., Karthik, C., Ubic, R., Ramachandra Rao, M. & Sudakar, C. Tunable bandgap in BiFeO3 nanoparticles: the role of microstrain and oxygen defects. *Applied Physics Letters* 103, 022910 (2013).
- Zhou, W. et al. Effects of Sm and Mn co-doping on structural, optical and magnetic properties of BiFeO3 films prepared by a sol-gel technique. Materials Letters 144, 93–96 (2015).
- 46. Zhang, Z., Wu, P., Chen, L. & Wang, J. Systematic variations in structural and electronic properties of BiFeO3 by A-site substitution. *Applied Physics Letters* **96**, 012905 (2010).
- Mcdarde, M. et al. High-pressure neutron-diffraction study of the metallization process in PrNiO3. Physical Review B 52, 9248 (1995).
- Tauc, J., Grigorovici, R. & Vancu, A. Optical properties and electronic structure of amorphous germanium. *physica status solidi (b)* 15, 627–637 (1966).
- Tang, H., Prasad, K., Sanjines, R., Schmid, P. & Levy, F. Electrical and optical properties of TiO2 anatase thin films. *Journal of applied physics* 75, 2042–2047 (1994).
- Ahamed, M., Khan, M. M., Akhtar, M. J., Alhadlaq, H. A. & Alshamsan, A. Ag-doping regulates the cytotoxicity of TiO2 nanoparticles via oxidative stress in human cancer cells. *Scientific reports* 7, 17662 (2017).
- Luo, J. & Maggard, P. A. Hydrothermal Synthesis and Photocatalytic Activities of SrTiO3-Coated Fe2O3 and BiFeO3. Advanced Materials 18, 514–517 (2006).
- Kim, Y. I., Atherton, S. J., Brigham, E. S. & Mallouk, T. E. Sensitized layered metal oxide semiconductor particles for photochemical hydrogen evolution from nonsacrificial electron donors. *The Journal of Physical Chemistry* 97, 11802–11810 (1993).
- Butler, M. & Ginley, D. Prediction of flatband potentials at semiconductor-electrolyte interfaces from atomic electronegativities. Journal of the Electrochemical Society 125, 228–232 (1978).
- Morgan, B. J. & Watson, G. W. Intrinsic n-type defect formation in TiO2: a comparison of rutile and anatase from GGA+ U calculations. *The Journal of Physical Chemistry C* 114, 2321–2328 (2010).
- 55. Walsh, A. & Catlow, C. R. A. Photostimulated reduction processes in a titania hybrid metal–organic framework. *Chem Phys Chem* 11, 2341–2344 (2010).
- Lee, E.-J. & Pyun, S.-I. Analysis of nonlinear Mott-Schottky plots obtained from anodically passivating amorphous and polycrystalline TiO2 films. *Journal of applied electrochemistry* 22, 156–160 (1992).
- Enache, C. S., Schoonman, J. & Van Krol, R. The photoresponse of iron-and carbon-doped TiO2 (anatase) photoelectrodes. *Journal of electroceramics* 13, 177–182 (2004).
- Manurung, P., Putri, Y., Simanjuntak, W. & Low, I. M. Synthesis and characterisation of chemical bath deposited TiO2 thin-films. *Ceramics International* 39, 255–259 (2013).
- 59. Rashti, M. & Brodie, D. The photoresponse of high resistance anatase TiO2 films prepared by the decomposition of titanium isopropoxide. *Thin Solid Films* **240**, 163–167 (1994).
- Perejón, A. et al. Preparation of phase pure, dense fine grained ceramics by conventional and spark plasma sintering of Lasubstituted BiFeO3 nanoparticles. Journal of the European Ceramic Society 35, 2283–2293 (2015).
- Perejón, A. et al. Single phase, electrically insulating, multiferroic La-substituted BiFeO 3 prepared by mechanosynthesis. Journal of Materials Chemistry C 2, 8398–8411 (2014).
- 62. Song, G. *et al.* Effect of Ho3+ doping on the electric, dielectric, ferromagnetic properties and TC of BiFeO3 ceramics. *Ceramics International* **40**, 3579–3587 (2014).
- 63. Neaman, D. A. Semiconductor physics and devices. (Irwin Chicago, 1992).
- 64. Yasuno, S. et al. Physical properties of amorphous In-Ga-Zn-O films deposited at different sputtering pressures. Japanese Journal of Applied Physics 52, 03BA01 (2013).

- Trinh, T. T., Jang, K., Velumani, S., Dao, V. A. & Yi, J. Role of Schottky barrier height at source/drain contact for electrical improvement in high carrier concentration amorphous InGaZnO thin film transistors. *Materials Science in Semiconductor Processing* 38, 50–56 (2015).
- Coey, J., Venkatesan, M. & Fitzgerald, C. Donor impurity band exchange in dilute ferromagnetic oxides. Nature materials 4, 173–179 (2005).
- 67. Zhao, T. et al. Electric field effect in diluted magnetic insulator anatase Co: TiO2. Physical review letters 94, 126601 (2005).
- Zhang, L. *et al.* Interfacial effect on Mn-doped TiO2 nanoparticles: from paramagnetism to ferromagnetism. *RSC Advances* 6, 57403–57408 (2016).
- 69. García, M. *et al.* Interface double-exchange ferromagnetism in the Mn-Zn-O system: new class of biphase magnetism. *Physical review letters* **94**, 217206 (2005).
- Choudhury, B. & Choudhury, A. Room temperature ferromagnetism in defective TiO2 nanoparticles: Role of surface and grain boundary oxygen vacancies. *Journal of Applied Physics* 114, 203906 (2013).

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#### Author Contributions

M.R.I. conducted the synthesis process and wrote whole manuscript, supervised by A.K.M.B.R. and M.A.Z. The Hall effect measurement was done by M.S.B.

### Additional Information

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