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Plasmonic Pt nanoparticles—TiO₂ hierarchical nano-architecture as a visible light photocatalyst for water splitting

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Visible light-driven water splitting (VLWS) into hydrogen and oxygen is attractive and depends on efficient photocatalysts. Herein, we demonstrate the first exploration of the capability to control the morphology of nanostructured TiO₂ in conjunction with the choice of a suitable plasmonic metal (PM) to fabricate novel photocatalysts that are capable of harvesting visible light for more efficient VL-fuel conversion. This methodology affords us to successful access to the novel plasmonic Pt/TiO₂-HA (large Pt nanoparticles (NPs) supported on TiO₂ hierarchical nano-architecture (TiO₂-HA)) photocatalysts that exhibit plasmon absorption in the visible range and consequent outstanding activity and durability for VLWS. Particularly, the Pt/TiO₂-HA shows an excellent photocatalytic activity for overall water splitting rather than only for hydrogen evolution (HE), which is superior to those of the conventional plasmonic Au/TiO₂ photocatalysts. The synergistic effects of the high Schottky barrier at the Pt-TiO₂-HA interface, which induces the stronger reduction ability of hot electrons, and intrinsic Pt catalytic activity are responsible for the exceptional photocatalytic performance of Pt/TiO₂-HA and simplify the composition of plasmonic photocatalysts.

Design and fabrication of high-performance photocatalysts with suitable architectures for VLWS into hydrogen and oxygen has been a topic of tremendous scientific interest in recent years because this process is promising for resolving today's increased global environmental crisis and energy shortage¹⁻⁴. The exploration of highly efficient photocatalysts with a strong response to visible light, a high activity and stability, and low cost is critical to construct an artificial leaf for realizing solar energy conversion. Among various promising photocatalysts for water splitting, TiO₂ is an ideal candidate and has received much more attention than all other known efficient ones owing to its unique and diverse favorable properties⁵, such as a facile manipulation of its nanoscale morphology, excellent long-term chemical stability, robust anti-photo-corrosion ability endowed by strong chemical bonding between Ti(IV) and O ions, low toxicity, and easy availability due to its Earth-abundance. However, the conversion efficiency of TiO₂ still remains at a very low level because of its limited absorption of solar radiation ($\lambda < 400$ nm) and the intractable recombination of photogenerated charge carriers in it.

Notably, a variety of methods, including doping with various elements⁶⁻⁸, engineering nano-heterostructures with semiconductor quantum dots or metal oxides⁹⁻¹², and combining with PM-NPs (typically, Au and Ag)¹³⁻³¹, have been reported to effectively enhance visible light harvesting of TiO₂ and concomitantly promote carrier separation and transfer. In particular, many PM-NPs/TiO₂ composites have shown remarkably plasmon-enhanced photocatalytic activities towards VLWS by extending light absorption to visible region, which is induced by greatly enhancing local electric field at the interface of PM-NPs and TiO₂ or by injecting hot electrons from photo-excited PM-NPs to TiO₂ via surface plasmon resonance (SPR) effects of the supported PM-NPs¹³⁻³⁰. Therefore, advanced strategies for producing nano-architectures consisting of PM-NPs and TiO₂ have been actively pursued by world-wide researchers because of the extraordinarily high stability of Au and Ag and the fascinating underlying physical mechanisms.

However, Au and Ag NPs show very low, and even not, chemical catalytic activities for water splitting. Concomitantly, the coverage of these PM-NPs could significantly diminish the catalytically active surface area²⁸.

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These drawbacks generally leads to the Au or Ag NPs-decorated semiconductor photocatalysts showing activity towards a half reaction of water splitting producing only H₂ or O₂ with the need of sacrificial reagents (SR)^{18–25} or additional driving bias^{26–28,32,33}. To date, there have been very few reports on photocatalysts used for the full VLWS (i.e., SR-free and unbiased)^{29,30}. Nonetheless, the activities of those plasmonic Au NPs-based multi-component photocatalysts are still too low to meet the challenging demands of renewable fuel production^{29,30}. In addition, the synthesis of TiO₂ nanostructures still remains challenging to improve light absorption and suppress the recombination events. In contrast, platinum shows the highest catalytic activity for hydrogen evolution (HE) and moderate activity for oxygen evolution (OE). So far, only small Pt NPs that merely serve as a chemical catalyst are loaded on TiO₂ to facilitate HE^{21,22,29,30,34}. The plasmonic absorption of Pt NPs typically in the UV region has been the main obstacle for use of Pt NPs as a plasmonic enhancement agent.

Most recently, Bigall *et al.* have established some very effective methods to extend the localized surface plasmonic (LSP) absorption of Pt NPs to the visible and even near-IR regions of spectrum by enlarging their size^{35–37}. The reported Pt NPs exhibit strong and broad plasmonic absorption features. In the report of Bigall *et al.*, the plasmonic absorption maximum can be tuned from ~400 to 494 nm by increasing Pt NP size from 73 to 107 nm³⁵. Jung and co-workers prepared uniform Pt nanorods showing a transverse LSP mode appearing at ~400 nm and a longitudinal mode redshifting from 800 to above 1600 nm with increasing the length³⁶. In addition, the LSP features of dimeric Pt NPs are different from a single NP due to plasmonic coupling resulting in dramatic redshift (from 650 to 750 nm), enhancement, and broadening of LSP peak³⁷. Inspired by these investigations, we aim to study the ability and functionality of large plasmonic Pt NPs in photosensitization of semiconductor. As a result, to overcome the drawbacks of PM-NPs/TiO₂, herein, we develop a facile strategy to obtain a new plasmonic Pt/TiO₂-HA heterostructure by fabricating the novel TiO₂-HA (TiO₂ branched nanowires (b-NWs) epitaxially grown on TiO₂ microtubes) decorated by large Pt NPs with SPR in the visible region, which enables us to fulfill a dual function, i.e., control over photosensitivity and improving surface catalytic function with the aid of the plasmonic Pt NPs. Thus, photocatalytic overall water splitting can be realized without using any SR and additional co-catalysts. In particular, it should be stressed that to realize the practical application of VLWS, the development of photocatalysts working in a system free from any chemicals except the photocatalyst and water, as in this work, is much more desirable than the exploration of conventional photocatalysts in combination with sacrificial agents, hole consumers, or scavengers^{38–41}.

Results and Discussion

Morphological and structural characterizations of TiO₂-HA and Pt/TiO₂-HA. The TiO₂-HA is produced by hydrolysis of K₂TiO(C₂O₄)₂ in the presence of the additive NaH₂PO₂ and capping reagent diethylene glycol at 180 °C under solvothermal conditions. NaH₂PO₂ serves as a structure-directing agent to control the morphology of TiO₂-HA (see Figs S1 and S2, and the related discussion in the Supporting Information (SI)). The overview scanning electron microscopy (SEM) image in Fig. 1a unambiguously displays that the TiO₂-HA is produced in a 100% morphological yield. The characteristic feature of the TiO₂-HA is that very thin, highly dense TiO₂ b-NWs are radially (epitaxially) grown from the whole length of the primary TiO₂ microtube trunk with a length of 0.5–2 μm (see Fig. 1b,c and the discussion below). The transmission electron microscopy (TEM) image in Fig. 1c presents that the b-NWs have a length and a diameter in the range of 100–200 and 4–10 nm, respectively. In particular, the TEM image in Fig. 1c reveals that each TiO₂ trunk has a hollow interior with a shell thickness of 20–30 nm, as evidenced by the different contrasts between the middle and the two sides (also see Fig. 1g). The broken TiO₂-HA shown in Fig. 1a also corroborates the tubular superstructure of each trunk. All the TiO₂-HAs have a complete microtube trunk with both closed ends of which one is round and the other is flat (Fig. 1a and the inset in Fig. 1b). The outer diameter of the microtube trunks ranges from 150 to 250 nm. The ultrathin TiO₂ b-NWs provide a high specific surface area (S_{BET} of 150.2 m² g⁻¹, Fig. S3) for harvesting light, short distances for migration of charge carriers, and a large contact area with electrolyte, while the large central trunks enhance light scattering. These features of TiO₂-HA result in a significant extension of the light travelling distance within the ensemble of the TiO₂-HAs and therefore increase the probability of photons being absorbed by TiO₂-HA. Such a binary synergistic function of the different parts of TiO₂-HA could greatly promote its photocatalytic capability.

The high resolution TEM (HRTEM) image in Fig. 1d shows the well-resolved, continuous lattice fringes with the same spacing, indicating that each branch is single crystalline and has the same growth direction. The *d*-spacing of adjacent fringes along the axial direction of the b-NWs is consistently 3.775 Å, corresponding to the (100) planes of anatase phase, suggesting that all the b-NWs of the TiO₂-HA epitaxially grow exclusively along [100] directions. The XRD pattern in Fig. 2a illustrates that all the Bragg peaks of the TiO₂-HA can be perfectly indexed to pure anatase phase of TiO₂ with lattice constants *a* = 3.784 Å and *c* = 9.514 Å (space group: I41/amd (141), JCPDF no. 78-2486). Furthermore, the strong diffraction peaks confirm the good crystallinity of the TiO₂-HA. The surface chemical composition and electronic state of the TiO₂-HA are further analyzed using X-ray photoelectron spectroscopy (XPS). The XPS data demonstrate the expected elements Ti(IV) and O species for TiO₂ (Fig. S1 and the related discussion in SI).

To improve visible light harvesting of the TiO₂-HA, the uniform Pt NPs with a mean diameter of 73 or 107 nm were directly grown on the TiO₂-HA. Correspondingly, the integrated composites of the TiO₂-HA and 73 or 107 nm Pt NPs are designated as Pt73/TiO₂-HA or Pt107/TiO₂-HA, respectively. The Pt NPs loading is varied from 1% and 3% to 7% by weight with respect to TiO₂-HA. For simplicity, the composite samples will be referred to Pt73-*x*/TiO₂-HA and Pt107-*x*/TiO₂-HA in the following, where *x* represents the weight percentage of Pt NPs, e.g., Pt107-7/TiO₂-HA refers to the Pt107/TiO₂-HA containing 7 wt% 107 nm Pt NPs. The structural characterization results for Pt107/TiO₂-HA and Pt73/TiO₂-HA are very similar and therefore for brevity we mainly focus on the Pt107/TiO₂-HA and the results of the Pt73/TiO₂-HA are not shown. SEM images show the homogeneous distribution of the ~107 nm Pt NPs over the TiO₂-HA and verify that the number density of Pt NPs on

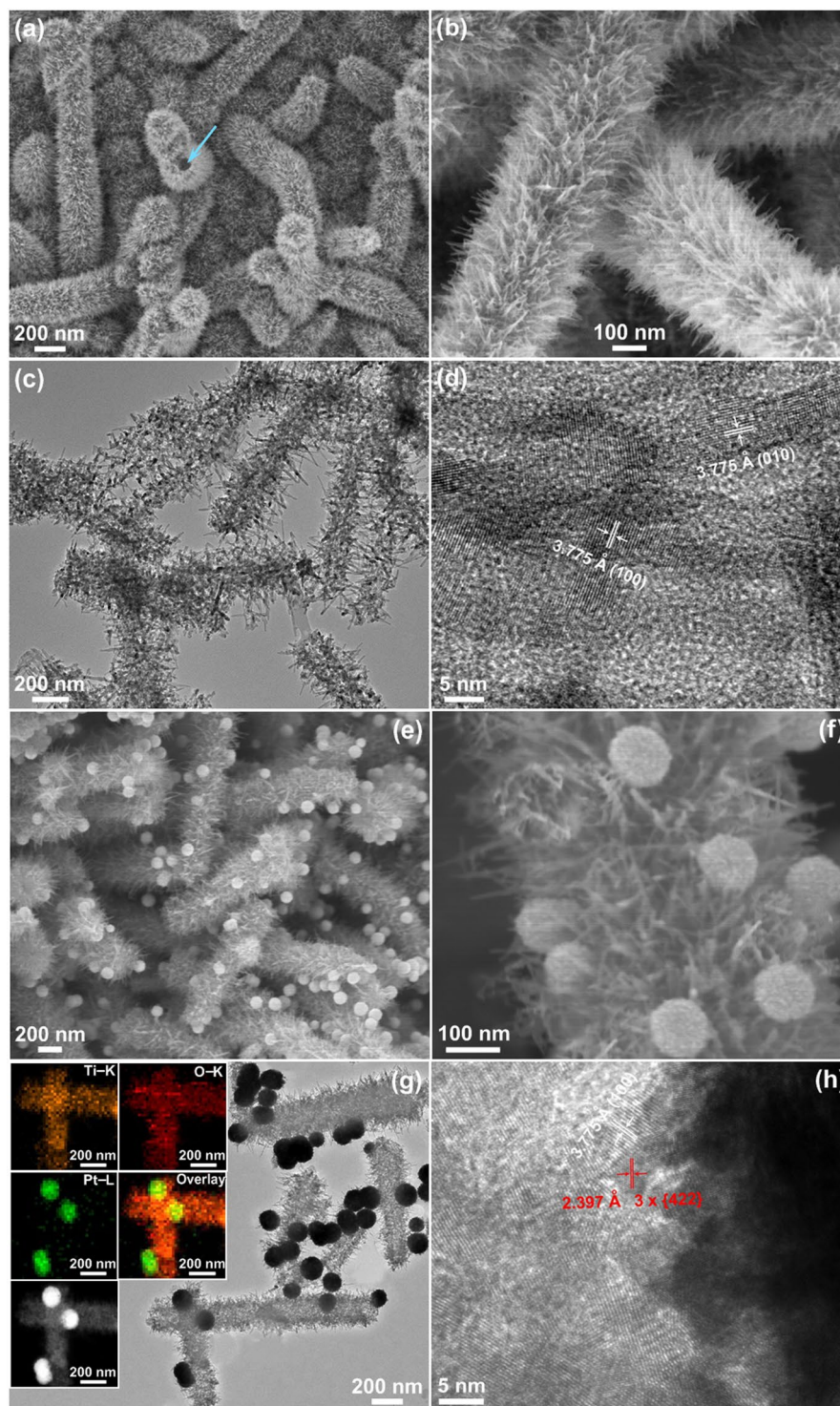


Figure 1. (a and e) Low- and (b and f) high-magnification SEM, (c and g) TEM, and (d and h) HRTEM images of (a–d) the TiO_2 -HA and (e–h) the Pt107-7/ TiO_2 -HA with 7 wt% loading of Pt NPs. The arrow in panel (a) points to the cavity of a broken hollow TiO_2 -HA. The insets: (b) SEM image showing the cross sections of several TiO_2 -HAs; (g) EDS elemental mapping images of Ti, O, and Pt, their overlay image, and the corresponding HAADF-STEM image (the bottom inset in panel (g)).

TiO_2 -HA increases with Pt NPs loading (Figs 1e,f and S4a,b). Representatively, the TEM image of the Pt107-7/ TiO_2 -HA in Fig. 1g reveals that Pt NPs are slightly embedded in the spacings among a number of b-NWs, as evidenced by the relatively slight contrast in the contact region between Pt NPs and the TiO_2 -HA compared to other domains of NPs. Such a embedding configuration could be advantageous for electric field enhancement

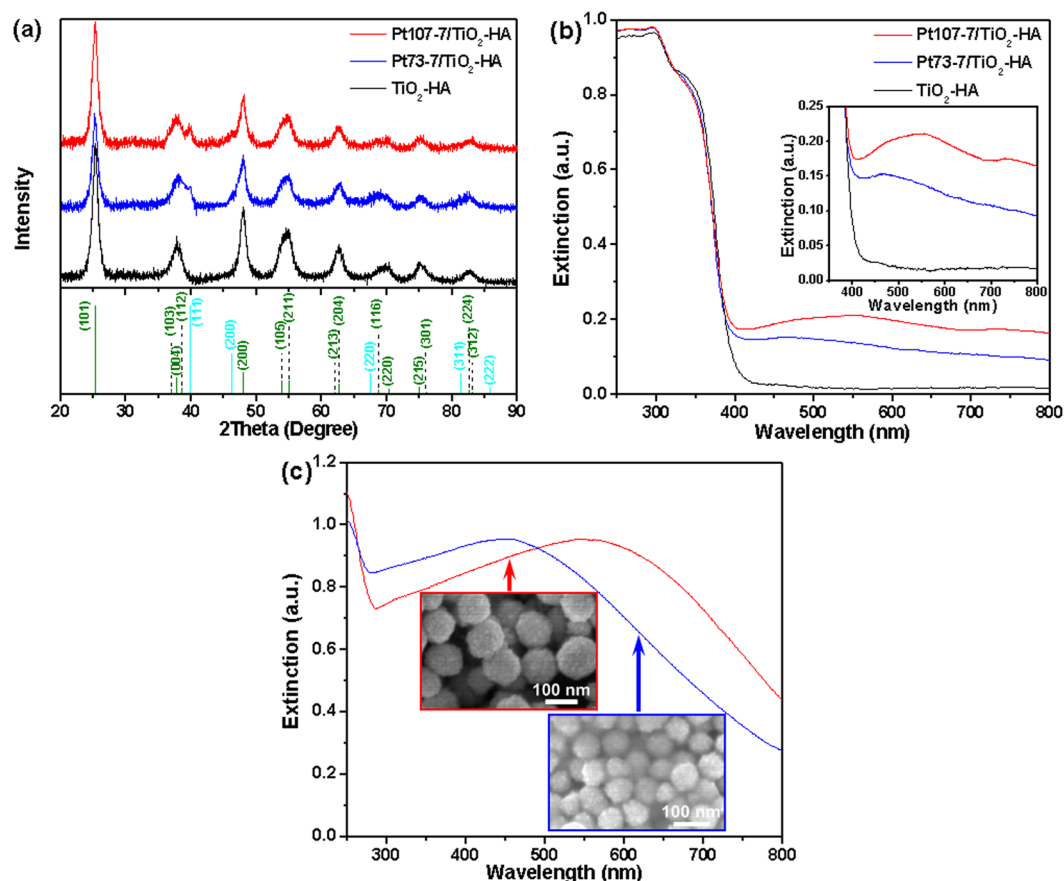


Figure 2. (a) XRD patterns of the TiO_2 -HA, Pt73-7/ TiO_2 -HA, and Pt107-7/ TiO_2 -HA. For comparison, the intensities and positions for the pure TiO_2 (olive bars, JCPDF No. 78-2486) and Pt (cyan bars, JCPDF No. 87-0646) references are given at the bottom according to the JCPDS database. UV-vis extinction spectra of (b) the TiO_2 -HA, Pt73-7/ TiO_2 -HA, and Pt107-7/ TiO_2 -HA and of (c) the Pt NPs with a mean diameter of 73 nm (blue curve, $\lambda_{\text{max}} = 455$ nm) and 107 nm (red curve, $\lambda_{\text{max}} = 542$ nm). The insets: (b) a zoom-in view of optical extinction in the visible region; (c) SEM images of the corresponding Pt NPs.

and more efficient transfer of hot electrons^{19,24,42}. The corresponding energy dispersive X-ray spectroscopy (EDS) elemental-mapping data in the insets in Fig. 1g and EDS spectrum in Fig. S4c support that the TiO_2 -HA is comprised of Ti and O atoms and decorated with Pt NPs. The quantitative EDS analysis also confirms the concentration of Pt NPs in each composite sample is consistent with the initial designated level (Table S1). In addition, the enhanced elemental contrast shown by the high annular dark-field scanning TEM (HAADF-STEM) image also implies that Pt NPs were successfully loaded on TiO_2 -HA (the bottom left inset in Fig. 1g).

Representatively, the close-up HRTEM image in Fig. 1h shows the interfacial area of the Pt107-7/ TiO_2 -HA composite. Two sets of distinctive lattice planes, which can be assigned to anatase and Pt, can be easily discerned. The well-resolved (100) planes of the continuous TiO_2 shell domain can be clearly identified, which proves the epitaxial growth of all the b-NWs along [100] directions. The lattice fringes observed in the margin of Pt NPs correspond to a periodicity of 2.40 Å, which can be ascribed to (1/3){422} forbidden reflections. Such forbidden reflections have been frequently observed for fcc Au and Ag NPs that are grown by seeded methods and can be attributed to the generation of twins and stacking faults in the fcc lattice⁴³. For comparison, the XRD patterns of the Pt107-7/ TiO_2 -HA and Pt73-7/ TiO_2 -HA present the same reflection peaks originating from TiO_2 and Pt constituents (Fig. 2a), further confirming the formation of the Pt/ TiO_2 -HA composite. However, no definitive diffraction peaks from Pt NPs can be identified for the Pt107/ TiO_2 -HA and Pt73/ TiO_2 -HA samples when the content of Pt is lower than 5 wt%. The formation of metallic Pt in Pt107-7/ TiO_2 -HA composite is further verified by XPS analysis (Fig. S5).

Figure 2b compares the extinction spectra of the TiO_2 -HA, Pt73-7/ TiO_2 -HA, and Pt107-7/ TiO_2 -HA. As expected, the pure TiO_2 -HA sample exhibits bandgap-dependent absorption only in the UV region ($\lambda < 400$ nm). The extended tail beyond 400 nm in the TiO_2 -HA extinction spectrum stems from its light scattering behavior rather than intrinsic optical absorption (see discussion in the later section). In contrast, the Pt107-7/ TiO_2 -HA sample exhibits a pronounced, broad optical absorption hump in the visible region. For comparison, large Pt colloids with a similar diameter of 73 or 107 nm, which were prepared according to the procedure reported by Bigall *et al.*³⁵, show a plasmonic absorption peak at 455 or 542 nm, respectively (Fig. 2c). It is worth noting that

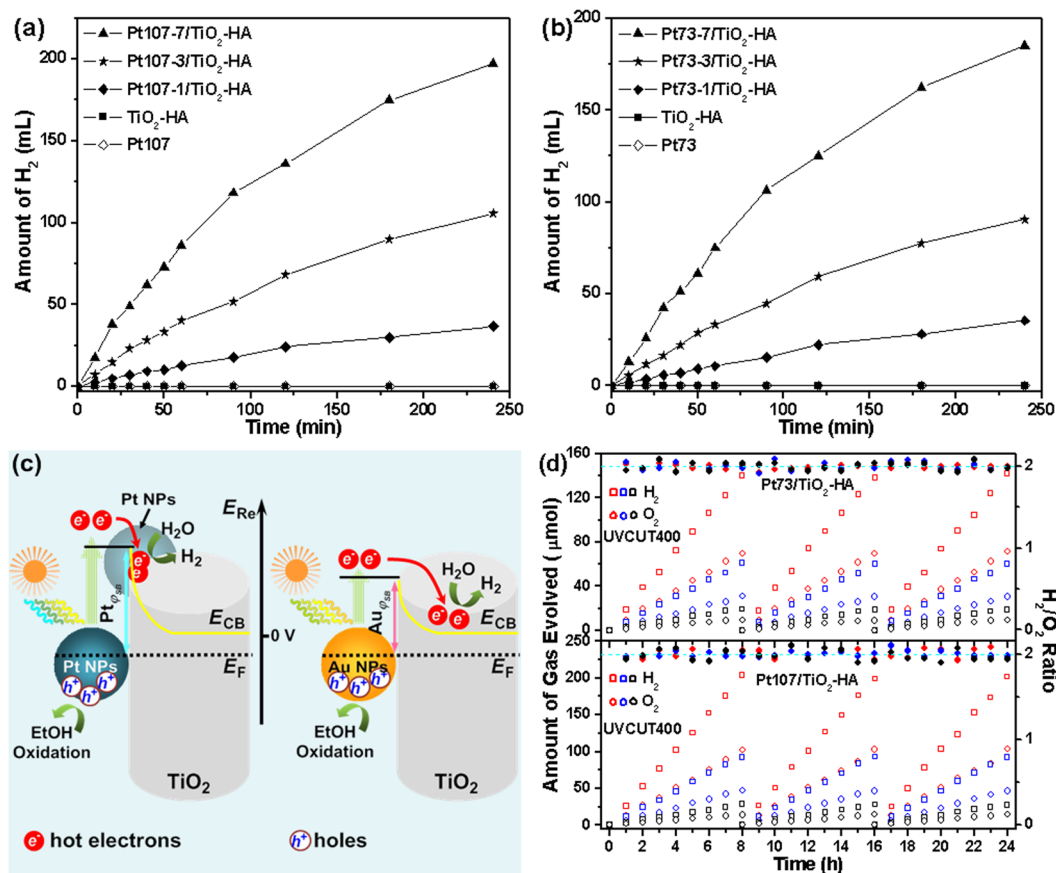


Figure 3. Time course of evolved H₂ from a 25 vol% water/ethanol system under visible light ($\lambda > 400$ nm) irradiation in the presence of (a) the Pt107/TiO₂-HA and (b) the Pt73/TiO₂-HA photocatalysts containing varied Pt NPs loading. (c) Schematic illustration of φ_{SB} -dependent hot electron generation and transfer from Pt or Au NPs to TiO₂. E_F , E_{CB} , and E_{Re} represent the energies of the Fermi level, TiO₂ conduction band, and chemical reduction potential, respectively. For clarity, the exciting lights are labeled as different colors approximately corresponding to the varied colored light absorbed by the plasmonic metal NPs. (d) Photocatalytic H₂ and O₂ evolution from pure water in 3 repeated cycles (8 h/cycle) using various Pt/TiO₂-HA (data point: red, Pt107-7 and Pt73-7; blue, Pt107-3 and Pt73-3; black, Pt107-1 and Pt73-1) photocatalysts under visible light ($\lambda > 400$ nm) irradiation. The testing system was evacuated between two cycles in turn.

compared to the plasmonic absorption of the corresponding colloidal Pt NPs, the extinction bands of the Pt73-7/TiO₂-HA and Pt107-7/TiO₂-HA are red-shifted by 11 and 13 nm, respectively. Furthermore, the band of the Pt107-7/TiO₂-HA is broadened and extended even further to the near-infrared (NIR) region ($\lambda > 700$ nm, the inset in Fig. 2c). These changes in the extinction spectra, especially for the Pt107-7/TiO₂, can be attributed to a superposition of the enhancement of local electric field near the interface between the Pt NPs and TiO₂-HA and scattering of photons due to the larger Pt NPs^{19,24}. Obviously, similar to the Au NPs^{19,22}, the bigger Pt NPs, the larger near field enhancement at the Pt NP-TiO₂-HA interface because the SPR intensity increases with particulate size. For both the Pt NPs-supported samples, their absorption intensity increases in the visible region with the concentration of supported Pt NPs in our experimental range (Fig. S6).

Photocatalytic performance of Pt/TiO₂-HA. Figure 3a,b illustrate the rates of photocatalytic HE from an aqueous solution of ethanol (4: 1 water/ethanol in v/v) for various Pt/TiO₂-HAs and pure TiO₂-HA as a function of time under visible light ($\lambda > 400$ nm) illumination. The utilization of a conventional sacrificial agent (ethanol) for evaluating the photocatalytic activity of various Pt/TiO₂-HAs provides a consistent comparison with those results from the Au/TiO₂ photocatalysts in the literature. Pt NPs and TiO₂-HA do not exhibit any activity (Fig. 3a,b), indicating that TiO₂-HA exhibits no optical absorption in the region of $\lambda > 400$ nm. No HE is also detected over the TiO₂-HA decorated by small Pt NPs with ~ 25 nm in size (Fig. S7), which enables adequate suppression of the e⁻-h recombination in TiO₂-HA, further confirming that TiO₂-HA exhibits no response to visible light. All of the Pt107/TiO₂-HA samples present a higher HE rate and a larger mass activity (MA) than the corresponding Pt73/TiO₂-HA sample with the same Pt NP loading (also see Table S2). Therefore, the 107 nm Pt NPs with a larger field enhancement appears to further promote the photocatalytic activity. For the samples with Pt loading not less than 3 wt%, their total MAs are much higher than those of the Au/TiO₂ photocatalysts reported previously (Table S2). Notably, the PM-based MAs of all the samples are at least 3-fold higher than the previous Au/TiO₂ (Table S2). The greater photocatalytic activity achieved by the Pt/TiO₂-HA indicates that the

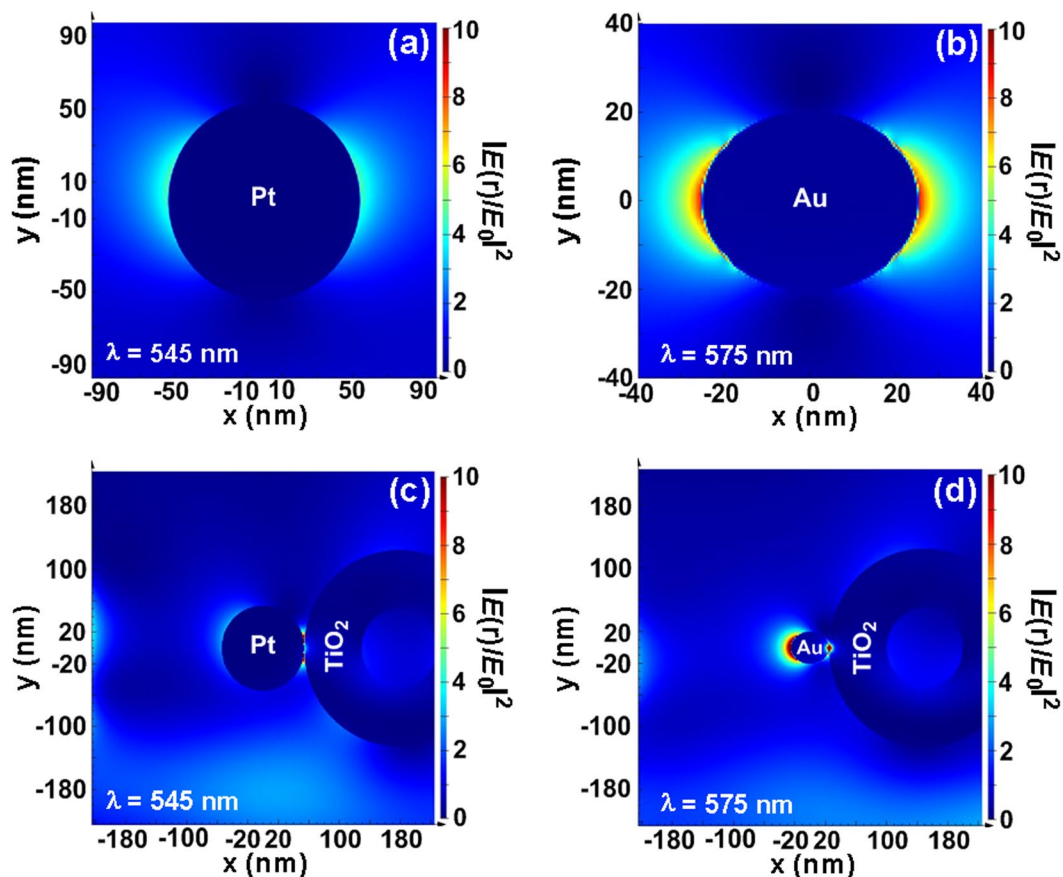


Figure 4. Plasmonic near-field distribution maps (cross-section view at $z=0$) of (a) Pt107 nanospheres and (b) ellipsoidal Au NPs with major axis of 60 nm and minor axis of 40 nm simulated using finite difference time domain methods, which show the evanescent electric near-field intensity enhancements $|E(r)/E_0|^2$ of the Pt and Au NPs at 545 and 575 nm, respectively. Electric field intensity enhancement contours of (c) the Pt107/TiO₂-HA and (d) Au_{50nm}/TiO₂-HA hybrid systems under the parallel polarized excitation. The color scale bar on the right side of each panel shows the relative increase in field enhancement. The geometrical shapes of the Pt107, Au_{50nm}, and TiO₂-HA are established based on the TEM images.

plasmonic Pt NPs could play an important role in enhancing photocatalysis besides the favorable TiO₂-HA morphology discussed above. In particular, unlike the Au/TiO₂ photocatalysts (a volcano curve for activity versus Au loading)²⁸, the HE rate monotonously rises with Pt loading within our experimental range, suggesting that the activity sites simultaneously increase because Pt also acts as a highly active HE catalyst. The PM-based MA of each Pt/TiO₂-HA with the same Pt size approves this proposition because it does not differ much by altering Pt loading (Table S2).

In the case of Au/TiO₂ photocatalysts, their enhanced activities endowed by Au NPs have recently been elucidated by the following mechanisms: (i) the enhanced optical absorption of TiO₂ driven by SPR-magnified electromagnetic fields, (ii) SPR-induced hot electron injection from Au NPs to TiO₂, and (iii) resonant photon scattering^{22,28}. However, the nature is different in the plasmonic Pt/TiO₂ photocatalysts. First, the plasmonic field enhancement from Pt SPR is lower than that from Au and Ag counterparts³⁵. In our case, obviously, the corresponding optical-absorption maps display that the Pt107 NPs exhibit a much weaker plasmonic near-field than the Au_{50nm} NPs (i.e., Au NPs with a mean diameter of 50 nm and an ellipsoidal shape to conform to the following synthetic results), as evidenced by performing numerical simulations using the finite-difference time-domain (FDTD) method (Fig. 4a,b). To further compare the nature of the plasmon mode at the Pt107/TiO₂-HA interface with that at the Au_{50nm}/TiO₂-HA interface, the FDTD simulations were conducted to calculate the electric field enhancement at each interface, in which parallel incident light polarization direction is applied to the Pt107 and Au_{50nm} NPs. The electric field intensity enhancement contour reveals that the enhancement of the local electric field is mediocre at the Pt107/TiO₂-HA interface, whereas a spatially confined “hot spot” appears at the Au_{50nm}/TiO₂-HA interface (Fig. 4c,d). Therefore, the local electric field enhancement has insignificant impacts on the photocatalytic activity of the Pt107/TiO₂-HA system. Second, the Schottky barrier height (φ_{SB}) of Pt/TiO₂ junction ($\varphi_{SB(Pt/TiO_2)} = 1.7$ eV)^{44,45} is higher than that of Au/TiO₂ ($\varphi_{SB(Au/TiO_2)} = 0.9$ –1.0 eV)⁴⁶. These two unique physical properties of Pt/TiO₂ suggest that the photocatalytic HE process for the Pt/TiO₂-HAs is mechanistically different from that of Au/TiO₂ photocatalysts. Despite the relatively low field enhancement afforded by the supported Pt NPs and no optical absorption for TiO₂-HA in the visible region, all the Pt/TiO₂-HA samples show a

much higher PM-based MA than the best Au-TiO₂ photocatalysts reported by different groups^{19,24}. Therefore, the mechanism of local enhanced electric field at the Pt-TiO₂-HA interface should not be responsible for the higher photocatalytic activity of the Pt/TiO₂-HA. In addition, the increase in absorbed photons scattered by TiO₂-HA itself and large Pt NPs plays a certain but not dominant role in boosting photocatalytic activity.

On the other hand, the larger $\varphi_{SB(Pt/TiO_2)}$ results in the hot electrons transferred from the Pt NPs to occupy the higher energy levels in the TiO₂-HA conduction band (CB) in comparison with Au/TiO₂ photocatalysts, giving rise to the more negative potential of the accumulated hot electrons (see Fig. 3c). This is because the higher $\varphi_{SB(Pt/TiO_2)}$ at the Pt-TiO₂-HA interface effectively prevents the injected hot electrons from striding over the Schottky barrier and transferring back to recombine with the holes left within Pt NPs and thereby allows the hot electrons to survive on a very long time scale. Thus, compared to Au/TiO₂, the increased population of hot electrons due to their prolonged lifetimes increases the probability of the injected electrons to occupy the higher energy levels in the TiO₂-HA CB and enhances their reduction potentials (Fig. 3c)²². This hypothesis is strongly supported by monitoring the open circuit voltage decay versus time characteristics, as recently introduced by DuChene *et al.*⁴⁷, which reveals the longer hot electron lifetime for the Pt107-7/TiO₂-HA electrode than that for the Au_{50nm}/TiO₂-HA (i.e., 50 nm Au NPs-decorated TiO₂-HA (8.3 wt% Au)) (Fig. S8 and the related discussion in SI). The energy of SPR in Pt NPs is sufficiently high and can be much larger than the $\varphi_{SB(Pt/TiO_2)}$ (1.7 eV). For example, it is estimated to be 3.26 eV based on the recent report of Lin *et al.*⁴⁸. In our case, the SPR absorption maxima of the Pt73 and Pt107 correspond to the photon energy of 2.7 and 2.3 eV, respectively. At the same time, the plasmonic hot electrons generated by our Pt NPs are excited by a large fraction of shorter wavelength light compared to Au/TiO₂. Accordingly, our experimental results suggest that the excited carriers are energetic and momentum enough to surmount the Pt-TiO₂-HA interface, so that a sufficient amount of the hot electrons can be injected into the TiO₂-HA CB and accumulated there to increase reduction potentials that enable the higher water reduction rate than that on Au/TiO₂.

Meanwhile, the Pt NPs with intrinsic excellent catalytic activity also significantly raise the rate of water photoreduction. Among various catalysts, Pt exhibits the best catalytic activity towards HE (i.e., the overpotential close to zero for H₂ formation and facile H₂ desorption) and thus the most favorable HE kinetics and a very high energetic efficiency. Therefore, the recombination of a considerable portion of the hot electrons in the TiO₂-HA CB and holes on Pt NPs can be further suppressed by the extremely high water reduction rate (i.e., rapid consumption of electrons at more negative potentials) considering the competition between hot electrons reducing water and recombining with those holes. In Fig. 3c, for clarity, an additional Pt NP is schematically depicted to highlight the HE reaction facilitated by coupling of the catalytic function of the plasmonic Pt NPs and the favourable reduction potentials of the hot electrons. Concomitantly, the injected hot electrons are inclined to transfer from their original host NP into other guest Pt NPs by tunneling due to the electron migration driven by the repulsive forces that are induced by the accumulated electrons (i.e., the electric field gradient) in the TiO₂-HA CB^{21,49}, and moreover, other Pt NPs are more accessible to the injected hot electrons than the original host NP (namely, the lower possibility of hot electrons diffusing back to the host NP) because of their larger population. In contrast, Au shows inappreciable catalytic activity towards HE, so that only a small part of hot electrons are able to participate in the water photoreduction, leading to the lower HE rates. Meanwhile, the hot electrons rapid relax from the high energy states determined by the $\varphi_{SB(Au/TiO_2)}$ to the level of the conduction band minimum and eventually the majority of hot electrons recombine with the holes on Au NPs^{47,50}. Unlike the plasmonic Pt/TiO₂-HA, water photoreduction proceeds on the surface of the TiO₂ moiety of Au/TiO₂ because of the transferred hot electrons lying in the TiO₂ CB and the catalytically inactive “by-stander” Au moiety.

To further confirm our new mechanism, two control samples including Au_{50nm}/TiO₂-HA and Au_{50nm}/TiO₂-HA-Pt_{25nm} (i.e., Au_{50nm}/TiO₂-HA containing 25 nm Pt NP co-catalyst (6.8 wt%)) were prepared and used as photocatalysts for water reduction under otherwise the same conditions as those used in testing of the Pt/TiO₂-HA (Figs S9, S10, and S11a). It should be mentioned that all of the TiO₂-HA supports show nearly the same absorption spectra, which lays a solid foundation to make a reliable comparison among photocatalytic activities of different samples (Figs 2b and S10b). In the case of the Au_{50nm}/TiO₂-HA-Pt_{25nm}, the incorporation of small Pt NPs provides the comparable chemical catalytic effect on HE to that of the Pt107-7(Pt73-7)/TiO₂-HA such that the net hot-electron effect can be exactly extracted from the Pt/TiO₂-HA. The PM-based MAS of all the Pt/TiO₂-HA samples are not only much higher than the control samples, especially the Au_{50nm}/TiO₂-HA-Pt_{25nm} (Table S2) but also essentially independent of Pt NP loading, which further validates the mechanism of hot electron reduction regulated by the Pt-TiO₂-HA interface.

Concomitantly, for our plasmonic Pt/TiO₂-HA samples, the population of holes with more positive potentials is expected to increase in comparison with Au/TiO₂, presumably due to the higher plasmonic energy obtained from a portion of exciting light with higher photon energies (Fig. 3c). On the other hand, the presence of large Pt NPs with small specific surface areas is favourable for concentrating charge carries (i.e., holes) in the vicinity of NP surface, that is to say, increasing the surface concentration of holes, because an order of magnitude increase in surface area can diminish the obtainable photovoltage by 59 mV^{1,51}. It can be deduced that the potential of a fraction of the accumulated holes on the large Pt NPs can be higher than the value of *ca.* 1.6 V (*vs.* reversible hydrogen electrode) required to achieve the onset of water oxidation on Pt⁵². In addition, the sufficiently high $\varphi_{SB(Pt/TiO_2)}$ can effectively inhibit the recombination of the hot electrons in the TiO₂-HA CB with the holes left on the Pt NPs and thus permits the holes to accumulate and occupy the higher energy levels in the valence band of Pt based on the time scale of the hot electron lifetime discussed above, which also contributes to the enhanced oxidation ability. Therefore, this increases the open circuit photovoltage within photocatalyst particles by decreasing the photocatalyst/electrolyte interfacial area, leading to the improvement of water oxidation rate. A further conclusion can be deduced from the experimental results discussed below.

The evaluation of photocatalysts for practical applications requires that photocatalytic activity should be measured in pure water without any sacrificial agents. Interestingly, significant gas generation from pure water is

observed over the Pt/TiO₂-HA samples under visible ($\lambda > 400$ nm) light irradiation. Figure 3d depicts the time courses of HE and OE obtained in a suspension of each Pt/TiO₂-HA photocatalyst in pure water. There is a proportional increase in the amounts of H₂ and O₂ with irradiation time within each illumination cycle and H₂/O₂ ratio is very close to 2, i.e., a stoichiometry, during the entire measurement period. As expected, the rate of water splitting increases with the loading of Pt NPs as well as with Pt NP size (Table S3). The largest HE and OE rates of 25.2 and 12.9 $\mu\text{mol h}^{-1}$, respectively, are obtained on the Pt107-7/TiO₂-HA sample and the corresponding total MAs and PM-based MAs are 2521 $\mu\text{mol h}^{-1} \text{g}^{-1}$ and 34.3 $\mu\text{mol h}^{-1} \text{mg}_{\text{PM}}^{-1}$, respectively. In a similar way, the Pt/TiO₂-HA samples have nearly the same PM-based MA when the size of Pt NPs is the same. In contrast, the Au_{50nm}/TiO₂-HA-Pt_{25nm} sample exhibits a far lower activity (i.e., a HE and an OE rate of 0.54 and 0.27 $\mu\text{mol h}^{-1}$, respectively, and a PM-based MA of 0.65 $\mu\text{mol h}^{-1} \text{mg}_{\text{PM}}^{-1}$) under visible light irradiation while the Au_{50nm}/TiO₂-HA does not show a perceivable activity in pure water (Fig. S11b and Table S3). Also, the Pt107-7(Pt73-7)/TiO₂-HA photocatalysts give a much higher HE/OE rate and show a larger PM-based MA than other Au-based plasmonic TiO₂ photocatalysts in the literature (see Table S3)^{29,30}. These results suggest that the plasmonic Pt NPs induce the hot electrons with more negative potentials as well as holes with more positive potentials and concomitantly serve as efficient chemical catalysts for HE and OE reactions. The Pt/TiO₂-HA photocatalysts are robust and therefore have a high stability for catalytic water splitting. The large specific area of TiO₂-HA and its branched features lead to a high affinity to Pt NPs so as to form a robust composite catalyst. Thus, various Pt/TiO₂-HA samples are highly stable and do not show any degradation in the photocatalytic activity after consecutive cycles of irradiation (Fig. 3d).

Note that it is possible for our Pt NPs to catalyze the reverse reaction of water splitting, i.e., the combination of H₂ and O₂ into H₂O³⁰. To test the net catalytic effects of the Pt107 and Pt73 NPs, the rate of the reaction between H₂ and O₂ is examined in the dark using these two kinds of Pt NPs as catalysts. A very small rate of H₂/O₂ consumption for each catalyst is obtained, as shown in Fig. S12. This indicates that the very large Pt NPs with low surface areas is rather inactive towards the reverse reaction, hence fully contributing to water splitting.

To further corroborate the Schottky barrier height hypothesis in the photocatalytic activity enhancement of the plasmonic Pt/TiO₂-HA in the visible region, the photoelectrochemical (PEC) properties of the Pt107-7/TiO₂-HA and Au_{50nm}/TiO₂-HA-Pt_{25nm} film electrodes are compared by conducting the photocathode half-reaction of the water-splitting process on the film electrodes and the water oxidation reaction on an un-illuminated platinum wire anode according to the method reported by Mubeen and co-workers²⁹. The photocurrent density of the Pt107-7/TiO₂-HA film electrode is considerably larger (1.8 times) than that of the Au_{50nm}/TiO₂-HA-Pt_{25nm} film electrode at 0 V vs. reversible hydrogen electrode (RHE) under visible light irradiation obtained with a cut-off filter L-42 (see Fig. S13a,b). Meanwhile, electrochemical impedance spectroscopy (EIS) was conducted on these two photoelectrodes under the same illumination. A single semicircle is obtained at an applied potential of 0 V vs RHE (Fig. S13c). Concomitantly, the complex nonlinear least square (CNLS) fitting of the EIS is performed with the Zview 3.1 software package (Fig. S13c). Obviously, there is a good agreement between the experimental data (symbols) and CNLS approximations (solid lines) when the Randles equivalent circuit model is applied, as indicated by the Chi-squared value, χ^2 (Table S4). The charge transfer resistance (R_{ct} , 16.7 $\text{k}\Omega \text{cm}^{-2}$) of the Pt107-7/TiO₂-HA photoelectrode is far lower than that (30.6 $\text{k}\Omega \text{cm}^{-2}$) of the Au_{50nm}/TiO₂-HA-Pt_{25nm} photoelectrode (i.e., the faster reaction rate for the Pt107-7/TiO₂-HA) (Table S4), consistently indicating the better photoreduction ability of plasmonic Pt hot electrons. Considering that the photocurrent density faithfully tracks the rate of hydrogen formation, the striking discrepancy in the photocurrent density between the Pt107-7/TiO₂-HA and Au_{50nm}/TiO₂-HA-Pt_{25nm} is presumably due to the fact that the higher Schottky barrier of Pt/TiO₂-HA boosts its photoreduction ability. Concurrently, to probe the oxidation ability of holes, the Pt107-7/TiO₂-HA and Au_{50nm}/TiO₂-HA-Pt_{25nm} photoanodes are tested for PEC oxidation of water and the water reduction reaction on an un-illuminated platinum wire cathode. The obtained order of photoanode current density support the proposition that there are at least a fraction of holes photogenerated from Pt107-7/TiO₂-HA have more positive potentials than those from Au_{50nm}/TiO₂-HA-Pt_{25nm} (Fig. S13d). Similarly, the corresponding EIS data reveal that R_{CT} of the Pt107-7/TiO₂-HA photoanode is much smaller than that of the Au_{50nm}/TiO₂-HA-Pt_{25nm} counterpart, confirming the aforementioned conclusion (data are not shown for brevity).

A more elaborate evaluation of the Pt NPs SPR-driven hot electron injection mechanism, the time course of water splitting under irradiation of visible light with longer wavelengths, which are obtained by L-42, Y-48, and O-54 cut-off filters, is shown in Fig. 5a. Evidently, the rates of HE and OE decrease with irradiation through the filter with the longer cut-off wavelength (i.e., in the order of L-42 > Y-48 > O-54), which can be attributed to the diminution of light absorption. As stated above, an increase in the loading of plasmonic Pt NPs increases both the hot electrons injected and catalytic active sites. As a result, there is an approximate linear correlation between the loading of plasmonic Pt NPs and the rates of H₂/O₂ formed, as shown in Fig. 5b. Figure 5c,d representatively illustrate the action spectra of the Pt107-7/TiO₂-HA and Pt73-7/TiO₂-HA photocatalysts, in which H₂ formation is achieved by conducting photocatalytic tests under monochromatic irradiation of visible light. Note that compared to the intensity of corresponding extinction spectrum, the relative HE rates of these two photocatalysts obtained under red and near-IR light illumination are lower. This discrepancy can be attributed to a superimposition of both the significant light-scattering caused by the larger Pt NPs⁴⁹ and the concurrent pronounced absorption from intraband transitions of Pt NPs^{15,51}, which induce no photochemical reactions, on the broad Pt LSPR peak. Thus, the HE rate is in good accord with the plasmon absorbance spectrum of the corresponding photocatalyst. These results undoubtedly support the proposition that the VLWS on Pt/TiO₂-HA photocatalysts stems from the hot electrons injection excited by the SPR of the Pt NPs on TiO₂-HA.

Apparent quantum efficiency (AQE) is an important parameter to evaluate the photocatalytic activity of a photocatalyst and is calculated in terms of the following equation:

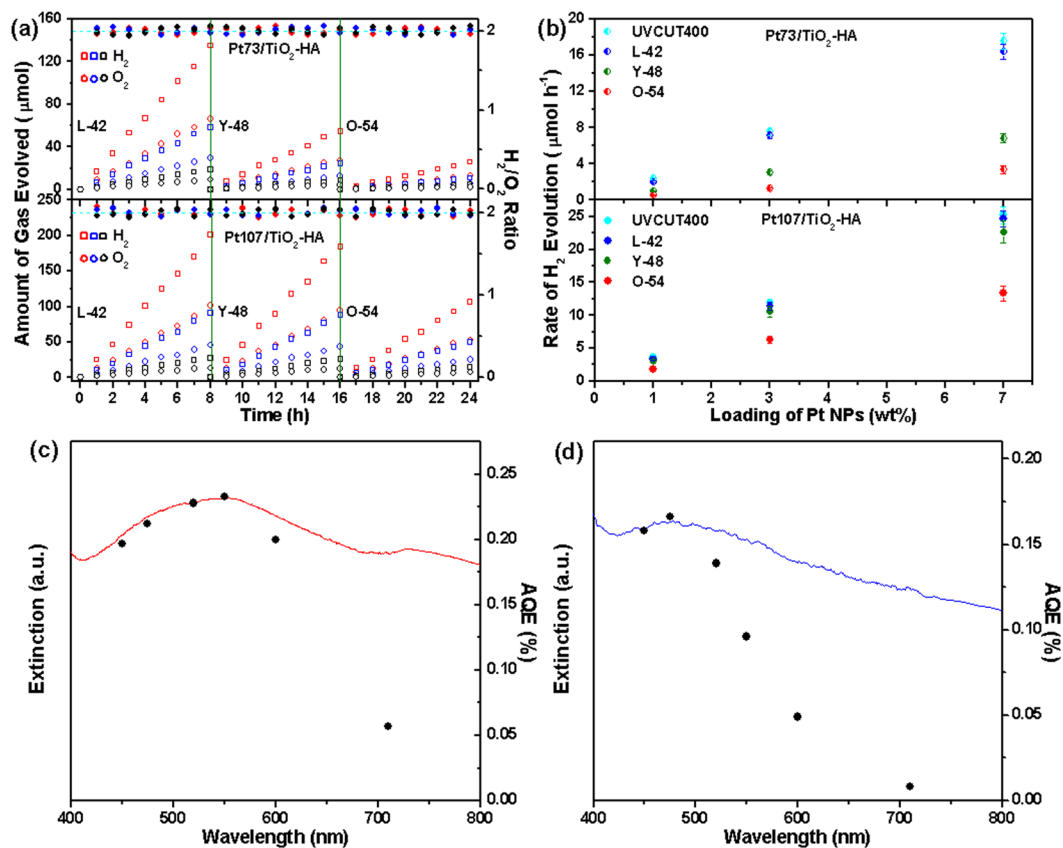


Figure 5. (a) Photocatalytic H₂ and O₂ evolution from pure water over time using various Pt/TiO₂-HA (data point: red, Pt107-7 and Pt73-7; blue, Pt107-3 and Pt73-3; black, Pt107-1 and Pt73-1) photocatalysts (i.e., aqueous suspensions of various Pt/TiO₂-HA samples) under illumination of visible light obtained from a Xe lamp equipped with different cut-off filters. (b) Effect of Pt loading amounts on the rate of H₂ evolution from pure water under irradiation of visible light obtained using different cut-off filters. Extinction spectra (left axis) and action spectra (right axis) of (c) Pt107-7 and (d) Pt73-7 in pure water. The net extinction spectra of Pt NPs in panel (c) and (d) are obtained by subtracting of TiO₂-HA from the corresponding extinction spectrum of Pt/TiO₂-HA.

$$\text{AQE} = \frac{2 \times (\text{number of H}_2 \text{ molecules})}{\text{number of incident photons}} \times 100 \quad (1)$$

The spectrally averaged photon energy (the visible portion ($\lambda > 420 \text{ nm}$)) is *ca.* $3.98 \times 10^{-19} \text{ J}^{29}$. In our case, the amount of photons irradiated is measured to be $0.5 \mu\text{mol cm}^{-2} \text{ s}^{-1}$, where the corresponding illumination intensity is 120 mW cm^{-2} . The average AQEs of the Pt107-7/TiO₂-HA and Pt73-7/TiO₂-HA photocatalysts are calculated to be 0.18% and 0.12%, respectively. The AQE values are much larger than the value (0.013%) in ref.²⁹ and comparable to the value (0.1%) in ref.²⁸. However, it is believed that the value (0.1%) in ref.²⁸ is questionable considering the extremely low rate of H₂ evolution ($0.25 \pm 0.05 \mu\text{mol h}^{-1}$) and higher total illumination intensity (300 mW cm^{-2}). Furthermore, the AQEs reach 0.23% at 550 nm and 0.17% at 475 nm for the Pt107-7/TiO₂-HA and Pt73-7/TiO₂-HA, respectively (Fig. 4c,d). Clearly, the rather broad SPR absorption of Pt NPs leads to the moderate increase in the AQE around the maximum in the plasmon spectrum in comparison with the corresponding average AQE. In addition, as anticipated, the AQE of the Pt107-7/TiO₂-HA is more than one order of magnitude higher than that of the Au_{50nm}/TiO₂-HA-Pt_{25nm} at the corresponding wavelength (Fig. S11c). However, it is not surprising that the AQE of the Pt107-7/TiO₂-HA is still too low, considering that the photocatalyst mainbody TiO₂-HA does not absorb visible light at all. Thus, usage of modified TiO₂-HA with significant visible-light absorption ability will be an effective pathway to further improve the AQE.

To further evaluate the long-term stability and durability of Pt/TiO₂-HA, the Pt107-7/TiO₂-HA and Pt73-7/TiO₂-HA photocatalysts are tested for 8 h each day and then used again by centrifugation of their aqueous suspension. The process is continually repeated for a week. The photocatalysts do not show any perceptible degradation in photocatalytic activity after one week of operation because of the highly robust nano-architecture of the Pt/TiO₂-HA photocatalysts as well as the exceptional chemical stability and anti-photocorrosion ability of Pt and TiO₂ (see Fig. S14).

Conclusions

In summary, the Pt/TiO₂-HA heterostructures are successfully fabricated and employed as efficient photocatalysts for plasmon-enhanced visible light hydrogen production. The plasmonic Pt NPs not only effectively activate the photocatalytic activity of pure TiO₂ in the visible region but also eliminate the use of additional chemical co-catalysts for HE and OE reactions in the VLWS process. Hence, the Pt/TiO₂-HA heterostructure is expected to behave as an emerging powerful photocatalyst with enhanced photocatalytic activity and a simple composition for sustainable development of solar energy conversion. Our studies pave the way for the development of efficient new plasmonic photocatalysts by extending and optimizing plasmonic species coupled with tailoring unique nanostructures. At the same, it is expected that the plasmonic Pt nanostructures can be extended to more types of semiconductor supports, such as α -Fe₂O₃, for further optimizing the performance of plasmonic photocatalysts to drive the unassisted full water splitting reaction.

Methods

Chemicals. Potassium titanium oxide oxalate dihydrate (K₂TiO(C₂O₄)₂·2H₂O, ≥98% Ti basis), hydrogen hexachloroplatinate(IV) hexahydrate (H₂PtCl₆·6H₂O, 99.9%), sodium borohydride (NaBH₄, 98%), sodium hypophosphite monohydrate (NaH₂PO₂·H₂O, 99%), diethylene glycol (99%), and urea (H₂NCONH₂, 99%) were commercially available from Sinopharm Chemical Reagent Co., Ltd. Hexadecyltrimethylammonium bromide (CTAB, >99%), L-ascorbic acid (L-AA, >99%), sodium citrate tribasic dihydrate (99%), and citric acid (99.5%) were purchased from Sigma-Aldrich. All reagents were used without any further purification. Ultrapure water (18.2 M Ω) produced with a Milli-Q purification system was used in the synthesis and photocatalytic measurements.

Synthesis of TiO₂-HA. Typically, K₂TiO(C₂O₄)₂·2H₂O (156 mg) and NaH₂PO₂·H₂O (42.4 mg) were added to a mixed solvent of water (5.0 mL) and diethylene glycol (15.0 mL), and then sonicated for 10 min. Subsequently, the reaction mixture was loaded into a 40 mL Teflon-lined stainless steel autoclave and vigorously stirred by a magnetic stirrer at ambient temperature for 30 min. Next, the autoclave was sealed and maintained at 180 °C for 12 h and subsequently cooled to room temperature naturally. Finally, the white sediments at the bottom of the autoclave were isolated by centrifugation of the resulting mixture at 2000 rpm for 5 min followed by washing with hot water and ethanol (80 °C) alternatively for three times to purify the final product. The purified sediments were dried at 100 °C overnight to obtain the anatase-structured TiO₂-HA for further characterization and preparing photocatalysts.

Synthesis of plasmonic Pt NPs and Pt/TiO₂-HA composites. For comparison, large plasmonic Pt colloidal NPs with a mean diameter of *ca.* 73 or 107 nm, which show maximum extinction wavelengths of 455 or 542 nm, respectively, were synthesized according to the procedure reported by Bigall *et al.*³⁵. Unlike the multi-step process described by Bigall *et al.*³⁵, we develop a one-step approach to achieve the composites of large plasmonic Pt NPs-decorated TiO₂-HA (Pt/TiO₂-HA). First, 5 nm Pt seeds were prepared using the procedure according to Bigall *et al.*³⁵. Briefly, an aqueous solution of H₂PtCl₆·6H₂O (36 mL, 0.2%) was added to 464 mL of boiling water. After one min, a solution of sodium citrate (1%) and citric acid (0.05%) was added to the above H₂PtCl₆ solution, and 0.5 min later a freshly prepared NaBH₄ solution (5.5 mL, 0.08%) containing sodium citrate (1%) and citric acid (0.05%) was rapidly injected into the above reaction solution. After reaction for 10 min, the Pt seed solution was cooled down to room temperature. To obtain the Pt/TiO₂-HA with a varied Pt loading (i.e., 1 wt%, 3 wt%, or 7 wt% Pt), a different amount of TiO₂-HA powder (1110, 370, or 158.6 mg) was dispersed in 27 mL of water in a three-necked flask (fitted with a magnetic stirrer, a reflux condenser, and a thermometer) under vigorous stirring for 30 min. Afterwards, small 5 nm Pt seeds in 0.25 (for the Pt107/TiO₂-HA) or 0.60 mL (for the Pt73/TiO₂-HA) of water was added to the TiO₂-HA suspension and fully mixed with each other under vigorous stirring for another 30 min. Next, to the above mixture, 0.45 mL of an aqueous H₂PtCl₆ solution (0.20 M) and 2.5 mL of an aqueous solution containing sodium citrate (1.0 wt%) and L-ascorbic acid (1.25 wt%) were added. The resultant reaction mixture was heated to the boiling point (~100 °C) at a rate of ~8 °C min⁻¹ and then refluxed at this temperature for 30 min. After the mixture was cooled to ambient temperature, the supernatant of the suspension became colorless and the pristine white TiO₂-HA powder turned into light grey or grey and precipitated at the bottom of the flask, indicating the complete attachment of Pt NPs onto the TiO₂-HA. Finally, the resulting Pt/TiO₂-HA precipitates were isolated by centrifugation and thoroughly washed with hot water (50 °C) three times to remove the adsorbed organic species and then were dried in a vacuum oven at 80 °C for 12 h. Note that Pt NPs show a strong adhesion to the surface of TiO₂-HA with ample spacings among b-NWs and very high surface energy so as to allow for the direct contact between Pt NPs and the TiO₂-HA surface after cleaning with water, thus enabling a higher photocatalytic activity. The final Pt/TiO₂-HA samples were further treated by oxygen plasma etching to completely remove the organic residues.

Characterization of materials. Scanning electron microscopy (SEM) was performed using a Hitachi S-4800 field-emission scanning electron microscope operating at 5 kV to investigate the morphology and nano/micro-structure of the samples. Transmission electron microscopy (TEM) micrographs were obtained using a FEI Tecnai G2 Spirit Bio TWIN transmission electron microscope operating at an accelerating voltage of 120 kV. Specimens for TEM observations were sonicated before dropping them onto 300 mesh carbon-coated copper grids. High resolution TEM (HRTEM) micrographs, high-angle annular dark field (HAADF) scanning transmission electron microscopy (STEM) images, and energy-dispersive X-ray spectroscopy (EDS) elemental maps were acquired using a FEI Tecnai G2 F20 S-Twin electron microscope operating at 200 kV. X-ray photoelectron spectroscopy (XPS) measurements were carried out using a PHI5000 VersaProbe (ULVAC-PHI) spectrometer with an energy analyzer, employing a monochromatized microfocused Al K α ($h\nu = 1,486.58$ eV) X-ray source.

Samples for XPS measurements were pretreated by repeated cycles of Ar⁺ ion sputtering to obtain clean sample surfaces. The binding energies (BEs) of the core levels were calibrated by setting the adventitious C 1s peak at 284.8 eV. Survey spectra of the samples in the BE range of 0–1,000 eV, and the core level spectra of the elemental signals were recorded at resolutions of 1 and 0.125 eV, respectively. The X-ray diffraction (XRD) patterns were recorded using a Rigaku SmartLab diffractometer with Cu K α radiation ($\lambda = 1.5406 \text{ \AA}$) operating at 40 kV and 100 mA at a scanning rate of $0.06^\circ \cdot \text{s}^{-1}$. N₂ adsorption-desorption isotherm analysis was conducted at 77 K using a BELSORP-max micropore analyzer. The Brunauer–Emmett–Teller (BET) specific surface area (SSA) and the pore size distribution (PSD) of the TiO₂-HA sample were obtained based on N₂ adsorption isotherms in the relative pressure (P/P_0) range from 0.04 to 0.50 and non-local density functional theory (NLDFT) calculations by using nitrogen adsorption data and assuming a slit pore model, respectively. The samples were degassed under high vacuum ($<0.01 \text{ mbar}$) at 200 °C for at least 6 h prior to the measurements. UV-vis extinction spectra were recorded using a Shimadzu UV-3600 UV-vis-NIR spectrophotometer equipped with a LISR-3100 150 mm integrating sphere. The diffuse reflection spectra of all the samples were obtained using BaSO₄ as a standard reference in the measurements.

Photocatalytic measurements. To assess the photocatalysis performance of various Pt/TiO₂-HA samples, hydrogen production from pure water was monitored in the presence or absence of ethanol. In the case of using ethanol as the model sacrificial agent, each Pt/TiO₂-HA sample (10 mg) was suspended in an aqueous solution of ethanol (100 mL, the volume ratio of water to ethanol is 4: 1) in a closed-gas circulation reactor. For overall water splitting, a 10 mg sample of Pt/TiO₂-HA powder was dispersed in 80 mL of pure water. The reactant suspension was evacuated under vacuum three times and then an argon flow (12.5 mL min^{-1}) was introduced into the reaction system to completely remove air from the reactor. Concomitantly, the argon flow also served as the carrier gas to carry the reaction products to the detector. Afterwards, the suspension was irradiated using a 300 W xenon lamp (HSX-UV300) equipped with various cut-off filters (UVCUT400, L-42, Y-48, and O-54) for irradiating the photocatalysts with different illumination wavelengths. The amounts of gaseous products (H₂) was analyzed by a gas chromatograph ((GC, Shimadzu, GC-8A) using a thermal conductivity detector (TCD).

Finite-difference time-domain (FDTD) analysis. The enhancement of electric field at the interface of Pt107/TiO₂-HA and Au_{50 nm}/TiO₂-HA was calculated by using a software package, FDTD Solutions 8.15 (Lumerical Solutions, Inc.). During simulations, an electromagnetic pulse in the wavelength range from 400 to 700 nm (for Pt107/TiO₂-HA) or 450 to 700 nm (for Au_{50 nm}/TiO₂-HA) was launched into a box containing a target nanostructure. A size of $1 \times 1 \times 1 \text{ nm}^3$ was chosen for the override mesh cell. The model was set up by using a Pt nanosphere of 107 nm in diameter supported on a TiO₂ nanoring of 250 nm in outer diameter and 100 nm in inner diameter. The optical constants of Pt and Au were adopted from tabulated values of bulk platinum and gold (Palik), respectively, while that of TiO₂ (anatase) as adopted from Hand book of Optical Constants of Solids⁵³.

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Author Contributions

Y.T. designed the project and wrote the manuscript. L.Q. performed the experiments. Y.T. and G.W. conducted the numerical simulations. All the authors discussed the results and commented on the manuscript.

Additional Information

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