SCIENTIFIC REPORTS

Received: 12 July 2017 Accepted: 15 August 2017 Published online: 11 September 2017

OPEN A Chinese Pane-Like 2D Metal-**Organic Framework Showing Magnetic Relaxation and** Luminescence Dual-Functions

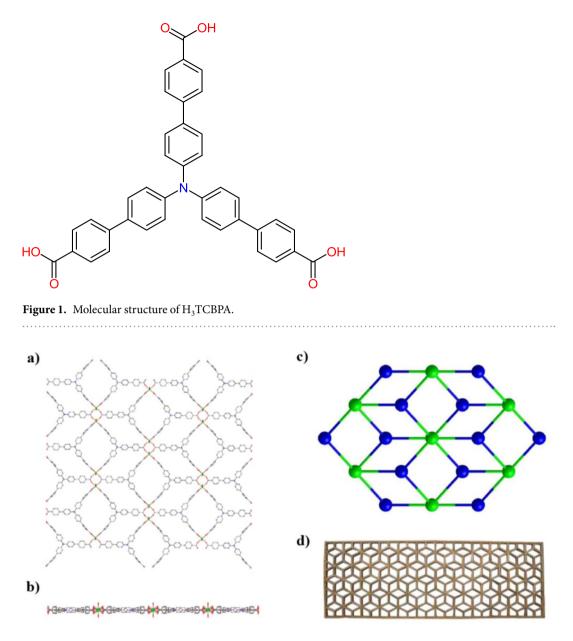
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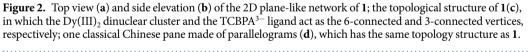
The discovery of graphene kicked off the curtain of atom-type two-dimensional (2D) materials. Layered metal-organic frameworks (MOFs) as parallel molecule-based 2D materials are more designable and more diverse, and magnetism may be induced by their metal ion nodes. However, the multifunctional 2D planelike MOFs are very difficult to obtain. Here we describe a Chinese pane-like 2D MOF constructed from the Ln³⁺ cation and the nanoscale luminescent tritopic ligand tris(4'-carboxybiphenyl)-amine, responding to the slow magnetic relaxation and luminescence properties, respectively. Notably, the Dy-Dy distances separated by the tritopic ligand are up to 2 nm. Such a 2D molecular material is expected to have potential applications in optoelectronics and multimodal sensing.

Two-dimensional (2D) materials have become most exciting and popular multifunctional materials since the discovery of graphene in 2004¹. Functional layered metal-organic frameworks (MOFs) can be viewed as molecule-based 2D materials²⁻⁴, whose pronounced advantages are more designable and more diverse, and their metal ion knots may induce good magnetism. Nevertheless, it is an extreme challenge to design the MOF-type 2D materials with multifunctions similar to the atom-type 2D materials. Recently, lanthanide(III) ions, especially the Dy(III) ion, which contain unique 4f orbital electrons showing strong spin-orbit coupling and high magnetic moment, have been chosen as the metal nodes to construct lanthanide-based MOFs (LnMOFs), endowing single-molecule magnet (SMM) behaviors^{5–7}. The SMMs feature magnetic hysteresis and slow magnetic relaxa-tion at low temperatures, with potential applications in magnetic devices for high-density information storage, quantum computing and spintronics⁸⁻¹⁴. The lanthanide (III) ions, each of which is equivalent to a single-ion magnet due to very weak magnetic exchange interactions between the lanthanide (III) ions containing shielded 4f electrons, may be arranged into highly ordered nodes with nanoscale-separation in the LnMOFs; such materials are convenient for assembly of corresponding molecular magnetic devices. Notably, the 2D layered LnMOF systems behaving as SMMs are especially attractive for this objective¹⁵.

Although some LnMOFs themselves may display photoluminescence by the aid of 'antenna effect' of ligands¹⁶, not all of LnMOFs are luminescent because the photoluminescence of lanthanide ions is subject to the energy's match rule. Alternatively, using a fluorescent spacer ligand to assemble LnMOFs should be an effective approach to the optical-magnetic multifunctional materials, this strategy allows both the SMM and luminescence properties from different building blocks to be combined together within a MOF system, guaranteeing both superparamagnet and emission functions. The motivation to explore the emissive quantum magnets is that they are potentially applicable in optoelectronics and multimodal sensing¹⁷. However, the reported luminescent SMMs are limited, which concentrate on mononuclear¹⁷⁻²¹, polynuclear²²⁻²⁶, and chainlike systems²⁷, the development of high-dimensional luminescent Ln(III) complexes showing SMM properties is still at a primitive stage²⁸; notably, the photoluminescence in these molecular materials mainly originates from the lanthanide(III) ions, the case with fluorescent bridging ligands has not been explored. Recently, we have adopted fluorescent tri s(4'-carboxybiphenyl)-amine²⁹ (H₃TCBPA, Fig. 1) as the spacer to construct multifunctional LnMOF-type 2D materials. Herein, we describe the solvothermal synthesis, crystal structure, luminescent and magnetic properties of a 2D LnMOF, $[Dy(TCBPA)(H_2O)_2]_n$ guest (guest: 2DMF·4H₂O, 1). Remarkably, 1 possesses an interesting

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plane-like network structure similar to one classical Chinese pane made of parallelograms, exhibiting not only slow magnetic relaxation responded to the Ln^{3+} knot but also ligand-based luminescence. Its erbium(III) analogue [Er(TCBPA)(H₂O)₂], guest (guest: 2DMF-4H₂O, **2**) is reported too.

Results and Discussion

Preparation. Pale yellow crystals of 1 or 2 were obtained as the only crystalline phase through solvothermal reaction between tris(4'-carboxybiphenyl)amine and $Ln(NO_3)_3 \cdot 5H_2O$ in an DMF/H₂O solution in the presence of 2-fluorobenzoic acid, which just serves as a directing agent and/or a modulator³⁰. It is noteworthy that this and other nanoscale tritopic ligands are inclined to construct 3D LnMOFs, acting as the linkers^{31–36}, whereas we successfully obtained the 2D layered LnMOFs with nanoscale tritopic ligands for the first time.

Structural description. Complex 1 crystallizes in the monoclinic space group $C_{2/m}$, its structure exhibits a 2D plane-like network, which is constructed from Dy(III)₂ dinuclear cluster nodes, TCBPA³⁻ bridging ligands and water terminal ligands (Fig. 2). The Dy atom is coordinated by two water O atoms and six carboxylate O atoms from four TCBPA³⁻ ligands. This [DyO₈] eight-coordination geometry was ascertained to be the biaugmented trigonal prism by the Shape software³⁷, and the calculation indicates that the deviation value from this polyhedron's ideal C_{2v} symmetry is 3.321 (Table S1, SI). Each TCBPA³⁻ ligand bridges four Dy³⁺ ions *via* three carboxylate groups (Fig. 2a):



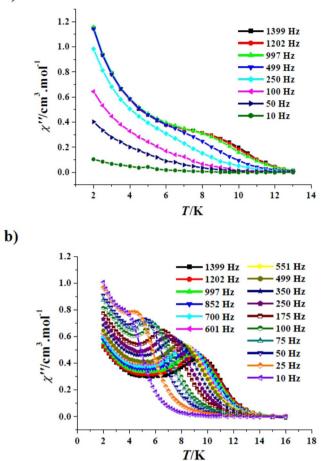


Figure 3. Plots of χ'' vs T for 1 ($H_{dc} = 0$ Oe, $H_{ac} = 2.5$ Oe) (**a**) and ($H_{dc} = 1$ kOe, $H_{ac} = 2.5$ Oe) (**b**).

two of which coordinate to two discrete Dy atoms with the chelating mode; whereas the third one links to two neighbouring Dy atoms with the *syn–syn* bridging mode, and the bridged Dy-Dy separation is 5.155 Å. Notably, the Dy-Dy distances separated by the TCBPA^{3–} bridging ligand are 21.555 Å, 21.555 Å and 19.641 Å,respectively. From a topological view, the Dy(III)₂ dinuclear cluster can be considered as a 6-connected node and the TCBPA^{3–} bridging ligand as a 3-connected node, therefore a simplified (3, 6)-connected network is generated (Fig. 2c), whose long topological (O'Keeffe) vertex symbol is 4.4.4.4.4.6.6.6.6.6.6 and the short (Schläfli) vertex symbol 4⁶.6⁶.8³, according to the analytical results by the Olex program. Such an interesting plane topology is reminiscent of one classical Chinese pane made of parallelograms (Fig. 2d). Complex **2** is of the same structure as **1**, but using Er instead of Dy (Fig. S1, SI). The average Dy-O bond distance of **1** is 2.357 Å (Table S2, SI), obviously larger than the mean Er-O bond distance of **2** (2.338 Å), being ascribed to the lanthanide contraction effect. Clearly, such a plane-like structure is totally different from those of the 3D LnMOFs based on the same ligand³¹.

Magnetic properties. The temperature dependence of magnetic susceptibility for complex 1 was investigated under 1 kOe in the range of 2–300 K. The room temperature χT value of 14.15 cm³ K mol⁻¹ for 1 is in agreement with the value of 14.17 cm³ K mol⁻¹ for one isolated Dy³⁺ ion in the ${}^{6}H_{15/2}$ state (Fig. S2, SI). The χT product declines gently with a temperature decrease in the range of 300-50 K, then falls steeply to reach 10.60 cm³ K mol⁻¹ at 2 K, which mainly originates from thermal depopulation of the excited Stark sublevels for the Dy³⁺ ions. The Curie-Weiss equation was used to fit magnetic susceptibility data, giving the C value of 14.27 cm³ K mol^{-1} and the θ value of -1.50 K; this very small negative θ value suggests that the possible dipole-dipole antiferromagnetic interactions between the Dy³⁺ ions is negligible. The field-dependent magnetization measured at 2-6 K revealed that there exists magnetic anisotropy in 1 as the M vs H/T plots are non-superimposed (Fig. S3, SI). Therefore, alternating-current (ac) susceptibility measurements as a function of temperature were performed to explore the dynamics of the magnetization of **1**. Besides the frequency dependence of the out-of-phase (χ') (Fig. S4, SI), the temperature dependence of the out-of-phase (χ'') determined in zero dc field shows significant signals below 12 K (Fig. 3a), indicating existence of SMM behaviours. However, no peaks could be detected owing to the obvious quantum-tunnelling effects. A dc field generally may suppress such quantum-tunnelling effects³⁸⁻⁴⁰. The field-dependent χ'' measured at 250 Hz and 7 K (Fig. S5, SI) indicates that the χ'' signals are enhanced with increased fields, but the change is small when the dc field is larger than 1 kOe. Therefore, a 1 kOe

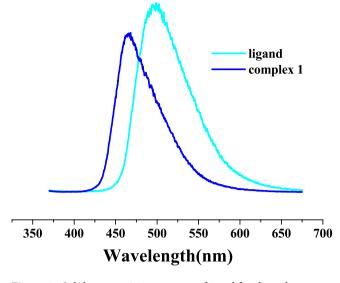


Figure 4. Solid-state emission spectra of 1 and free ligand at room temperature ($\lambda_{ex} = 350 \text{ nm}$).

dc field was applied to overcome the quantum-tunnelling effects, and not only χ' (Fig. S6, SI) but also χ'' (Fig. 3b)

peaks actually appeared at frequencies more than 25 Hz. In order to assess the effective barrier for the magnetic relaxation process of 1, the frequency-dependent χ'' peak temperature, in the format of the $\ln(\tau)$ vs 1/T plot, was fitted to the Arrhenius law, $\tau = \tau_0 \exp(U_{\text{eff}}/kT)$, affording U_{eff}/k value of 70.1(1.1) K and τ_0 value of 6.6(0.1) × 10⁻⁸ s (Fig. S7, SI). The τ_0 value is normal for the SMMs/SIMs (10⁻⁵- 10^{-11} s). The U_{eff}/k value of 1 is at a moderate level of carboxylate containing complexes⁴¹. The SMM behaviours of 1 were also supported by the parameter $\Phi \{= (\Delta T_t/T_t)/\Delta(\log f)\}$ calculation⁴², where T_t is the χ'' peak temperature, and *f* the frequency; the Φ value of 0.29 for **1** is actually in line with the superparamagnet value ($\overline{\Phi}$ > 0.1), but much larger than the spin glass value ($\Phi \approx 0.01$)⁴². Furthermore, under a dc field of 1 kOe, the frequency-dependent χ'' signals of 1 display well-defined temperature-dependent peaks at 4-9 K, being indicative of slow magnetic relaxation too (Fig. S8, SI). At 7 K and 8 K, two classical semicircle Cole-Cole diagrams were observed ($\chi'' vs \chi'$, Fig. S9, SI), which could be fitted by a Debye model for a single relaxation time (τ)^{43,44}, giving α values of 0.096 and 0.066 for 7 and 8 K, respectively. The small lpha values suggest narrow distributions of relaxation times. Furthermore, no loop was observed for 1 at 1.9 K (Fig. S10, SI). Based on the very weak magnetic interactions among the Dy³⁺ cations and the low symmetry of the Dy³⁺ cation, we assume that a Ising ground state of the Dy³⁺ cation exists in 1, the electrostatic model⁴⁵ was therefore used to estimate the magnetic axis of the dysprosium(III) ion. As shown in Fig. S11 (SI), the calculated magnetic axis directionality is very close to the line defined by Dy1-C13, with a small angle of 2.48°.

The magnetic relaxation of complex **2** was also studied for comparison. The χT value of 14.46 cm³ K mol⁻¹ at room temperature for **2** is concordant with the theory value of 11.48 cm³ K mol⁻¹ for one Er³⁺ ion (⁶I_{15/2}, Fig. S12, SI). Unlike complex **1**, the erbium(III) analogue **2** didn't display any obvious χ'' signals when the dc field was zero (Fig. S13, SI). Although the 2 kOe dc field can let complex **2** present both frequency-dependent χ' signals (Fig. S14, SI) and frequency-dependent χ'' signals (Fig. S15, SI), no peaks could be observed in the χ'' curves even at 1399 Hz and at above 2.0 K, which suggest a very small effective energy barrier value that could not be calculated in the above mentioned way based on the ln(τ) vs 1/T plot. This result indicates that the biaugmented trigonal prism of the [ErO₈] coordination geometry is not beneficial to present SMM behaviors⁴⁶.

Luminescence. The typical narrow luminescence bands arising from the Ln³⁺ cation are absent in the solid state photoluminescence spectrum of 1 at room temperature (Fig. 4). Nevertheless, upon excitation with 350 nm UV light, complex 1 displays a broad emission band at 465 nm, which is similar to the large emission band at 498 nm for the free H₃L ligand due to the π - π * transition, but a 33 nm of blue-shift was observed. Obviously, the ligand's coordination to the metal ion induces the luminescence's blue-shift³¹. The average fluorescence lifetime of 1 was determined to be 0.79 ns ($\lambda_{ex} = 360$ nm, Fig. S16, SI), shorter than 2.06 ns ($\lambda_{ex} = 360$ nm, Fig. S17, SI) of the free ligand. Furthermore, complex 1 shows a smaller absolute emission quantum yield (2.69%) than the free ligand (24.32%) ($\lambda_{ex} = 360$ nm), which can be mainly attributed to the solvent effect because 1 contains coordinated water molecules and some solvent molecules within the crystalline host lattices.

Similar to 1, only ligand-based luminescence spectrum could be observed for 2 in the solid state at room temperature. A broad emission band at 462 nm was excited by 360 nm UV light, which has a 36 nm of blue-shift comparing to free H_3L ligand because of the ligand to metal coordination role again (Fig. S18, SI). The absolute emission quantum yield (1.11%) of 2 is smaller than that of 1, and the fluorescence lifetime of 2 is too fast (<0.7 ns) to run lifetime measurement with the Edinburgh Analytical Instruments FLS980.

Conclusions

In summary, we present the first 2D layered LnMOF based on nanoscale tritopic ligands, which possesses an interesting plane-like structure and shows both slow magnetic relaxation and luminescence properties. We demonstrate a promising approach to the multifunctional MOFs using the magnetic Ln^{3+} cation and the

luminescent bridging ligand as the node and the linker, respectively. Our work also opens up a designable avenue to novel MOF-type 2D materials with interesting optical/magnetic multifunctions, such molecule-based 2D materials are expected to have potential applications in many fields such as optoelectronics and multimodal sensing.

Methods

Materials and instrumentation. Tris(4'-carboxybiphenyl)amine was synthesized following reported procedures²⁹. Other chemicals and solvents were obtained commercially and used as received. The elemental analyses (C, H, N) were accomplished on a Vario ELIII elemental analyser. The FT-IR spectra were determined on a Bruker/Tensor-27 spectrophotometer with pressed KBr pellets in the range 4000–400 cm⁻¹. The fluorescence spectra, the fluorescence lifetime and the absolute emission quantum yield were measured on an Edinburgh Analytical Instruments FLS980. Both the dc and ac magnetic susceptibility measurements were carried out on a Quantum Design MPMS-XL5 SQUID magnetometer using Pascal's constants diamagnetic corrections.

Preparation of 1 and 2. A mixture of tris(4'-carboxybiphenyl)amine (0.05 mmol), 2-fluorobenzoic acid (1.0 mmol), $Ln(NO_3)_3 \cdot 5H_2O$ (0.05 mmol), DMF (2.5 mL) and H_2O (0.75 mL) in a 25 mL Teflonlined stainless steel autoclave was maintained at 105 °C for 3 days. After the autoclave had been cooled slowly to 20 °C during 10 hours, light yellow plate crystals of 1 (Ln = Dy, 65% yield based on Dy), or light yellow plate crystals of 2 (Ln = Er, 58% yield based on Er) were harvested. These crystals were washed with water and dried at ambient temperature in air. Anal. Calcd (%) for $C_{45}H_{50}DyN_3O_{14}$ (1): C 53.02; H 4.94; N 4.12. Found: C 53.06; H 4.97; N 4.09. IR (KBr, cm⁻¹): 3443(b, s), 3033(w), 2928(w), 1663(m), 1600(m), 1522(m), 1492(w), 1424(s), 1324(w), 1280(w), 1188(w), 1112(w), 1088(w), 862(w), 835(w), 788(m), 731(w), 708(w), 685(w), 656(w), 576(w), 558(w), 475(w), 429(w). Anal. Calcd (%) for $C_{45}H_{50}ErN_3O_{14}$ (2): C 52.77; H 4.92; N 4.10. Found: C 52.73; H 4.95; N 4.08. IR (KBr, cm⁻¹): 3401(b, s), 3033(w), 2932(w), 1666(m), 1599(m), 1523(m), 1493(w), 1426(s), 1324(w), 1282(w), 1188(w), 1111(w), 1089(w), 864(w), 835(w), 788(m), 732(w), 709(w), 685(w), 656(w), 577(w), 560(w), 477(w), 429(w).

X-ray collection and structure determination. A single crystal with dimensions $0.14 \times 0.11 \times 0.06$ mm³ of **1** or $0.26 \times 0.09 \times 0.06$ mm³ of **2** was chosen to collect data on a Bruker SMART APEX-CCD diffractometer with Mo-K_a radiation ($\lambda = 0.71073$ Å) at 173(2) K. Empirical absorption corrections from φ scan were applied. Cell parameters were obtained by the global refinement of the positions of all collected reflections for two complexes. The structures were solved by direct methods and refined by a full matrix least-squares technique based on F^2 using with the SHELX-2014 program package. The DMF and H₂O solvent molecules in the complexes are highly disordered. The SQUEEZE subroutine of the PLATON software suit was used to remove the scattering from the highly disordered guest molecules. All non-hydrogen atoms were refined anisotropically, and all hydrogen atoms except those in coordinated water molecules were refined as riding atoms. Crystallographic data and structure determination parameters of complexes **1** and **2** are given in Table S3. CCDC-1557227 (**1**) and 1557228 (**2**) contain the supplementary crystallographic data, which can be obtained free of charge from The Cambridge Crystallographic Data Centre via www.ccdc.cam.ac.uk/data_request/cif.

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Acknowledgements

This work was supported by the National Key Basic Re-search Program of China (2013CB933403), the National Natural Science Foundation of China (21471154 and 91022014), and the Strategic Priority Research Program of the Chinese Academy of Sciences (XDB12010103).

Author Contributions

C.-M. Liu and D.-Q. Zhang designed the concept. C.-M. Liu wrote the paper. C.-M. Liu, D.-Q. Zhang, X.-Hao and D.-B. Zhu carried out the experimental work, and analyzed the results and revised the paper.

Additional Information

Supplementary information accompanies this paper at doi:10.1038/s41598-017-11006-5

Competing Interests: The authors declare that they have no competing interests.

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