# ARTICLE OPEN Check for updates Pit-to-crack mechanisms of 316LN stainless steel reinforcement in alkaline solution influenced by strain induced martensite

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The pit-to-crack transition of AISI 316LN stainless steel reinforcement exposed to stress corrosion cracking (SCC) in chlorides contaminated alkaline environment, was studied by a combination of slow strain rate testing (SSRT) and electrochemical impedance spectroscopy (EIS). The phase angle shift ( $\Delta \varphi$ ) obtained by EIS at low frequencies was utilized to determine the pit-to-crack transition, differentiating from crack nucleation and propagation as identified by shifts in the frequency range of phase angle ( $\theta$ ) peaks. The pit-to-crack transition was developed once the maximum  $\theta$  value shifted from the low to high frequencies. EIS analysis was corroborated by assessment of repassivation rates and pit growth, in addition to calculating  $\Delta G^{v \to \alpha'}$ . Crack nucleation at lath martensite developed transgranular SCC. Strain-induced martensitic transformation was associated with the brittle failure of AISI 316LN stainless steel, where  $\alpha'$ -martensite phase preferentially incubated the pit, and favored crack nucleation, thus promoting pit-to-crack transition.

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### INTRODUCTION

Stress corrosion cracking (SCC) occurs in the presence of a corrosive environment and mechanical stimulus, leading to the development of cracks that can propagate and reduce service lifetimes. SCC can develop without significant signs of damage accumulation, and is therefore critical threat for engineering asset management<sup>1,2</sup>. Austenitic stainless steels, while used in many industries for outstanding corrosion resistance, are not immune to SCC, even under atmospheric conditions and ambient temperatures<sup>3</sup>. One of the most common causes of SCC for austenitic stainless steel is chloride-induced stress corrosion cracking (CI-SCC), that develops as transgranular SCC (TG-SCC); which is triggered when a critical chloride threshold is reached in the presence of residual and/or applied loading. Several mechanisms have been proposed to describe CI<sup>-</sup>-SCC phenomena, including the slip-dissolution model by Newman, where both crack nucleation and propagation are promoted by the dissolution of slip planes, as well as surface film cracking on newly developed fresh planes<sup>4,5</sup>. Additionally, the cleavage-dissolution model proposes that crack nucleation and propagation stages of SCC are governed by the synergistic effect of brittle rupture of lowsurface-energy cleavage planes in conjunction with anodic dissolution<sup>6</sup>.

In marine and coastal environments, where aerosols containing high chloride concentrations are present, environmentally assisted cracking (EAC) develops due to thin-film electrolyte deposition or water droplets<sup>7–9</sup>. As a result of the higher chloride concentrations, the redox couple  $Fe^{2+}/Fe^{3+}$  resulting from iron acid hydrolysis locally reduces the pH below 2, enhancing the cathodic hydrogen evolution and promoting hydrogen embrittlement (HE)<sup>10</sup>. By concentrated sat droplet experiments under tension, it has been proven that the crack growth was related to hydrogen,

as seen with AISI 304 U-bends stainless steel samples exposed to MgCl<sub>2</sub> droplets at 30 °C and 28% relative humidity<sup>11,12</sup>. The HE involves the chemisorption of hydrogen at active sites with high stress, like the crack tip and the crack front. The HE mechanisms that better explain this behavior are hydrogen enhanced decohesion (HEDE), adsorption-induced dislocation emission (AIDE), hydrogen-enhanced localized plasticity (HELP), or a combination of some of them due to the complexity of the processes<sup>13,14</sup>. When hydrogen is present in austenitic stainless steel, the velocity of hydrogen migration through the microstructure was found to be higher than the crack propagation, thus allowing enough time for the hydrogen to diffuse<sup>10</sup>. The results of the HE are plastic deformation ahead of the tip of the crack developing the formation of micro-voids (HELP) or, the weakening of interatomic bonds facilitating decohesion (HEDE), both lowering the minimum energy required for the crack advance<sup>15,16</sup>.

In addition to the EAC aspects playing a role on the SCC mechanisms, austenitic stainless steel can develop microstructural phase transformation during the manufacturing processes or from in-service straining. The strain-induced phase transformation developed solid-state transformation of the austenitic phase (yphase) to martensite (either  $\epsilon$ - and/or  $\alpha'$ -phase)<sup>17</sup>. The presence of strain-induced martensite lowers the pitting potential and can act as SCC initiation sites, as well as promote crack propagation, thus increasing susceptibility to Cl<sup>-</sup>–SCC<sup>18</sup>. The anodic dissolution model of Newman et al.<sup>19</sup> proposed that the crack propagation rate increased when the strain-induced martensite was fractured at grain boundaries ahead of the crack tip. The SCC susceptibility of strain-induced martensite was shown by both electrochemical and optical microscopy characterization in chloride containing solution to lead to a faster crack nucleation, where preferential pitting attack on strongly slip banded austenite and  $\alpha'$ -martensite

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Fig. 1 Cyclic potentiodynamic polarization curves of AISI 316LN SS core and corrugate sections immersed in concrete pore solution. (a) 0 wt.% Cl<sup>−</sup>, (b) 4 wt.% Cl<sup>−</sup>, and (c) 8 wt.% Cl<sup>−</sup>.

Table 1. Electrochemical parameters obtained from the CPP curve for AISI 316LN reinforcement in simulated concrete pore solution at different Cl<sup>-</sup> concentrations without any applied stress.

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Sample	[CI <sup>-</sup> ]	$E_{\rm corr}~{\rm mV}_{\rm SCE}$	i <sub>corr</sub> A/cm <sup>2</sup>	$E_{\rm pit}~{\rm mV}_{\rm SCE}$	$\beta_{\rm a}$ mV/dec	$\beta_{\rm c}$ mV/dec	B mV	PSF
Core	0 wt.%	-226	$2.12 \times 10^{-9}$	_	276	445	74	_
Corrugated		-281	1.14 × 10 <sup>-9</sup>	_	212	245	49	_
Core	4 wt.%	-245	3.67 × 10 <sup>-7</sup>	_	351	375	79	_
Corrugated		-208	8.93 × 10 <sup>-8</sup>	651	325	215	56	0.89
Core	8 wt.%	-250	6.80 × 10 <sup>-7</sup>	600	485	456	102	0.97
Corrugated		-250	$8.75 \times 10^{-7}$	465	314	131	40	1.13

was observed<sup>20</sup>. The work from Ghosh et al. correlated the effect of strain-induced martensitic transformation to the CI-SCC of AISI 304 L at room temperature by U-bend testing<sup>21</sup>. The crack path started at the martensite formed in the outermost surface and followed through the slip bands. Even more, when the alloy had a low stacking fault energy (SFE), the likelihood of promoting straininduced martensite increased, thus being more prone to suffer from TG-SCC, lowering the chloride concentration threshold to nucleate the first pit<sup>22</sup>. Alyousif et al. studied the influence of SCC and HE in AISI 304 U-bended stainless steel samples in boiling MgCl<sub>2</sub> solutions, finding that the mechanism responsible for the SCC failure was a combination of strain-induced martensitic transformation and HELP<sup>23</sup>. The diffusion of hydrogen atoms ahead of the crack tip promoted a high plastic region, increasing the local stress by an interaction between  $\alpha'$ -martensite and the dislocation motion by HELP, prematurely reaching the stress for the crack advance. This was found to be a cyclic process, where the  $\alpha'$ -martensite accelerated both the crack nucleation and crack propagation.

However, there is still a lack of knowledge in the effect of straininduced martensitic transformation in SCC for high alkaline environments, and whether the strain-induced martensitic transformation affects only the rate and path for crack propagation, or if it also accelerates the pit initiation, which eventually would trigger the crack nucleation<sup>24</sup>. As seen on literature, martensite is more electrochemically active than austenite, becoming the most favorable anodic site, hence initiating the pit nucleation. Furthermore, the high chloride susceptibility of martensite combined with the more brittle phase, can make for a reduction of the minimum pit size for crack nucleation, lowering the stress susceptibility index ( $K_1$ ) to trigger the pit-to-crack transition<sup>25</sup>. In this work, the effect of the strain-induced martensitic transformation on the pit-to-crack transition of AISI 316LN reinforcing stainless steel in alkaline solution containing chlorides was studied using the slow strain rate technique (SSRT). Electrochemical monitoring was also utilized to elucidate the pit-to-crack transition by analysis of current transients and EIS measurements. The crack nucleation and propagation path were correlated with the strain-induced  $\alpha'$ -martensite and its influence on the electrochemical SCC mechanism.

#### **RESULTS AND DISCUSSION**

## Effect of the microstructure and surface on the electrochemical reaction kinetics

Figure 1 shows the cyclic potentiodynamic polarization (CPP) curves for the core and corrugated sections at the three different chloride concentrations. The corresponding corrosion potential  $(E_{corr})$ , corrosion current density  $(i_{corr})$ , pitting potential  $(E_{pit})$ , anodic ( $\beta_a$ ) and the cathodic ( $\beta_c$ ) Tafel slopes and, the Stern-Geary constant (B) are listed in Table 1. Up to 4 wt.% Cl<sup>-</sup> (chloride concentration below the critical chloride threshold of 4.9 wt.% Cl<sup>-</sup> for SS in concrete)<sup>26</sup>, the anodic kinetics for the corrugated samples had similar or lower current densities than the core for electrochemical potentials below the Cr<sub>2</sub>O<sub>3</sub> oxidation potential  $(Cr_2O_3 + 5H_2O \Rightarrow 2CrO_4^{2-} + 10H^+ + 6e^-, E_{Cr6+/Cr3+} = +225 \text{ mV}_{SCE}$ at 12.6 pH)<sup>20</sup>. However, as the anodic polarization increased, more electrons were supplied, activating the chloride induced pitting mechanism. While the passive oxide film is dissolved due to the higher rate of generation of oxygen vacancies compared to the rate of annihilation of oxygen vacancies supplied by the metal



Fig. 2 SEM micrograph of as-received AISI 316LN SS reinforcement. (a) core section 150x, and (b) corrugated section 150x.

surface, the *i* raised. At this point, the passive film thinning was triggered until passivity breakdown finally occurred<sup>27</sup>.

The difference in chloride susceptibility was seen by the development of the  $E_{pit}$  at 4 wt.% Cl<sup>-</sup> for the corrugate section (+698 mV<sub>SCE</sub>), while the core sample experienced a negative hysteresis, thus no  $E_{pit}$  develops. Furthermore, after the chloride concentration threshold was reached (8 wt.% Cl<sup>-</sup>), the core sample developed a  $E_{pit}$  at +600 mV<sub>SCE</sub>, while the corrugated section depicted a  $E_{pit}$  value of +465 mV<sub>SCE</sub>, thus indicating a higher pitting susceptibility of the corrugated section. The  $\beta_a$  values of the core sample revealed a lower electrochemical corrosion kinetics, in addition of the higher *B* values per chloride concentration of the core than the corrugated samples, indicating a more passive behavior for the core samples. To further compare the dissimilar pitting susceptibility properties between the core and the corrugated sections, the pitting susceptibility factor (PSF) was calculated (see Eq. (1))<sup>28</sup>:

$$\mathsf{PSF} = \frac{E_{\mathsf{pit}} - E_{\mathsf{rep}}}{E_{\mathsf{pit}} - E_{\mathsf{corr}}} \tag{1}$$

where the PSF is a probabilistic parameter that takes into account the  $E_{\rm corr}$ , the  $E_{\rm pit}$  and the repassivation potential ( $E_{\rm rep}$ ), which is the potential of the backward scan that exhibits a current density of 1  $\mu$ A/cm<sup>2</sup><sup>29</sup>.

The PSF values below 0.5 indicate low pitting susceptibility, while values equal or greater than 1 indicate high pitting susceptibility. The corrugated sample exposed to 4 and 8 wt.% Cl<sup>-</sup> had a PSF of 0.89 and 1.13 respectively, while the core sample only developed a PSF of 0.97 at 8 wt.% Cl<sup>-</sup>; showing the higher pitting susceptibility of the corrugated section compared to the core section.

To evaluate the effect the deformation mechanism had on the electrochemical kinetics, the microstructure of the core and corrugated sections was revealed. Figure 2a shows a  $\gamma$ -phase matrix with some minor deformation twins and slip bands characteristic of face centered cubic (FCC) systems. The corrugated section had a higher density of slip bands, and revealed the development of single and double slip crossover bands within grains (see Fig. 2b). The higher density of planar defects at the corrugated section reduced the mean free path of dislocation

glide, therefore increasing the dislocation pile ups and hindering the corrosion protection properties, which indicates the higher pitting susceptibility of the corrugated section over the core<sup>30</sup>. The enhanced anodic dissolution kinetics of the corrugated section can be explained by the mechanical energy stored inside the plastically deformed grains, which can be quantified by the residual stress<sup>31,32</sup>. By performing XRD analysis on the free residual stress peak (found at  $2\theta = 147^{\circ}$  for austenite), a tensile residual stress of 36 ± 12 MPa was obtained, indicating the higher mechanical energy of the surface compared to the core, which had no residual stress. As reported elsewhere, the higher the tensile residual stress the higher the pitting susceptibility, as well as the easier the pit-to-crack transition<sup>33,34</sup>. In addition, grain refinement was found due to manufacturing process, where the grain size values were  $13 \pm 2$  and  $10 \pm 3 \,\mu\text{m}$  for the core and corrugated sections, respectively.

## Effect of the strain-induced martensitic transformation on the pit-to-crack transition

The stress-strain curve including the current transient as a function of the time is presented in Fig. 3a, where the yield strength ( $\sigma_v$ ) and the ultimate tensile strength ( $\sigma_{UTS}$ ) are labeled with dashed lines<sup>35,36</sup>. Two regions can be distinguished, resilience from the preload to the  $\sigma_{\rm v}$  (elastic region, colored in blue) and toughness being from the  $\sigma_{\rm v}$ to failure (plastic region, colored in green)<sup>37</sup>. The initial electrochemical current density measurement with an applied stress circa 120 MPa experienced a sharp increase in *i* starting at  $66 \,\mu$ A/cm<sup>2</sup>. At this initial loading the nucleation of pits is metastable, as it has not yet reached the  $E_{pit}$  (with a corresponding  $i = 200 \,\mu\text{A/cm}^2$ ). Before the  $\sigma_v$  point (270 MPa), at 192 MPa the recorded *i* reached the 200 µÅ/cm<sup>2</sup>, which in combination with the induced localized microplastic deformation promoted a stable pit growth<sup>38</sup>. Once the stress approached the  $\sigma_{v}$ , the *i* rose to 430  $\mu$ A/cm<sup>2</sup>, remaining at this value with some minor fluctuations until the  $\sigma_{UTS}$ , where finally the *i* increased up to the failure. The pit-to-crack transition was calculated by the stress intensity factor to promote SCC ( $K_{ISCC}$ ), which for stainless steel is around 50 MPa  $\sqrt{m^{39}}$ . To calculate the K<sub>ISCC</sub> a hemispherical pit geometry was selected, yielding Eq.  $(2)^{40}$ :

$$K_{\rm ISCC} = 1.1215(K_{\rm t}\sigma)\sqrt{\pi r_{\rm pit}}$$

(2)



Fig. 3 Stress corrosion cracking testing curves for AISI 316LN SS reinforcement in simulated concrete pore solution with 8 wt.% CI<sup>-</sup>. (a) stress-current transient curves, resilience (blue) and, toughness (green), inset showing sketch for the current transient decomposition, and (b) SCC stages with growth ( $v_{\text{grow}}$ ) and repassivation ( $v_{\text{rep}}$ ) rates as a function of strain, showing the different stages for the pit and crack processes, and error bars representing the standard deviation.

where  $K_t$  is the stress concentration factor (3 for semicircle),  $r_{pit}$  is the radius of the pit, and  $\sigma$  is the normal stress.

The  $r_{pit}$  was calculated integrating the monitored current density and converting it to mass loss based on Eq. (3)<sup>41</sup>:

$$r_{\rm pit} = \sqrt[3]{\frac{3EW\int i\ dt}{2\pi\rho F}} \tag{3}$$

where *EW* is the equivalent weight, *i* is the current density, *t* is the time and  $\rho$  is the density.

By assuming the monitored electrochemical current corresponds to the governing anodic dissolution process (i.e., the growth of the dominant pit), the pit-to-crack transition was triggered at 295 MPa (~110%  $\sigma_{\rm v}$ ).

An in-depth analysis of the current transient values shown in Fig. 3a allowed for the pit growth and repassivation rates to be determined (see Fig. 3b). Following the work from Feng et al., each of the current transient peaks can be divided into growth of the pit and its repassivation, where the fitting of each region gives the corresponding rates of growth ( $v_{grow}$ ) and repassivation ( $v_{rep}$ ) (Eqs.

(4) and (5))<sup>42</sup>:

$$u_{\rm grow} = \frac{i_{\rm peak} - i_{\rm bg}}{(t_1 - t_0)^2}$$
 (4)

$$v_{\rm rep} = \frac{i_{\rm peak} - i_{\rm bg}}{(t_2 - t_1)^2}$$
(5)

where  $t_0$  is the start time of the pit,  $t_1$  is the time when maximum peak current occurs,  $t_2$  is the end time of the pit,  $i_{peak}$  is the maximum peak current and,  $i_{bg}$  is the background current (see Fig. 3a inset).

The  $v_{arow}$  always stayed above the  $v_{rep}$  during the entire straining, showing that the sample was active for the entire testing and that once the passivity breakdown occurred, no repassivation was possible. From the preload the  $v_{\text{grow}}$  grew and made a plateau at 0.1  $\mu$ A cm<sup>-2</sup> s<sup>-2</sup>, while the  $v_{rep}$  immediately decreased and plateaued around 0.03  $\mu$ A cm<sup>-2</sup> s<sup>-2</sup>. After  $E_{pit}$  value is reached, the formation of pits is developed (green region Fig. 3b), where the  $v_{\text{grow}}$  increased while the  $v_{rep}$  remained unchanged. Once the  $\sigma_y$  was reached, the activation of dislocations glide begun, and the strain-induced martensite transformation was triggered. By increasing the density of planar defects, which enhances the  $\alpha'$ -phase nucleation, the chloride ions will be easily adsorbed at those anodic sites, thus decreasing the surface energy facilitating the development of the first crack after the KISCC is reached (crack nucleation stage, yellow region). After the pit-to-crack transition completed, the crack propagation became the governing process and continued until failure (crack propagation stage, red region). For stresses close to the  $\sigma_{\rm UTS}$ , the  $v_{\rm rep}$  started decreasing while the  $v_{\rm grow}$  kept its increasing trend. This phenomenon is explained by the combination of increased stress and the chloride attack, which make the dissolution and cracking rate to be much greater than the repassivation<sup>43</sup>. This is in agreement with previous works where due to the increase in dislocation density the dissolution rate increases, depleting the  $v_{\rm rep}^{44}$ . After reaching  $\sigma_{\text{UTS}}$  the stresses have reached the maximum possible strain-induced martensitic transformation and necking of the sample develops, as observed by the sudden increase of  $v_{arow}$ , whereas  $v_{rep}$ is almost negligible. The higher anodic dissolution activity than repassivation at the crack tip allows for the crack propagation, which is influenced by the  $\alpha'$ -phase strain-induced martensite transformation and the acid hydrolysis of iron at the crack tip<sup>45</sup>.

To further prove the beginning of the strain-induced martensitic transformation at the  $\sigma_{y}$ , the calculations of the  $\Delta G_{\text{total}}^{\gamma - \alpha'}$  were made following the work from Patel and Cohen<sup>46</sup>. Assuming that the mechanical stimulus and the chemical thermodynamics governed by the Gibbs free energy are the main factors influencing the transformation, the following formula is obtained (see Eq. (6)):

$$\Delta G_{\text{total}}^{V \to \alpha'} = -V_m [210.9 - 0.732(M_s + T)] + V_m \left[\frac{\sigma}{2} \times \gamma_0 \times \sin 2\theta + \frac{\sigma}{2} \times \varepsilon_0 \times (1 + \cos 2\theta)\right]$$
(6)

where  $V_{\rm m}$  is the molar volume,  $M_{\rm s}$  is the temperature at which the athermal martensite transformation begins,  $\gamma_0$  is the shear strain,  $\theta$  is the angle of the applied stress to a potential habit plane normal, and  $\varepsilon_0$  is the volume expansion normal to the habit plane<sup>47</sup>.

With  $M_s$  being -258 °C (calculated based on the elemental composition<sup>48</sup>), *T* being 25 °C,  $\gamma_0$  being 0.229 and  $\varepsilon_0$  being 0.012 both based on variables calculated from the X-ray diffraction (XRD) pattern,  $V_m$  being 7.09 cm<sup>3</sup>/mol for iron and,  $\sigma$  ranging from the  $\sigma_y$  to the  $\sigma_{UTS}^{49}$ . Using Eq. (6), the  $\Delta G^{\gamma \to \alpha'}$  ranged from -2677 to -2738 J/mol. Following the work from Hauser et al., the minimum  $\Delta G$  value for the strain-induced martensite to form was calculated to be the  $\Delta G^{\gamma \to \alpha'}$  at 0 °C, which yielded – 2643 J/mol<sup>49</sup>. Thus, stress values around the  $\sigma_y$  show the strain-induced martensitic transformation is thermodynamically favored, and it is from the  $\sigma_y$  onward that the  $\alpha'$ -phase nucleates.



Fig. 4 EIS plots for AISI 316LN SS strained reinforcement in simulated concrete pore solution with 8 wt.% Cl<sup>-</sup>. (a) Nyquist diagrams with insert electrical equivalent circuit (EEC) with two time constants, (b) magnification for loading percentages above the preload, (c) Bode plot with the phase angle ( $\theta$ ), and (d) phase shift for low and high frequencies at every load percentage.

## Pit-to-crack transition analysis by electrochemical impedance spectroscopy

The EIS results are plotted in the Nyquist diagrams of Fig. 4a, b, where the experimental data was recorded for every percentage of the loading. The Bode plots with the change in the phase angle ( $-\theta$ ) were plotted in Fig. 4c. Before the EIS data was further analyzed, the Kramer-Kronig transforms were performed to show the robustness of the data, proving that there were no artifacts<sup>50</sup>. The electrical equivalent circuit (EEC) used to fit EIS data is shown in the insert of Fig. 4a, having a hierarchically distribution and, being comprised of two times constant in parallel. A good matching was found between EIS fitting data and the experimental values recorded, where the total error was below 10% and chi-squared ( $\chi^2$ ) values in the order of 10<sup>-4</sup>, proving the goodness of the fitting (see Table 2).

The pseudo-capacitance of the passive oxide film ( $Y_{\text{film}}$ ) progressively increased its value, suggesting a more doped passive film with high density of defects, hence more prone to corrosion – as the pseudo-capacitance values were in the 21–365 µS cm<sup>-2</sup> s<sup>nfilm</sup>. The  $n_{\text{film}}$  values were around 0.7, indicating a rough and defective film due to the *n* values not reaching the unity<sup>51</sup>. The charge transfer resistance ( $R_{\text{ct}}$ ), which accounts for the properties of the electrochemical double layer, started at 108 k $\Omega$  cm<sup>2</sup>, three orders of magnitude higher than resistance of the film ( $R_{\text{film}}$ )<sup>52</sup>. The high starting values of  $R_{\text{film}}$  accounted for the developed passive film at the open circuit potential (OCP) prior to

the electrode straining, and as the strain increased in combination with the chloride present in the electrolyte hinders the passive film growth due to the flow of vacancies towards the electrode/ electrolyte interface. As the process continued, the higher density of oxygen vacancies doping the passive film promoted the passivity breakdown, as indicated by the lower values of the  $R_{\rm ct}$ .

The selection of the *CPE* over the ideal capacitor was taken upon the non-ideality of the electrochemical interface, first because of the roughness of the corrugated area and also because of the polycrystalline grains forming the microstructure, all having different size and different orientation. As a consequence, the capacitance cannot be directly read as the *CPE* depicts a branched ladder *RC* network, therefore, the effective capacitance (*C*<sub>eff</sub>) for the passive film and the electrochemical double layer can be obtained (see Eq. (7) for *C*<sub>eff,flm</sub><sup>53</sup> and Eq. (8) for *C*<sub>eff,dl</sub><sup>54</sup>):

$$C_{\rm eff, film} = Y_{\rm film} \left( \omega_{\rm m}^{''} \right)^{n_{\rm film} - 1} \tag{7}$$

$$C_{\text{eff,dl}} = \left[ Y_{\text{dl}} \left( \frac{1}{R_{\text{s}}} + \frac{1}{R_{\text{ct}}} \right)^{(n_{\text{dl}}-1)} \right]^{\frac{1}{n_{\text{dl}}}}$$
(8)

The  $C_{\text{eff,film}}$  started at values above the typical  $\mu$ F cm<sup>-2</sup>, which is the order of magnitude expected for a protective passive film, instead it begun at 53  $\mu$ F cm<sup>-2</sup> and raised up to 2650  $\mu$ F cm<sup>-2</sup>.

Applied load	$R_{\rm s} \ \Omega \ {\rm cm}$	$R_{\rm film} \ \Omega \ {\rm cm^2}$	Y <sub>film</sub> µS cm <sup>−2</sup> s <sup>nfilm</sup>	n <sub>film</sub>	$C_{\rm eff, film}~\mu {\rm F}~{\rm cm}^{-2}$	$R_{\rm ct}~{ m k}\Omega~{ m cm}^2$	$Y_{\rm dl}~\mu S~{ m cm}^{-2}~{ m s}^{n { m dl}}$	n <sub>dl</sub>	$C_{\rm eff,dl}~\mu {\rm F}~{\rm cm}^{-2}$	$\chi^{2}$ (a)
Preload	3	870	21	0.80	53	108	38	0.88	11	$1.13 \times 10^{-4}$
80% σ <sub>y</sub>	3	611	50	0.70	203	24	253	0.76	25	$1.26 \times 10^{-4}$
115% σ <sub>y</sub>	3	465	61	0.69	254	17	308	0.74	26	$1.93 \times 10^{-4}$
90% σ <sub>UTS</sub>	3	452	71	0.69	297	13	431	0.72	36	$1.75 \times 10^{-4}$
100% σ <sub>υτs</sub>	3	431	77	0.68	337	11	457	0.73	40	$2.69 \times 10^{-4}$
Failure	3	373	365	0.66	1750	9	932	0.71	86	$2.55 \times 10^{-4}$



**Fig. 5** XRD patterns of AISI 316LN SS reinforcement, as-received (AR) and strained samples.

Such high C<sub>eff,film</sub> values imply that the passive film is extremely doped with point defects, hence the electron transfer from the electrolyte through the passive film requires less energy. Even more, those Ceff,film values, if converted to thickness, would give less than 1 nm, meaning that no passive film is present (as an oxide crystal structure would have a lattice parameter  $\approx 1$  nm). This coincides with the high pitting susceptibility that the corrugated section showed, as well with the high *i* values recorded during the current transient monitoring. The obtained C<sub>eff dl</sub> values show a defective electrochemical double layer, while being between one and two orders of magnitude lower than the C<sub>eff.film</sub>, which had high capacitance values in the order of tenths of µA/cm<sup>2</sup>. The electrochemical double layer, which structure can be explained by Helmholtz model, indicated that the band gap had considerably thinned, almost overlapping the conduction band  $(E_c)$  and the valence band  $(E_v)$ , meaning that the time constant of the passive film was no longer related to a semiconductor but a conductor. This depletion of the space charge is due to the applied potential which raised the OCP above the flatband potential  $(E_{\rm fb})$ , forcing the electrons to flow towards the electrolyte.

In addition, EIS data were analyzed to reveal the change of the –  $\theta$  as a function of the frequency was monitored for the different applied stress regions to further elucidate the pit-to-crack

transition, as it has been proven to have a dependency<sup>55-60</sup>. Figure 4d shows the evolution of the phase angle shift ( $\Delta \phi$ ) as a function of the applied load for the high and low frequencies (taken from Fig. 4c), where the  $\Delta \phi$  is the subtraction of the preload to the other applied load points<sup>61</sup>. From the preload to the 80%  $\sigma_{\rm v}$ the  $\Delta \phi$  for both low and high frequencies increased 20 and 10 degrees, respectively; indicating the passivity breakdown developed at similar stress values compared to the ones observed by the current transient monitoring<sup>56</sup>. From that point on, the  $\Delta \phi$  of high frequencies did not vary, suggesting that the low frequencies were the responsible for indicating the cracking process; in this case the range was from 0.1 to 1 Hz, in accordance with the literature<sup>62</sup>. The increase of the  $\Delta \phi$  after the  $\sigma_y$  for the low frequencies were influenced by the formation of the  $\alpha'$ -phase, enhancing the anodic kinetics as also seen from the CPP analysis, but it was not until the maximum value of  $\theta$  ( $\theta_{max}$ ) shifted from the low frequencies to the high frequencies, that the pit-to-crack transition developed. Similar cracking behavior is seen in literature, where the increase in  $\Delta \phi$  denoted the crack nucleation<sup>57,58</sup>.

The decrease of  $\theta_{max}$  indicates the crack nucleation and growth, showing the decrease in  $R_{ct}$  (as both  $\theta_{max}$  and  $R_{ct}$  are related)<sup>63</sup>. Furthermore, as the crack incubated, the frequency of the  $\theta_{max}$  ( $f(\theta_{max})$ ), shifted towards higher frequencies (change from low to high frequencies, see Fig. 4c)<sup>61</sup>. The crack propagation was also proof by the  $\Delta \phi$  towards lower frequencies of the low frequency peak (both  $\sigma_{UTS}$  points and Failure). Finally, after the increase of the  $\Delta \phi$  for low frequencies, the last point before the failure depicted the formation of a third peak at intermediate frequencies (1 Hz), as well as abruptly increasing both low and high  $\Delta \phi$ . The development of this third relaxation process was due to the large crack propagation and as a consequence, both the  $\theta_{max}$  and  $f(\theta_{max})$  shifted to this third peak (see Failure in Fig. 4c). This third phase angle peak denotes an active cracking process, as explained by Ghobadi et al.<sup>60</sup>.

## Effect of the strain-induced martensite on the fracture and crack propagation

Following failure of the AISI 316LN stainless steel reinforcement, XRD was conducted in the as-received (AR) sample as well as in the strained sample (close to the corrugated section) to compare the phase composition (see Fig. 5). The AR sample only had peaks corresponding to the  $\gamma$ -phase, indicating a FCC structure, while the strained sample developed two peaks for the  $\alpha'$ -phase having a body centered cubic (BCC) crystal structure<sup>64</sup>. The absent of  $\alpha'$ -phase in the AR sample could be due to the overlap of  $\alpha'$ -phase (110) with the peak of  $\gamma$ -phase (111) at  $2\theta = 43^{\circ}$  or due to a density of  $\alpha'$ -phase below 1 vol.% (limit definition for the equipment). The  $\alpha'$ -phase of the strained sample was elucidated and quantified by an approximate value of 13.4 vol.% by the integration of the intensity peaks of each phase<sup>65</sup>. The lattice parameter of the  $\gamma$ - ( $a_{\gamma}$ ) and  $\alpha'$ -phase ( $a_{\alpha'}$ ) were calculated,



Fig. 6 Surface SEM micrographs of AISI 316LN SS strained reinforcement in simulated concrete pore solution with 8 wt.% CI<sup>-</sup>. (a) fracture plane 100×, (b) cleavage plane with martensitic pack inside a pit 1,000×, (c) strain-induced martensitic packets 14,000×, (d) detail of cracked strain-induced martensitic packet 20,000×, and (e) transgranular crack 50×.

obtaining a value of 3.575 and 2.849 Å, respectively<sup>66</sup>. As it can be seen the *a* of the  $\gamma$ -phase is larger than the one from the  $\alpha'$ -phase due to the compression stresses that are required for the phase transformation, however the packing factor for the FCC is 74% while for the BCC is only 68%, indicating a higher void volume for the  $\alpha'$ -phase<sup>67</sup>. In addition, the calculated tetrahedral void radius of BCC crystal structure was larger than the one from the FCC, being 0.29  $r_{\rm Fe}$  compared with the 0.225  $r_{\rm Fe}$  ( $r_{\rm Fe}$  = radius of the Fe atom)<sup>68</sup>. As seen from the work of Zhu, the adsorption of chloride ions on the interstitials minimizes the surface energy of the cleavage planes, thus the  $\alpha'$ -phase with its higher void radius will facilitate the formation of cleavage planes (brittle fracture)<sup>6</sup>.

The surface morphology of the fractured sample can be seen in Fig. 6a, where the two sides of the strained sample are shown in the inset. The notch area is delimited by the dotted line, as well as the outer section of the corrugate surface. A brittle fracture mode is shown on the outer section, where the formation of cleavages planes along the corrugate surface are identified (see Fig. 6b). Moreover, those areas presented numerous pits, proving the high chloride susceptibility of the corrugated section. This agrees with the CPP analysis where the corrugated area had a much higher pitting susceptibility, with a PSF > 1 for 8 wt.% Cl<sup>-</sup>. In addition, this can be explained by the tensile residual stress obtained from the XRD analysis ( $36 \pm 12$  MPa), favoring the pit nucleation, and further crack nucleation<sup>33</sup>. The reason of the premature nucleation and fast pit propagation was the strain-induced martensite phase transformation, which acted as a preferential site for the pit nucleation; seen as martensitic packets inside the pits (see Fig. 6c). This more active and brittle martensite  $\alpha'$ -phase also enhanced the crack nucleation and propagation, seen by the dissolution of the martensitic packets (gray stripes), while no dissolution was found on the  $\gamma$ -phase matrix (white stripes) (see Fig. 6d). This finding agrees with previous studies of austenitic SS where the strain-induced martensite enhanced the crack propagation<sup>23</sup>. The grain morphology depicted in Fig. 6d resembled an austenitic grain with the formation of lath martensite, where the stripes shape had different directions within the same grain, due to the combination of complex stresses<sup>69</sup>. The cracking of the lath martensite is due to the angular misorientation at the interfaces, this crack propagation is also favored by the anodic dissolution kinetics, thus promoting the SCC mechanism<sup>70</sup>.

Figure 6e indicates a well-defined perpendicular crack to the surface, initiated on the outermost surface where the straininduced martensite was located (Fig. 6b), propagating towards the inside through the austenite grains following slip bands and further  $\alpha'$ -phase (TG–SCC morphology)<sup>18</sup>. The transgranular crack was measured by ImageJ software to propagate over 222 µm depth with an average width of 6.84 µm (based on Fig. 6e). This width was smaller than the average grain size of the  $\gamma$ -phase matrix (13 µm), indicating that the dissolution process did not completely dissolve the grains.

The inverse pole figure (IPF) map of the AR (Fig. 7a) and strained samples (Fig. 7d notch and Fig. 7g corrugate) indicates polygonal shape  $\gamma$ -phase grains with some aligned small local misorientation domains, which were attributed to the plastic strain forming dislocation channels where the strain-induce martensite would be promoted (see Fig. 7d, g)<sup>71</sup>. These misoriented domains formed specific chain-like formations with different contrast to the parent  $\gamma$ -phase matrix. To verify the nature of the phase transformation, the phase maps were also added, revealing that no hexagonal close-packed  $\epsilon$ -phase matrix was colored in green, and as it can be seen for the AR sample, no traces of  $\alpha'$ -phase (red) where seen (see Fig. 7b). In the case of the strained samples,  $\gamma$ -phase and  $\alpha'$ -phase where found, coinciding with the crystallographic analysis from the XRD (see Fig. 7e, h)<sup>72</sup>. The strain-induced



Fig. 7 EBSD analysis of AISI 316LN reinforcement; core area of as-received (AR) sample. (a) Inverse pole figure (IPF) mapping, (b) phase map, and (c) kernel average misorientation (KAM) mapping; notched area of strained sample (d) IPF mapping, (e) phase map, (f) KAM mapping, corrugate area of strained sample (g) IPF mapping, (h) phase map, and (i) KAM mapping.

martensitic transformation commenced at slip bands, and followed a parallel path along such bands, as observed seen in the misoriented domains (Fig. 7d, g)<sup>73</sup>. Kernel average misorientation (KAM) mapping was carried out to evaluate local changes in alloy structure; whilst for this study KAM was calculated using nearest neighbors with a maximum misorientation of 5°. As observed from Fig. 7c, the majority of the KAM map revealed low to no deformation (green to blue color), with some minor regions of greenish-yellow areas attributed to deformation twins, sites with gathered plastic energy. Nonetheless, for the strained samples, where the slip bands and  $\alpha'$ -phase domains where located higher KAM values were found, this is due to the dislocation and planar defects acting as obstacles for dislocation

motion (red color) (Fig. 7f, i). As a conclusion, by combining the higher strains observed on the  $\alpha'$ -phase by EBSD with its higher chloride pitting susceptibility seen by the electrochemical testing, it is to expect that SCC mechanism is governed by the  $\alpha'$ -phase martensite transformation. Which agrees with previous literature studies where strain-induce martensite was seen to be the main cause for SCC development on austenitic stainless steel<sup>74</sup>.

In summary, the study herein revealed that the SCC mechanisms responsible for the brittle fracture of AISI 316LN stainless steel rebars in alkaline environment was the strain-induced martensitic transformation, where  $\alpha'$ -phase preferentially incubated the pit, and promoted the pit-to-crack nucleation. The combination of the higher strains observed on the  $\alpha'$ -phase by EBSD with its higher chloride pitting susceptibility seen by the electrochemical testing, agrees with previous literature studies where strain-induced martensite was seen to be the main cause for SCC development on austenitic stainless steel<sup>74</sup>. It was also proved that stress values around the  $\sigma_y$  made the strain-induced martensitic transformation thermodynamically favored, and it is from the  $\sigma_y$  onward that the  $\alpha'$ -phase nucleated.

### METHODS

### Materials

The material used was AISI 316LN stainless steel rebar grade 75, the elemental composition of which can be seen in Table 3. AISI 316LN reinforcing stainless steel was manufactured by hot-rolling process at a temperature of 890 °C with a post processing of pickling and passivation as per the manufacturer specifications (NAS, North American Stainless). The reinforcing stainless steel specimens were so-called size 3 (3/8 inches, 9.5 mm diameter). The AISI 316LN reinforcement was cut and prepared following ISO 7539-6:201875. To accelerate the crack nucleation process the specimens were machined with a sharp 60° V-notch in the center of the 60 mm corrugated section (see Fig. 8)<sup>76</sup>. Before the samples were attached to the electrochemical cell for the SCC test, they were rinsed with deionized (DI) water, degreased with acetone and blown-dried with air. A 3  $cm^2$  exposed area was selected by coating with epoxy lacquer (Midas 335-009 non-conductive paint) (see Fig. 8b).

#### Slow strain rate testing

Reinforcement steel specimens were subjected to a uniaxial tensile test by slow strain rate testing (SSRT) while being immersed in a corrosive media following the ASTM-G129 standard<sup>77</sup>. The testing solution was simulated concrete pore solution (pH 12.6) made out of a saturated Ca(OH)<sub>2</sub> contaminated with 8 wt.% Cl<sup>-</sup> by means of CaCl<sub>2</sub> additions. The pH of the electrolyte solution (Ca(OH)<sub>2</sub> + 8 wt.% Cl<sup>-</sup>) was monitored before and after the testing, proving that the bulk pH did not vary during

Table 3.Elemental composition (wt.%) of AISI 316LN SSreinforcement, Fe balance.										
с	Cr	Mn	Мо	Ν	Ni	Р	S	Si		
0.016	17.55	1.18	2.12	0.167	10.1	0.03	0.001	0.25		



the testing. SSRT experiments were conducted with a strain rate of  $1 \times 10^{-6}$  s<sup>-1</sup> to enable an increased number of environmental interactions<sup>78</sup>. SSRT were under aerated conditions at room temperature (RT, 25 °C).

#### **Electrochemical testing**

Before the SCC testing, and without any applied stress, cyclic potentiodynamic polarization (CPP) tests were performed to better understand the corrosion performance of the corrugated (outermost surface influenced by the manufacturing process) and core sections (bulk of the sample, purely austenite matrix) at RT in aerated conditions. Samples were cut and isolated both ends to measure the influence of the manufacturing process, while epoxy mounted samples with the edges covered in epoxy lacquer were used to measure the core section. The potential scan ranged from  $-300 \text{ mV}_{\text{OCP}}$  to  $+1000 \text{ mV}_{\text{OCP}}$  at a potential scan rate of , in accordance to ASTM G61-86<sup>79</sup>. Three chloride 1.667 mV s<sup>-1</sup> concentrations were tested to the corrugated and core sections; 0, 4 and 8 wt.% of Cl<sup>-</sup> by means of CaCl<sub>2</sub> additions to the concrete pore solution (pH 12.6) made out of a saturated Ca(OH)<sub>2</sub> following ASTM G1-03<sup>80</sup>

The electrochemical tests conducted during the straining of the sample were performed using a three-electrode configuration cell setup with a Gamry Reference 600 potentiostat. The reference electrode (RE) used in this test was a saturated calomel electrode (SCE), the counter electrode (CE) used was a graphite rod, and the working electrode (WE) was the AISI 316LN stainless steel reinforcement. All tests were performed in triplicate to ensure reproducibility.

The sequence for the different electrochemical test was the following: first open circuit potential (OCP) was monitored until a steady state OCP value was reached, and then electrochemical impedance spectroscopy (EIS) measurements were performed. The EIS measurements were recorded at the OCP, in a frequency range of  $10^5 - 10^{-2}$  Hz with an applied 10 mV AC excitation signal and at a step rate of 5 points per decade. This set of experiments was repeated with 3 h delay between each sequence, during these 3 h, a potentiostatic scan test monitored the current passing through the AISI 316LN stainless steel reinforcement with an applied potential ( $E_{app}$ ) of +200 mV<sub>SCE</sub>. This  $E_{app}$  was kept constant throughout the entire loading.

#### Microstructure and fracture surface characterization

Microstructural analysis of AISI 316LN stainless steel reinforcement was conducted on cross-section area. The samples were epoxy



Fig. 8 Schematic representation of the SCC testing. (a) electrochemical cell, and (b) sketch of the notched AISI 316LN SS rebar.

mounted and polished to mirror finishing by SiC paper and diamond powder (1 µm). The etchant solution used to reveal the microstructure contained 30 ml HCl, 20 ml HNO<sub>3</sub>, 10 ml deionized water and 3 ml  $H_2O_2$ , samples were exposed for  $10 s^{81}$ . The metallographic study was performed using scanning electron microscope SEM in a Tescan Lyra 3 XMU. The surface morphology of the broken samples was also studied by SEM. The SEM analysis was performed in secondary electron mode (SE) at an accelerating voltage of 15 kV and at a working distance of 10 mm. The Tescan Lyra 3 XMU SEM was also equipped with electron backscattered diffraction (EBSD), which was used to analyze the fraction of the phases and the crack path. The EBSD analysis was carried put at 10 kV accelerating voltage and at 0.5 µm in step size. Finally, X-ray diffraction (XRD) analysis was performed using a Rigaku SmartLab-3kW X-ray diffractometer, with a Cu target ( $K_{\alpha} = 1.5406$  Å), and a scan step of 2 °/min over the 2 $\theta$  range of 35°–95°. The  $\gamma$ – and  $\alpha'$ -phases were elucidated in the XRD patterns. Residual stresses were determined by XRD technique in the core and corrugate regions using the  $sin(2\psi)$  method. A voltage and current inputs of 40 kV and 40 mA were used for the Cu  $K_{\alpha}$  radiation source<sup>82</sup>. The XRD system was calibrated with a stress-free sample, and the ASTM E915-96 was followed to check the X-ray alignment<sup>83</sup>. The free residual stress peak for the XRD analysis was the austenitic (420) Bragg reflection found at  $2\theta = 147^{\circ 84}$ .

#### DATA AVAILABILITY

The data that support the findings of this study are available on request from the corresponding author upon reasonable request.

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#### AUTHOR CONTRIBUTIONS

U.M. conducted the experimental work under the supervision of D.M.B. U.M., N.B., D.D.M. and D.M.B. evaluated and visualized the results. U.M., N.B., D.D.M. and D.M.B. drafted, reviewed and edited the paper. D.M.B. provided funding support.

#### **COMPETING INTERESTS**

The authors declare no competing interests.

#### **ADDITIONAL INFORMATION**

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