

Bulk evidence of anisotropic *s*-wave pairing with no sign change in the kagome superconductor CsV₃Sb₅

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The recently discovered kagome superconductors AV₃Sb₅ (A = K, Rb, Cs) exhibit unusual charge-density-wave (CDW) orders with time-reversal and rotational symmetry breaking. One of the most crucial unresolved issues is identifying the symmetry of the superconductivity that develops inside the CDW phase. Theory predicts a variety of unconventional superconducting symmetries with sign-changing and chiral order parameters. Experimentally, however, superconducting phase information in AV₃Sb₅ is still lacking. Here we report the impurity effects in CsV₃Sb₅ using electron irradiation as a phase-sensitive probe of superconductivity. Our magnetic penetration depth measurements reveal that with increasing impurities, an anisotropic fully-gapped state changes to an isotropic full-gap state without passing through a nodal state. Furthermore, transport measurements under pressure show that the double superconducting dome in the pressure-temperature phase diagram survives against sufficient impurities. These results support that CsV₃Sb₅ is a non-chiral, anisotropic *s*-wave superconductor with no sign change both at ambient and under pressure.

The kagome lattice, a motif consisting of corner-sharing triangles and hexagonal holes, provides a platform for a rich variety of novel quantum phases of matter. Due to its strong geometrical frustration, it has long been studied in quantum spin systems as a playground for realizing quantum spin liquids¹. Recently, however, significant efforts have been devoted to exploring topological metals and semimetals in kagome-lattice systems, in which unique band structures such as flat bands, Dirac cones, and van Hove singularities (vHSs) can lead to Dirac/Weyl fermions^{2,3}, spin/charge ordering, and unconventional superconductivity^{4–6}. Although various topological kagome materials have been reported so far^{2,3}, superconductors with kagome lattices are rarely found⁷.

The recently discovered AV₃Sb₅ (A = K, Rb, Cs) is a new family of kagome superconductors with the superconducting transition temperature T_c of 0.9–2.5 K^{8–10}. The alkali *A* atoms are intercalated between sheets consisting of the two-dimensional (2D) kagome networks of V atoms and triangular and hexagonal networks of Sb atoms (Fig. 1a, b). The electronic band dispersions near the Fermi energy E_F share characteristic features predicted for an ideal kagome-lattice system such as a vHS at the M point and a Dirac point at the K point^{8,10–13}. In these materials, E_F is located near the vHS point, and multiple Fermi surfaces are formed by the V *d*-orbitals and Sb *p*-orbitals (Fig. 1c). Such unique band structures in AV₃Sb₅ give rise to unconventional charge-density-wave (CDW) orders with the transition

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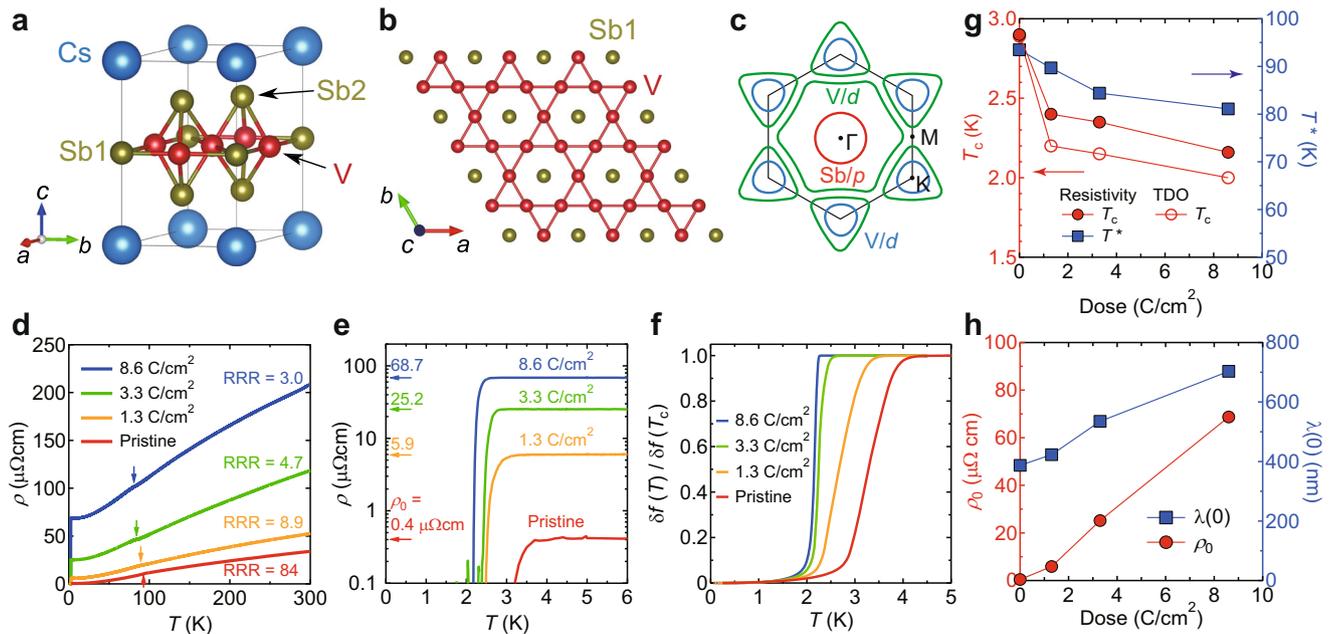


Fig. 1 | Crystal structure, Fermi surface, and electron irradiation effects on the charge-density-wave and superconducting transition temperatures in CsV_3Sb_5 . **a** Crystal structure of CsV_3Sb_5 . **b** V-Sb1 plane viewed from the c -axis direction. Whereas the V atoms form a two-dimensional kagome network, the Sb1 atoms are located at the hexagonal centers. **c** Schematic of the 2D Fermi surface of CsV_3Sb_5 . The circular (red) and hexagonal (green) Fermi surfaces around the Γ point are composed of the Sb p -orbitals and V d -orbitals, respectively, while the two triangular Fermi surfaces (blue and green) around the K point are formed by the V d -orbitals. **d** Temperature dependence of resistivity $\rho(T)$ in CsV_3Sb_5 single crystals with irradiated doses of 0 (pristine, red), 1.3 (orange), 3.3 (green), and 8.6 (blue) C/cm^2 . The RRR values for each sample are listed. Arrows indicate the CDW transition temperatures determined from the temperature derivative of $\rho(T)$ (see Fig. 4b, f, j). Note that the $\rho(T)$ curves for the irradiated samples do not shift

parallel to that of the pristine sample, which can be naturally understood by considering that CsV_3Sb_5 is a multiband system (see Supplementary Information Sec. III). **e** Low-temperature resistivity below 6 K on a logarithmic scale. Arrows indicate the residual resistivity ρ_0 for each irradiated sample. **f** Temperature dependence of normalized frequency shifts of the TDO for each sample. **g** Superconducting and CDW transition temperatures T_c (left axis) and T^* (right axis) as a function of irradiation dose. T_c is defined as the temperature at which the resistivity becomes zero (filled red circles), and the superfluid density becomes finite (open red circles). T^* is defined as the temperature at which the derivative $d\rho/dT$ shows an abrupt change or dip (filled blue squares). **h** ρ_0 (left axis) and $\lambda(0)$ (right axis) as a function of irradiation dose. For the estimation of $\lambda(0)$, see Supplementary Information Sec. IV.

temperature $T^* = 78\text{--}103\text{ K}^{8-10}$ driven by electron correlation^{5,14-17}. The CDW transition is accompanied by a $2a_0 \times 2a_0 \times 2c_0$ or $2a_0 \times 2a_0 \times 4c_0$ superlattice composed of modulated star-of-David and inverse star-of-David patterns (where a_0 and c_0 indicate the lattice constants above T^*), which breaks translational symmetry^{13,18-20}. More intriguingly, it has been reported that additional symmetries, such as time-reversal symmetry (TRS) and rotational symmetry (RS), can be broken below $T^{18,21-25}$. Since the superconducting transition takes place inside the unusual CDW phase, a fundamental question arises as to whether the superconducting pairing mechanism in AV_3Sb_5 is conventional or not²⁶.

Theories on the kagome lattice near van Hove filling have proposed that unconventional superconductivity beyond the electron-phonon mechanism can be realized by electron correlation effects^{4-6,15,27}. Spin and charge fluctuations can lead to spin-triplet p - and f -wave^{6,15,27} and chiral d -wave superconductivity¹⁵, whereas bond-order fluctuations can promote anisotropic s -wave²⁷ and chiral d -wave superconductivity⁵. In support of the above, first-principle calculations have pointed out that the Bardeen-Cooper-Schrieffer (BCS) theory (electron-phonon mechanism) cannot explain the experimental T_c values, suggesting an unconventional pairing mechanism in AV_3Sb_5 ²⁸.

Experimentally, however, the superconducting gap symmetry of AV_3Sb_5 is highly controversial, and whether TRS is broken or not is still elusive. Thermal conductivity measurements in CsV_3Sb_5 ²⁹ and μSR measurements in $\text{Rb}/\text{KV}_3\text{Sb}_5$ ³⁰ have suggested a nodal gap structure. In contrast, magnetic penetration depth³¹ and scanning tunneling spectroscopy (STS)³² studies in CsV_3Sb_5 have suggested a

nodeless gap structure. Nuclear magnetic/quadrupole resonance (NMR/NQR) measurements in CsV_3Sb_5 ³³ have shown a finite Hebel-Slichter coherence peak in $1/T_1T$ and a decrease in Knight shift below T_c , which exclude spin-triplet superconductivity. Regarding the TRS in the superconducting state, Josephson STS measurements in CsV_3Sb_5 ³⁴ have suggested a possible roton pair-density wave, corresponding to an unconventional superconducting state with TRS breaking. Contrastingly, μSR studies in CsV_3Sb_5 ³⁰ have reported that TRS is not broken in the superconductivity state. In addition to the above results at ambient pressure, high-pressure studies^{35,36} have revealed that the CDW phase is suppressed by the application of pressure, accompanied by the emergence of a superconducting dome, indicating the close relationship between the CDW and superconductivity. Moreover, recent μSR experiments under pressure^{37,38} have suggested that TRS is broken in the superconducting state when the CDW phase is suppressed by applying pressure. Therefore, to clarify the pairing mechanism of the kagome superconductors, it is crucial to pin down the superconducting gap symmetry of AV_3Sb_5 both at ambient and high pressure, including whether TRS is broken or not.

In general, the conventional phonon-mediated pairing mechanism leads to a superconducting gap opening all over the Fermi surface, while unconventional pairing mechanisms, such as spin fluctuations, can lead to an anisotropic gap with nodes where the superconducting gap becomes zero. Thus, experimental observations of the low-energy quasiparticle excitations determine whether the gap structure has nodes or not. In addition to clarifying the presence or absence of nodes in the gap, determining the sign of the gap function also provides a

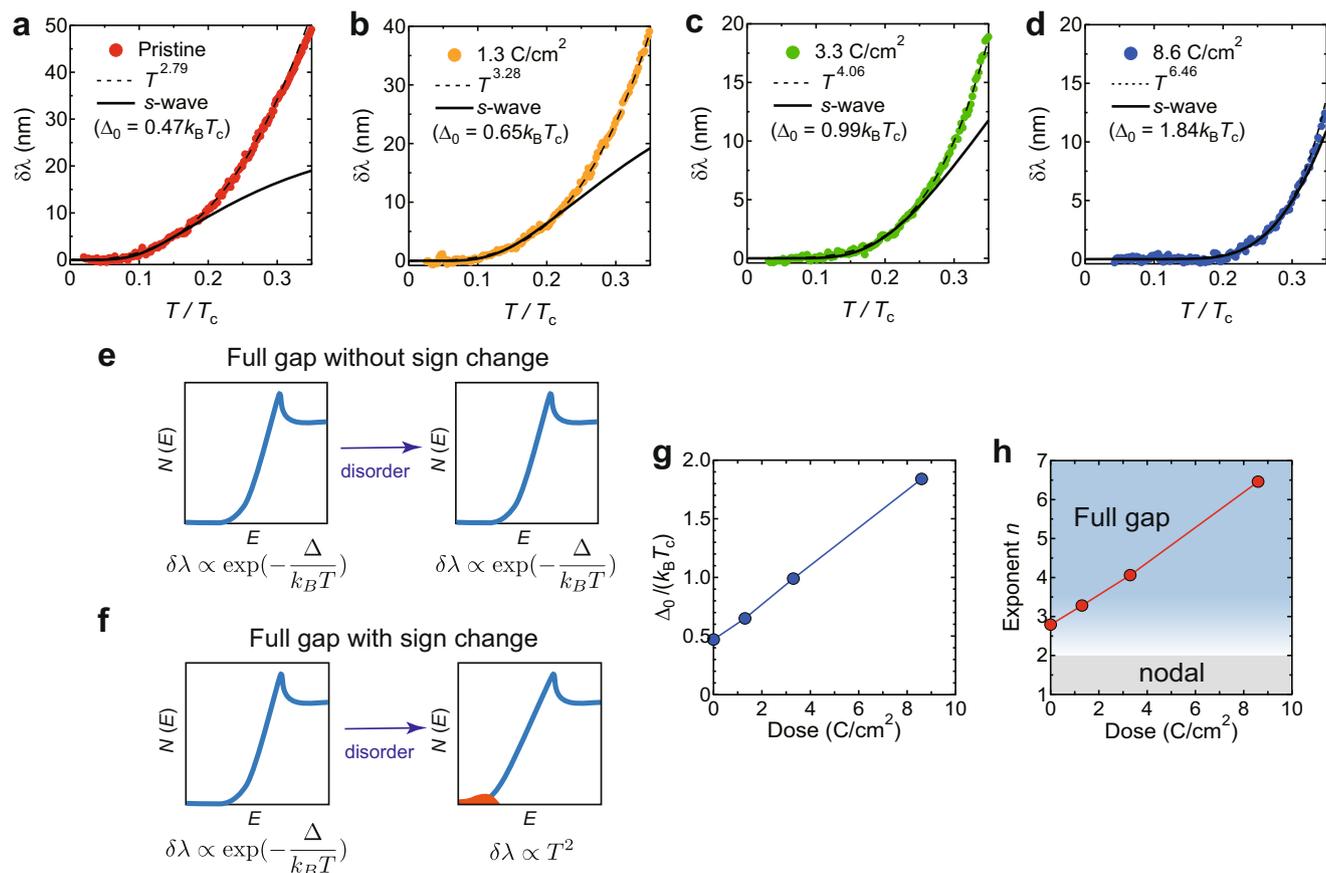


Fig. 2 | Electron irradiation effects on the low-temperature penetration depth. **a–d** Temperature dependence of the change in the penetration depth $\delta\lambda(T)$ for the pristine (**a**), 1.3 (**b**), 3.3 (**c**), and 8.6 (**d**) C/cm² irradiated samples. Solid circles are the experimental data. Black solid and dashed lines are fitting curves of the fully gapped *s*-wave and power-law analysis. **e, f** Schematics of the change in the quasi-particle density of states against disorder in the case of fully gapped states without

(**e**) and with (**f**) sign change. **g** Gap value $\Delta_0/k_B T_c$ obtained from the *s*-wave fit to $\delta\lambda(T)$ at low temperatures as a function of irradiation dose. **h** Exponent value n obtained from the power-law fitting to $\delta\lambda(T)$ up to $0.3 T_c$ as a function of irradiation dose. The exponent $n \leq 2$ (gray shaded region) indicates a nodal gap structure, while $n \geq 3$ implies a fully gapped state (blue shaded region).

strong constraint on the superconducting pairing symmetry. Especially in kagome superconductors, theories have predicted that two degenerated superconducting order parameters, $d_{x^2-y^2}$ and d_{xy} , give rise to chiral $d_{x^2-y^2} \pm id_{xy}$ -wave symmetry, where a finite gap opens all over the Fermi surface, but the phase of the order parameter changes by 4π in momentum \mathbf{k} space³⁹. Therefore, phase-sensitive probes are highly required to determine the pairing symmetry of AV_3Sb_5 .

There are several experimental probes that are sensitive to the sign of gap functions, such as Josephson junction⁴⁰, quasiparticle interference⁴¹, and neutron scattering techniques⁴². In general, however, the analysis of such interference effects is complicated in multiband systems due to the complexity of the scattering processes. In addition, most of them require good surface/interface conditions. In contrast, the non-magnetic impurity effect, on which we focus here, is one of the phase-sensitive probes that are applicable to multiband systems and reflect the bulk superconducting properties^{43,44}. In *s*-wave superconductors with no sign-changing order parameter, the Cooper pairs are not destroyed by non-magnetic impurity scattering, and both T_c and quasiparticle density of states (DOS) are little affected by disorder (the so-called Anderson's theorem⁴⁵) (Fig. 2e). In contrast, in the case of nodeless superconductors with sign-changing order parameters, such as chiral *d*-wave and s_{\pm} -wave superconductors (note that considering the electronic structure of the present kagome system, one cannot expect sufficient interband interactions to induce the s_{\pm} -wave superconductivity), the Cooper pairs are destroyed by impurity

scattering, which suppresses T_c rapidly and induces impurity states associated with the Andreev bound state (Fig. 2f). In this case, additional low-energy quasiparticle excitations appear near the zero energy, e.g., leading to a change in the temperature dependence of the magnetic penetration depth λ from exponential to T^2 ^{43,46}.

Here, we show that superconductivity in CsV_3Sb_5 is robust against impurities both at ambient and under high pressure. Our magnetic penetration depth measurements reveal that with increasing impurities, a highly anisotropic fully gapped superconducting state changes gradually to an isotropic full-gap state without showing impurity-induced Andreev bound states, which excludes any of sign-changing symmetries. Moreover, transport measurements under high pressure show that the superconducting dome in the pressure-temperature (*P*-*T*) phase diagram survives against sufficient impurities. These results suggest that the superconducting gap function in CsV_3Sb_5 is non-chiral and non-sign-changing *s*-wave.

Results

Electron irradiation effects on the CDW and superconducting transition temperatures

In this study, we used electron irradiation to systematically introduce non-magnetic impurities into CsV_3Sb_5 single crystals (see Methods and Supplementary Information Sec. I). In this method, high-energy electron beam irradiation creates vacancies in the crystal⁴³, acting as point defects without changing the electronic structure and lattice constants (see Supplementary Information Sec. II). Figure 1d, e

shows the temperature dependence of resistivity $\rho(T)$ at ambient pressure in samples with irradiated doses of 0 (pristine), 1.3, 3.3, and 8.6 C/cm². The residual resistivity ratio (RRR) of the pristine sample is ~84, indicating the high quality of our crystals. As the dose increases, the residual resistivity ρ_0 increases (also see Fig. 1h), and the RRR value decreases. The change in $\rho(T)$ with impurities is successive for the irradiation dose in the whole temperature range. Furthermore, our X-ray structural analysis and Hall coefficient measurements show the absence of any change in the lattice parameters and carrier density induced by electron irradiation. These results indicate that the non-parallel shift with impurities in $\rho(T)$ is most likely due to the multiband nature of the present kagome system (for more details, see Supplementary Information Sec. III). Along with this, both the CDW and superconducting transition temperatures T and T_c shift to a lower temperature (Fig. 1g). In general, non-magnetic impurity scattering can suppress long-range orders because the introduced defects shorten the coherence length. Indeed, the suppression of CDW order by impurities has been theoretically studied⁴⁷. The suppression of T_c has also been confirmed by the Meissner effect measured by the normalized frequency shift of a tunnel diode oscillator (TDO) (Fig. 1f). Note that the superconducting transition becomes sharper with increasing dose, which may be related to the suppression of superconducting phase fluctuations^{7,48} or the change in skin depth due to impurity scattering. The sharp superconducting transition width in the 8.6 C/cm² irradiated sample with sufficient disorder indicates that the defects are introduced quite uniformly inside the crystals.

Non-magnetic impurity effects on low-energy quasiparticle excitations in the superconducting state

Next, we turn to the impurity effect on low-energy quasiparticle excitations in the superconducting state. Magnetic penetration depth λ is one of the most fundamental properties of superconductors sensitive to low-energy quasiparticle excitations^{7,31,43}. In this study, we measured the magnetic penetration depth of the pristine and irradiated CsV₃Sb₅ single crystals down to 50 mK by using the TDO in a dilution refrigerator (see Methods). Figure 2a–d shows the change in the magnetic penetration depth $\delta\lambda(T) \equiv \lambda(T) - \lambda(0)$ (where $\lambda(0)$ is the absolute value of the penetration depth at 0 K) at low temperatures for the pristine and 1.3, 3.3, and 8.6 C/cm² irradiated samples. In the pristine sample, $\delta\lambda(T)$ shows a flat temperature dependence at low temperatures below $0.1T_c$ (Fig. 2a), indicating a fully gapped superconducting state in CsV₃Sb₅. To examine the low-energy quasiparticle excitations in the pristine sample, we applied a power-law fit $\delta\lambda(T) \propto T^n$ to the experimental data. In general, in the case of nodal superconductors with line and point nodes, the exponent value n gives 1 and 2 in the clean limit, respectively. We obtained $n \sim 2.8$ from the fitting (Fig. 2a), indicating the absence of nodes in the gap (or conversely, the presence of a finite gap). Then, to quantitatively evaluate the gap value, we tried to fit the data with the fully gapped *s*-wave model $\delta\lambda(T) \propto T^{-1/2} \exp(-\Delta_0/k_B T)$, where k_B is the Boltzmann constant and Δ_0 is the superconducting gap. We obtained an extremely small gap value $\Delta_0 = 0.47k_B T_c$ (which is consistent with the previous study³¹), suggesting the existence of gap minima Δ_{\min} coming from the anisotropic gap nature of CsV₃Sb₅, as discussed later. One of our key findings is that the fully gapped behavior in $\delta\lambda(T)$ is robust against disorder (Fig. 2b–d). The flat temperature region at low temperatures expands to a higher temperature region with increasing dose. In the case of fully gapped superconductors with sign-changing order parameters, $\delta\lambda(T)$ is expected to change from an exponential to a T^2 dependence with increasing impurities because of the impurity-induced DOS (Fig. 2f)⁴³. In sharp contrast, our experimental observations show that Δ_0 and n obtained from the fitting rather increase with increasing dose (Fig. 2g, h), indicating no impurity-induced DOS in the superconducting gap. These results

provide strong bulk evidence that the superconducting gap structure of CsV₃Sb₅ is nodeless without a sign-changing gap.

For a more detailed analysis of the superconducting gap structure, we derived the normalized superfluid density $\rho_s(T) \equiv \lambda^2(0)/\lambda^2(T)$. We used $\lambda(0) = 387$ nm for the pristine sample estimated in the previous study³¹ and calculated $\lambda(0)$ for the irradiated samples by using the relation $\lambda(0) = \lambda_L(0)(1 + \xi/l)^{1/2}$ (Fig. 1h), where the London penetration depth $\lambda_L(0)$ is assumed to be equal to $\lambda(0) = 387$ nm for the pristine sample, and l and ξ are the mean free path and coherence length, respectively (for more details, see Supplementary Information Sec. IV). Figure 3a shows the obtained $\rho_s(T)$ curve as a function of T/T_c for each sample. In all the samples, ρ_s shows a flat temperature dependence at low temperatures, which extends to a higher temperature region with increasing dose. This is again inconsistent with a nodal gap structure.

Here, we consider a multigap model to analyze the overall temperature dependence of ρ_s . In CsV₃Sb₅, the Fermi surfaces are formed by two different orbitals: one is derived from the *d*-orbitals of V, forming a hexagonal Fermi surface around the Γ point and two triangular Fermi surfaces around the K point, while the other is from the *p*-orbitals of Sb, forming a circular Fermi surface around the Γ point (Fig. 1c)¹³. The Fermi surfaces derived from the V *d*-orbitals determine the physical properties of this material, and three equivalent **q** vectors^{18,21,23,24} are considered to give rise to anisotropic pairing interactions^{5,27}. Indeed, recent STM measurements³² have reported the emergence of an anisotropic superconducting gap as well as an isotropic gap below T_c . We therefore consider a multigap model with an anisotropic but nodeless superconducting gap with six-fold symmetry ($\Delta_1 \propto 1 + \alpha \cos(6\phi)$) and an isotropic superconducting gap ($\Delta_2 = \text{const.}$) on two cylindrical Fermi surfaces (Fig. 3b). Note that an isotropic two-gap model cannot produce reasonable results for irradiated samples (for more details, see Supplementary Information Sec. V). We fitted the experimental data with the anisotropic multigap model (Fig. 3a) and obtained the gap values Δ_1 and Δ_2 as a function of dose (Fig. 3c). As the dose increases, the difference between the maximum and minimum values of Δ_1 decreases, and eventually, all the gaps become almost identical. This is due to the averaging effect between the two gaps introduced by impurity-induced intra/interband scattering, and a very similar behavior has been observed in the prototypical multigap superconductor MgB₂⁴⁹. This evidences nodeless multigap superconductivity with a sign-preserving order parameter in CsV₃Sb₅, which excludes the possibility of spin-triplet *p*- and *f*-wave and chiral *d*-wave superconductivity.

Pair-breaking effect

To discuss the impurity effect on T_c more quantitatively, we next introduce a pair-breaking parameter $g = \hbar/(\tau_{\text{imp}} k_B T_{c0})$, where $\tau_{\text{imp}} = \mu_0 \lambda_L^2(0)/\rho_0$ is the impurity scattering time and T_{c0} is the superconducting transition temperature of the pristine sample^{44,50}. The suppression of T_c is plotted as a function of g and compared to other superconductors with and without sign-changing order parameters (Fig. 3d). T_c of CsV₃Sb₅ is rapidly suppressed at a low irradiation dose but starts to saturate at moderate irradiation doses. The initial rapid suppression of T_c is considered to be related to the reduction of the anisotropy of Δ_1 (Fig. 3c), as discussed later. The T_c suppression above 1.3 C/cm² irradiation dose is much slower than those in superconductors with sign-changing order parameters such as *d*-wave, rather similar to those in *s*-wave superconductors without sign-changing gaps. These results also support that multigap *s*-wave superconductivity with no sign change is realized in CsV₃Sb₅ at ambient pressure.

Impurity effects on the high-pressure superconducting phase

To further investigate the non-magnetic impurity effect on the superconducting phase of CsV₃Sb₅ under pressure, we constructed the

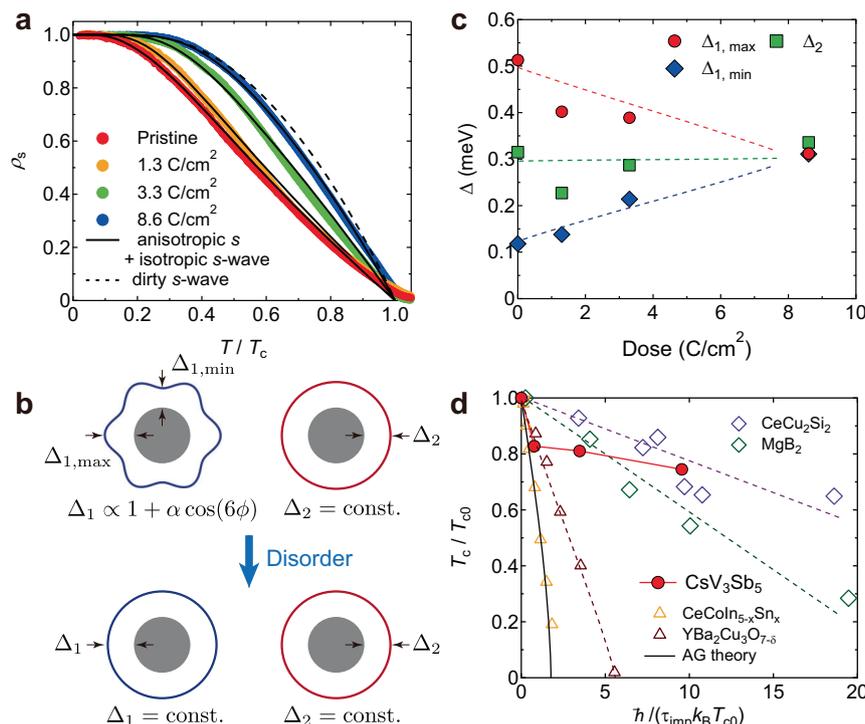


Fig. 3 | Pair-breaking effect in CsV₃Sb₅. **a** The Temperature dependence of normalized superfluid density $\rho_s(T) \equiv \lambda^2(0)/\lambda^2(T)$ for the pristine (red), 1.3 (orange), 3.3 (green), and 8.6 (blue) C/cm² irradiated samples. Black solid lines are the fitting curves of the multigap model (clean limit). The black dashed line is the curve for the dirty *s*-wave (single gap) case. Note that $\rho_s(T)$ for the 8.6 C/cm² irradiated sample with the relatively large value of $\xi/l = 2.3$ approaches the dirty limit (see Supplementary Information Sec. V). **b** Schematic picture of the change in the superconducting gap structure against disorder. The introduction of electron irradiation changes the anisotropic superconducting gap structure to an isotropic one. **c** Gap sizes obtained from the fitting analysis of the superfluid density as a function of dose. Red circles and blue diamonds represent the maximum and minimum values

of the anisotropic gap, $\Delta_{1, \text{max}}$ and $\Delta_{1, \text{min}}$, respectively. Green squares represent the gap values of the isotropic gap Δ_2 . Dotted lines are guides for the eyes.

d Suppression of T_c in CsV₃Sb₅ (red circles) as a function of pair breaking parameter $g = \hbar/(\tau_{\text{imp}}k_B T_{c0})$. For comparison, the results for Sn-substituted CeCoIn₅ (yellow triangles)⁵¹ and electron-irradiated YBa₂Cu₃O_{7- δ} (brown triangles)⁵² are plotted as examples of *d*-wave superconductors. Black solid line represents the suppression of T_c expected in the Abrikosov-Gor'kov (AG) theory. Also, the results of neutron-irradiated MgB₂ (green diamonds)⁴⁹ and electron-irradiated CeCu₂Si₂ (purple diamonds)⁵⁰ are plotted as examples of multigap *s*-wave superconductors. Dotted lines are guides for the eyes.

P-T phase diagram in the pristine sample and 4.8 and 8.6 C/cm² irradiated samples. Figure 4a shows the $\rho(T)$ curve of the pristine sample at several pressures. The CDW transition temperature T , which is determined from a jump or dip in $d\rho(T)/dT$ (Fig. 4b), decreases monotonically with increasing pressure (Fig. 4d). In contrast, the superconducting transition temperature T_c shows a non-monotonic pressure dependence (Fig. 4c), and a double superconducting dome is observed, as reported in previous high-pressure studies^{35,36} (Fig. 4d). The first peak of the superconducting double dome locates at $P_1 \sim 0.7$ GPa inside the CDW phase, while the second peak locates at $P_2 \sim 2$ GPa near the CDW endpoint. We conducted the same experiments for the 4.8 and 8.6 C/cm² irradiated samples (Fig. 4e–l) and obtained the *P-T* phase diagrams as shown in Fig. 4h, i. T is suppressed in the 4.8 and 8.6 C/cm² irradiated samples, and the CDW endpoint shifts to lower pressure with increasing irradiation dose. T_c is also suppressed by disorder, but the superconducting dome survives even after 8.6 C/cm² irradiation.

As already discussed in Fig. 3d, the irradiation dose of 8.6 C/cm² introduces enough defects to suppress superconductivity with a sign-changing order parameter. To investigate the impurity effect on the high-pressure superconducting phase, we evaluated the pair-breaking parameter g at P_2 for each dose (Fig. S5). Our results display that the suppression of T_c at the second dome is slower than that of the *d*-wave case with a sign-changing order parameter; it rather traces the trend in MgB₂. Therefore, these results suggest that the superconducting gap symmetry of CsV₃Sb₅ at high pressure is also non-sign-changing *s*-wave.

Discussion

Recent μ SR measurements under pressure³⁷ have reported that the superconducting pairing symmetry near P_2 has a finite superconducting gap across the Fermi surface and breaks TRS. As a possible symmetry, chiral $d_{x^2-y^2} \pm id_{xy}$ or $p_x \pm ip_y$ -wave superconductivity has been discussed^{37,38}. However, such unconventional superconductivity is expected to be sensitive to disorder, because the chiral states have sign-changing order parameters which would produce Andreev bound states by impurities. Our present results show that the superconducting state under high pressure is robust against disorder. These findings seem to be inconsistent with the chiral $d_{x^2-y^2} \pm id_{xy}$ and $p_x \pm ip_y$ -wave states. To fully understand the relationship between our impurity effects and the μ SR results, further theoretical and experimental studies on the high-pressure phase in the kagome systems are highly desired. We note that the TRS breaking has been observed in the CDW phase, and thus the possible fluctuations of chiral CDW order in the high-pressure phase on the time scale of μ SR measurements may be an important issue.

Another important aspect of our findings is that the CDW endpoint shifts to lower pressure with irradiation, followed by the second peak of the superconducting double dome (Fig. 4d, h, i), suggesting that the CDW is closely related to the superconductivity in the present system. Recent theoretical calculations in AV₃Sb₅²⁷ have proposed that bond-order fluctuations originating from the triple- \mathbf{q} vectors corresponding to the (inverse) star of David pattern induce anisotropic pairing interactions, leading to anisotropic *s*-wave superconductivity. This theory can explain the relatively high T_c in AV₃Sb₅ that cannot be

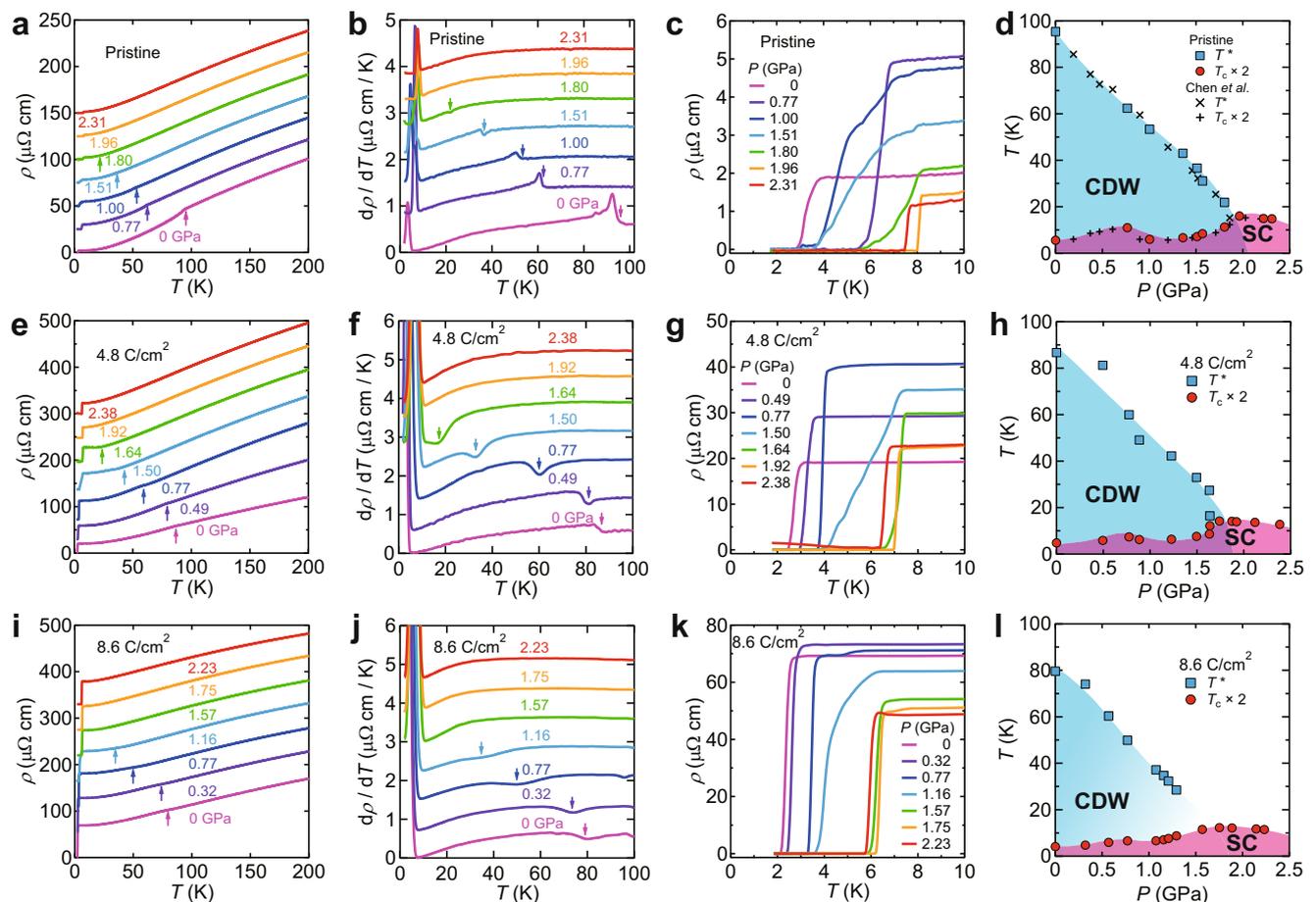


Fig. 4 | Pressure versus temperature phase diagrams of pristine and irradiated CsV_3Sb_5 . **a, e, i** $\rho(T)$ curves below 200 K at several pressures for the pristine (**a**), 4.8 (**e**), and 8.6 (**i**) C/cm^2 irradiated samples. Arrows indicate the CDW transitions determined from the $d\rho/dT$ curves in **b, f, j**. **b, f, j** Temperature dependence of $d\rho/dT$ below 100 K for the pristine (**b**), 4.8 (**f**), and 8.6 (**j**) C/cm^2 irradiated samples. Arrows indicate the CDW transitions. **c, g, k** Low-temperature $\rho(T)$ curves below 10 K at several pressures for the pristine (**c**), 4.8 (**g**), and 8.6 (**k**) C/cm^2 irradiated samples. Note that the resistivity at 2.38 GPa does not reach zero at low temperatures in **g**. This may come from distortion such as microcracks to the sample and damage to the terminals caused by solidification of daphne 7373 above 2 GPa.

d, h, l P - T phase diagrams of the pristine (**d**), 4.8 (**h**), and 8.6 (**l**) C/cm^2 irradiated samples. For clarity, T_c is doubled. The phase diagram of the pristine sample includes data from Chen et al.³⁵. The CDW (superconducting (SC)) phase is shaded in blue (red). Note that the double superconducting dome becomes broader for the irradiated samples. In the intermediate pressure region between P_1 and P_2 , a significant broadening of the superconducting transition has been observed, and T_c shows a strong sample dependence^{35,36}. These behaviors indicate that the superconducting state in the intermediate pressure region is inhomogeneous and sensitive to the microscopic disorder within the crystal. Therefore, the broadening of the double superconducting dome may come from the introduction of disorders.

reproduced by the BCS theory²⁸ and the anisotropic superconducting gap structure in CsV_3Sb_5 obtained in the present study. Moreover, in such anisotropic s -wave superconductivity, the introduction of impurity scattering averages out the anisotropic gap, changing to the isotropic gap. In this case, T_c drops rapidly at an initial introduction of impurities, but as the gap becomes isotropic, the reduction of T_c saturates and becomes much slower than that expected in the Abrikosov-Gor'kov (AG) theory. These expectations are in good agreement with our observations of the T_c suppression in CsV_3Sb_5 . We note that a possible transition from a p -wave to an s -wave state caused by impurities²⁷ may explain the initial rapid suppression of T_c . However, our present results exclude the possibility of a nodal superconducting state in the pristine sample, which is at odds with the p -wave state. The gradual change in the superconducting gap inferred from the temperature dependence of the superfluid density suggests that an impurity-induced transition from a full-gap chiral state to a non-chiral s -wave state is also unlikely. This is reinforced by the μSR measurements at ambient pressure³⁰, which report that chiral superconductivity in the pristine sample at ambient pressure can be ruled out. Thus, our present results support a new type of unconventional superconductivity due to bond-order fluctuations on the kagome

lattice in CsV_3Sb_5 , where the gap function is non-sign-changing s -wave. In the present kagome superconductors, the possibility of a loop-current phase with broken TRS and a nematic phase with broken RS has been pointed out above the superconducting phase^{18,21–25}. Therefore, elucidating the intertwining of these unusual normal and superconducting phases, which is commonly seen in high- T_c cuprates and iron-based superconductors, will pave the way to understanding novel quantum phases of matter in condensed matter physics.

Methods

Single crystal growth

High-quality single crystals of CsV_3Sb_5 were synthesized using the self-flux method. All sample preparations are performed in an argon glovebox with oxygen and moisture < 0.5ppm. The flux precursor was formed through mechanochemical methods by mixing Cs metal (Alfa 99.98%), V powder (Sigma 99.9%), and Sb beads (Alfa 99.999%) to form a mixture which is ~50 at.% $\text{Cs}_{0.4}\text{Sb}_{0.6}$ (near eutectic composition) and 50 at.% VSb_2 . Note that prior to mixing, as-received vanadium powders were purified in-house to remove residual oxides. After milling for 60 min a pre-seasoned tungsten carbide vial, flux precursors are extracted and sealed into 10 mL alumina crucibles. The crucibles are

nested within stainless steel jackets and sealed under argon. Samples are heated to 1000 °C at 250 °C/h and soaked for 24 h before dropping to 900 °C at 100 °C/h. Crystals are formed during the final slow cool to 500 °C at 1 °C/h before terminating the growth. Once cooled, the crystals are recovered mechanically. Samples are hexagonal flakes with a brilliant metallic luster. The elemental composition of crystals was assessed using energy-dispersive X-ray spectroscopy (EDS) using an APREO-C scanning electron microscope.

Electron irradiation

Electron irradiation with the incident electron energy of 2.5 MeV was performed on SIRIUS Pelletron accelerator operated by the Laboratoire des Solides Irradiés (LSI) at École Polytechnique. To prevent defect migration and agglomeration, the sample temperature was kept at -20 K during irradiation which produces stable vacancy-interstitial Frenkel pairs. Subsequent warming to room temperature causes annealing of interstitials, which have a lower migration energy, leaving a uniform population of vacancy type defects. The electron irradiation of 1.3, 3.3, and 8.6 C/cm² was performed at the same beam time (run#1), while the irradiation of 4.8 C/cm² was conducted at another beam time (run#2).

Electrical resistivity measurements

The electrical resistivity was measured at ambient and high pressure by the 4-terminal method using a Physical Property Measurement System (PPMS) from Quantum Design with the lowest temperature of about 1.8 K. The resistivity under pressure was measured using a piston cylinder cell to generate pressure up to -2.5 GPa and daphne 7373 as a pressure medium in PPMS. The pressure value in the sample was determined from the superconducting transition temperature T_c of Pb under pressure, using the relation $P = (7.20 - T_c)/0.365$. Note that when the 4.8 C/cm² irradiated sample was set in the piston cell, even before pressure was applied, the resistivity value changed, probably due to cracks, so we have corrected it to the value before pressure was applied.

Magnetic penetration depth measurements

The temperature variation of the in-plane magnetic penetration depth $\delta\lambda(T) = \lambda(T) - \lambda(0)$ was measured by using a tunnel diode oscillator technique (TDO) with the resonant frequency of -13.8 MHz in a dilution refrigerator down to -50 mK. The sample was mounted on a sapphire rod with Apiezon N grease, then inserted into a copper coil in the LC circuit. The frequency shift δf in the TDO is related to the change of magnetic susceptibility $\delta\chi$ by the following equation, $\delta f = -(f_0 V_s / (2V_c(1 - N)))\delta\chi$, where f_0 is the resonant frequency without the sample, V_s and V_c are the sample and coil volume, respectively, and N is the demagnetization factor. $\delta\chi$ is related to $\delta\lambda$ by the following equation, $\delta\chi = \delta\lambda/R$, where R is a constant determined by the geometry of the sample from the calculation. Thus, δf is related to $\delta\lambda$ by the following equation, $\delta f = -(f_0 V_s / (2RV_c(1 - N)))\delta\lambda$.

Data availability

All data supporting the findings of this study are available within the paper and its Supplementary Information. Source data are provided with this paper.

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Author contributions

K.H. and T.S. conceived the project. M.R., K.I., K. Ogawa, K. Okada, Y.M., K.H., and T.S. performed magnetic penetration depth measurements and analyzed the data. Y.T., K. Okada, K.M., and Y.U. carried out high-pressure measurements. M.R. and S.L. performed electrical transport and X-ray diffraction measurements. R.G. and M.K. conducted electron irradiation experiments. B.R.O. and S.D.W. carried out sample preparation. M.R., K.H., and T.S. prepared the manuscript with inputs from R.G., M.K., B.R.O., and S.D.W. All authors discussed the experimental results.

Competing interests

The authors declare no competing interests.

Additional information

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