

ARTICLE

DOI: 10.1038/s41467-018-05307-0

OPEN

In situ topographical chemical and electrical imaging of carboxyl graphene oxide at the nanoscale

Weitao Su^{1,2}, Naresh Kumar^{3,4}, Andrey Krayev⁵ & Marc Chaigneau⁶

Visualising the distribution of structural defects and functional groups present on the surface of two-dimensional (2D) materials such as graphene oxide challenges the sensitivity and spatial resolution of the most advanced analytical techniques. Here we demonstrate mapping of functional groups on a carboxyl-modified graphene oxide (GO-COOH) surface with a spatial resolution of \approx 10 nm using tip-enhanced Raman spectroscopy (TERS). Furthermore, we extend the capability of TERS by measuring local electronic properties in situ, in addition to the surface topography and chemical composition. Our results reveal that the Fermi level at the GO-COOH surface decreases as the I_D/I_G ratio increases, correlating the local defect density with the Fermi level at nanometre length-scales. The in situ multi-parameter microscopy demonstrated in this work significantly improves the accuracy of nanoscale surface characterisation, eliminates measurement artefacts, and opens up the possibilities for characterising optoelectronic devices based on 2D materials under operational conditions.

¹ College of Materials and Environmental Engineering, Hangzhou Dianzi University, 310018 Hangzhou, China. ² Key Laboratory of RF Circuits and Systems (Hangzhou Dianzi University), Ministry of Education, Hangzhou 310018, China. ³ National Physical Laboratory, Hampton Road, Teddington, Middlesex TW11 OLW, UK. ⁴ Inorganic Chemistry and Catalysis group, Debye Institute for Nanomaterials Science, Utrecht University, Universiteitsweg 99, 3584 CG Utrecht, The Netherlands. ⁵ HORIBA Instruments Incorporated, Novato, CA 94949, USA. ⁶ HORIBA France, Avenue de la Vauve, Passage Jobin Yvon, 91120 Palaiseau, France. These authors contributed equally: Weitao Su, Naresh Kumar. Correspondence and requests for materials should be addressed to W.S. (email: suweitao@hdu.edu.cn) or to N.K. (email: naresh.kumar@npl.co.uk) or to M.C. (email: Marc.CHAIGNEAU@horiba.com)

n recent years, graphene oxide (GO) has attracted much attention in the field of two-dimensional (2D) materials research because of its many potential applications, including lithium batteries¹, energy harvesting^{2,3} and novel medicines⁴. A GO flake typically comprises a few-nanometre thick sheet of carbon atoms bonded to functional groups, such as carboxyl, carbonyl, hydroxyl or nitro⁵⁻⁷. These functional groups can be further linked to RNA-aptamers⁸ or even peptides⁹ resulting in functionalised GO⁶ with fascinating tunable properties. However, visualising the distribution of functional groups on a GO surface at the nanoscale challenges the chemical sensitivity and spatial resolution of most conventional analytical techniques, such as Xray photoelectron spectroscopy (XPS) (millimetre spatial resolution)^{7,10}, confocal Raman spectroscopy (200-300 nm spatial resolution)^{7,10} and aberration-corrected transmission electron microscopy (TEM) (limited molecular information)¹¹.

Tip-enhanced Raman spectroscopy (TERS) has emerged as a powerful analytical technique providing high chemical sensitivity for surface molecular mapping with nanoscale spatial resolution under ambient conditions¹²⁻¹⁴. In TERS, a metallic scanning probe microscopy (SPM) probe placed at the focal point of a laser undergoes localised surface plasmon resonance (LSPR), which together with the lightening rod effect leads to the enhancement and confinement of the electric field at the TERS probe apex. This effect simultaneously improves the sensitivity as well as the spatial resolution of Raman microscopy by enhancing the Raman signal from analyte molecules directly underneath the TERS probe. In previous reports, TERS has been successfully used to characterise ultrathin organic films¹⁵, single molecules¹⁶ and various material properties of graphene, such as the number of layers¹⁷⁻¹⁹, local strain²⁰, edge orientation^{14,21}, surface adsorbates^{13,22} and artificial defects²³. However, to the best of our knowledge, the capability of TERS to detect functional groups on GO or a functionalised GO surface has not been demonstrated^{13,17,18,24}.

Herein, we present nanoscale surface mapping of structural defects and functional groups on few-layer carboxyl-modified graphene oxide (GO–COOH) flakes with ≈ 10 nm spatial resolution using TERS. Furthermore, we take the surface characterisation a step further by combining TERS with Kelvin probe force microscopy (KPFM) and demonstrate in situ topographical, chemical and electrical nanoscopy of a GO–COOH surface. These measurements reveal an inverse correlation between the local defect density and contact potential difference (CPD) at GO–COOH surface. This work paves the way for optimisation of optoelectronic devices based on 2D materials via simultaneous, non-destructive and multi-parameter characterisation of their surface properties in situ.

Results

High-resolution TERS mapping of GO-COOH. Microchemical analysis of the GO-COOH sample used in this study was carried out using XPS (Supplementary Fig. 1), Fourier-transform infrared (FTIR) spectroscopy (Supplementary Fig. 2a) and Raman spectroscopy (Supplementary Fig. 2b). The XPS binding energies and the fingerprint FTIR absorption bands indicate that the GO-COOH sample consists of mainly carbon and oxygen, along with the C-O-C, C-O, C-CH₃, C=O, COOH and C-H functional groups. See Supplementary Note 1 for details.

A schematic diagram of the reflection mode atomic force microscopy (AFM)-TERS setup used in this work is shown in Fig. 1. Figure 2a shows the topography map of a GO–COOH sample containing multilayer flakes, which was recorded during the TERS mapping measurement. A height profile across the topography map shown in Supplementary Fig. 3 indicates that the sample contains both few-layer (1–2 layers) and thick-layer

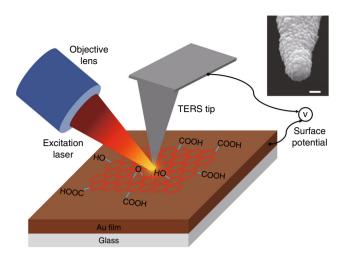


Fig. 1 Schematic diagram of the experimental setup. The side illumination AFM-TERS setup used in this work. Au coated AFM probe in combination with 638 nm excitation laser is used for in situ topographical, chemical and electrical nanoscopy of a GO-COOH sample. Scanning electron microscopy (SEM) image of a representative Au coated TERS probe is shown inset. Scale bar: 100 nm

(about 5 layers) GO-COOH flakes. The spatial resolution of this topography map is relatively poor due to the large radius (\approx 100 nm) of the Au coated TERS probe apex (inset of Fig. 1) compared to typical AFM probes that typically have an apex radius of $<10 \text{ nm}^{14}$. High-resolution TERS mapping was conducted on the GO-COOH sample area shown in Fig. 2a. The stacked TERS spectra measured at each pixel in the TERS map are shown in Fig. 2b. These stacked TERS spectra exhibit graphene D and G bands, which arise from the scattering by a defect phonon and the first-order scattering by the E_{2g} phonon, respectively²⁵, along with the Raman bands corresponding to the functional groups observed in the XPS and FTIR measurements presented in Supplementary Figs. 1 and 2, respectively. Assignment of the Raman bands of six functional groups observed in the TERS spectra along with the D and G bands is presented in Table 1. See Supplementary Note 1 and 2 for details about the assignment of these bands.

Each spectrum in the TERS map was fitted using eight Lorentzian curves at the band positions listed in Table 1 and eight TERS maps were generated. The TERS intensity maps of the D and G bands are shown in Fig. 2c, d, respectively. In these TERS maps, big (~1 µm) and small (100-300 nm) GO-COOH flakes can be clearly visualised. The edges and steps of different flakes in the TERS maps are much clearer than in the corresponding topography map shown in Fig. 2a indicating a higher TERS signal enhancement at these locations. The most likely cause of the higher signal enhancement at the steps and edges is the disruption of the GO-COOH lattice leading to a better alignment of the Raman polarisability tensor with the axial polarisation of the tip-enhanced near-field resulting from the gap-mode configuration^{16,26-28}. Furthermore, high TERS signal enhancement for the D and G bands is restricted to a few pixels at the steps and edges of the GO-COOH flakes confirming the high spatial resolution of these TERS maps.

During TERS mapping, markedly different spectra were measured in different areas of the GO-COOH flakes indicating the heterogeneity of surface functionalisation and non-uniform enhancement of the D and G bands. Fig. 2e shows four representative TERS spectra measured at the locations marked 1–4 in Fig. 2c. The TERS spectrum of location 1 represents an averaged spectrum of 10 pixels from a low intensity region. This

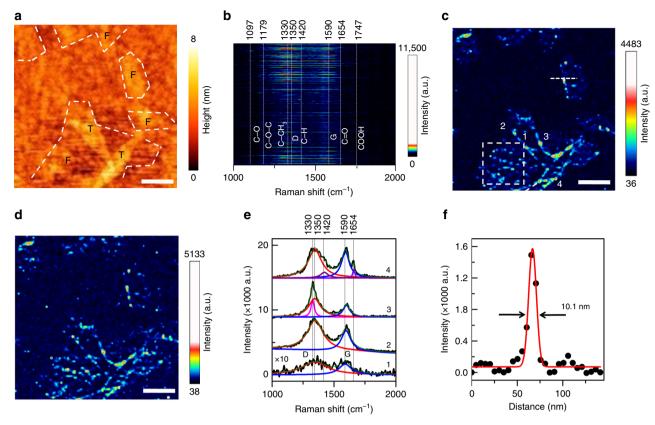


Fig. 2 High-resolution TERS mapping of a GO-COOH flake. **a** Topography map of a multilayer GO-COOH sample obtained whilst TERS mapping. "T" and "F" refer to thick-layer and few-layer GO-COOH flakes, respectively. **b** Image of the stacked TERS spectra measured from 100 × 100 pixels across the GO-COOH sample area shown in (**a**). TERS spectra are stacked in the order of their acquisition with the 1st and the 10,000th TERS spectrum presented at the bottom and top of the image, respectively. TERS maps of **c** D band (1350 cm⁻¹) intensity and **d** G band (1590 cm⁻¹) intensity measured from the GO-COOH sample area shown in (**a**). Nominal pixel dimension: 10 nm. Integration time: 0.4 s. **e** TERS spectra measured at the locations marked 1-4 in (**c**) along with the fitted Lorentzian curves. The intensity of the averaged TERS spectrum from location 1 has been multiplied by 10 for easier visualisation. **f** TERS intensity profile along the white dashed line marked in (**c**) fitted with a Gaussian curve. All scale bars: 200 nm

Table 1 Proposed assignment of Raman bands observed inthe TERS spectra in Fig. 2b	
Raman band position (cm ⁻¹)	Tentative assignment
1097	C-O ^{30,32} (_{ν_{s, C-O})}
1179	$C-O-C^{30,32}(\nu_{s}, c_{-O-C})$
1330	C-CH ₃ ^{30,32} (δ _{5, C-CH₂})
1350	D band (GO) ²⁹
1420	C-H ^{30,32} (β _{C-H})
1590	G band (GO) ²⁹
1654	C=O ^{30,32} (<i>v</i> _{s, C=O})
1747	$COOH^{30,32}$ ($\nu_{s, COOH}$)
$\nu_{\rm s}$ symmetric stretching mode, β bending mode, $\delta_{\rm s}$ symmetric deformation mode	

spectrum contains pronounced Raman bands around 1350 cm⁻¹ and 1590 cm⁻¹ correlating with the D and G band positions of GO–COOH²⁹, confirming that this region contains the GO–COOH lattice. The TERS spectrum at location 2 is similar to that of location 1; although, the signal intensity is around 20 times stronger. However, the TERS spectrum at location 3 shows a very sharp and asymmetric D band and an additional Raman band at 1330 cm⁻¹ overlaps with the D band, which could be attributed to the symmetric deformation mode ($\delta_{s, C-CH_3}$) of C–CH₃^{30,31}. In the TERS spectrum for location 4, two new Raman bands are observed at 1654 cm⁻¹ and 1420 cm⁻¹, which

could be assigned to the symmetric stretching mode ($\nu_{s, C=O}$) of C=O³² and bending mode (β_{C-H}) of C-H³⁰, respectively.

To investigate the distribution of functional groups on the GO-COOH surface, we generated TERS intensity maps of the Raman bands corresponding to six different functional groups (Table 1), which are presented in Supplementary Fig. 4a-f. In these maps, a heterogeneous distribution of different functional groups is observed on the GO-COOH surface (see Supplementary Note 2 for details). TERS maps of the 1097 cm^{-1} , 1179 cm^{-1} , 1654 cm^{-1} , and 1747 cm^{-1} bands shown in Supplementary Fig. 4a, b, e, f, respectively indicate that the C-O, C-O-C, C=O and COOH groups are present only at a few locations on the GO-COOH surface. Furthermore, these functional groups are distributed over the entire GO-COOH surface. These experimental observations cannot be explained by the Lerf-Klinowski theoretical model of GO, according to which, COOH should be located on the edges, whereas other groups, such as C=O and C-O-C should be distributed over the entire surface³³. We propose that the discrepancy between the Lerf-Klinowski theoretical model and our experimental results could be caused by the two-step method used in the preparation of the GO-COOH sample, whereby only a fraction of the C-O and C-O-C groups obtained on the GO flake in the first step are converted to COOH in the second step³⁴. Further understanding of this discrepancy would require a detailed analysis of the sample preparation process and its influence on the lattice structure of GO-COOH, which is beyond the scope of this study. From the

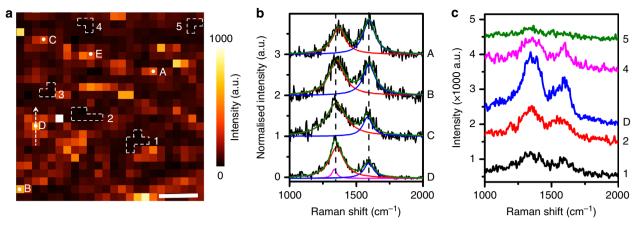


Fig. 3 Visualising structural defects and functional groups on the GO-COOH surface at the nanoscale. **a** Zoomed-in TERS map of D band intensity from $250 \times 250 \text{ nm}^2$ area marked with a dashed square in Fig. 2c. Scale bar: 50 nm. **b** TERS spectra measured at the locations marked as A-D in (**a**). **c** TERS spectra measured at five pixels along the arrow marked across location D in (**a**). Note that the area of the TERS map in (**a**) is only around 0.5 % of the diffraction limited laser spot area for our confocal Raman microscope

TERS maps, the area percentages of the C-O, C-O-C, C=O and COOH groups on GO-COOH surface are calculated to be 1.3%, 1.0%, 0.5% and 0.6%, respectively. However, in the TERS maps of the 1330 cm⁻¹ and 1420 cm⁻¹ bands (Supplementary Fig. 4c, d) a stronger TERS intensity is observed at a relatively larger number of areas on the GO-COOH flakes, indicating that the C-CH₃ and C-H groups have a more dominant presence on the GO-COOH surface. Furthermore, the TERS maps in Supplementary Fig. 4 indicate that the different functional groups are distributed differently over the GO-COOH surface. In particular, in the TERS map of 1330 cm^{-1} band intensity in Supplementary Fig. 4c, C-CH₃ groups show a clearly higher enhancement at the steps and edges of the GO-COOH flakes, whereas the TERS maps for other functional groups indicate a random distribution over the entire GO-COOH surface (see Supplementary Note 2). On the GO-COOH sample used in this study, the functional groups having Raman polarisability tensors aligned with the axial polarisation of the tip-enhanced near-field are expected to experience a much higher signal enhancement due to the gapmode TERS surface selection rule^{16,26-28}, rendering them more visible. This suggests that compared to the other functional groups, C-CH₃ groups likely bind to the steps and edges of GO-COOH in a more well-defined orientation with their Raman polarisability tensor aligned with the near-field axial polarisation.

The spatial resolution of the TERS map in Fig. 2c can be calculated by fitting a line profile across a sharp feature with a Gaussian curve³⁵⁻³⁷. An example of this is shown in Fig. 2f, where the spatial resolution of the TERS map is calculated to be around 10.1 nm, from the FWHM of the Gaussian curve fitted to the line profile. In order to obtain a statistical estimate, we performed Gaussian fitting of 5 line profiles from different locations in the TERS map and calculated the average spatial resolution to be 10.5 ± 1.7 nm (Supplementary Fig. 5). Further discussion of the obtained spatial resolution and comparison with previous studies^{28,38-41} is presented in Supplementary Note 3. To demonstrate the reproducibility of nanoscale spatial resolution we conducted TERS mapping on the same GO-COOH flakes using three different TERS tips, the results of which are presented in Supplementary Fig. 8. All three TERS tips provided a similar spatial resolution in the TERS maps (comparable to the step size) and displayed reproducibility of chemical imaging as discussed in Supplementary Note 4.

Visualising structural defects and functional groups. To demonstrate the capability of TERS for providing rich chemical

information with a high spatial resolution, we present further analysis of the TERS map shown in Fig. 2c. A zoomed-in TERS map of D band intensity from the area marked by the dotted square in Fig. 2c is presented in Fig. 3a. In this TERS map, areas with high TERS D band intensity (>900 counts) are observed interspersed with several areas having a low signal intensity (30-40 counts), which is close to the noise level. Averaged TERS spectra from five low intensity areas marked using white dashed lines in Fig. 3a are presented in Supplementary Fig. 9 (see Supplementary Note 5 for details). In the averaged spectra measured at these locations, no D or G bands are observed indicating a discontinuity in the GO-COOH lattice. Previously, using TEM, Erickson et al. have shown that few-layer GO-COOH flakes are composed of functionalised areas with large vacancies inside¹¹. Therefore, these low signal intensity areas with unobservable D and G bands can be attributed to the vacancies present within the GO lattice. From inspection of individual TERS spectra, the size of these vacancies was determined to be $100-600 \text{ nm}^2$ (1-6 pixels) correlating well with the size reported by Erickson et al.¹¹.

Furthermore, in the TERS map, remarkably different spectra were observed at different areas with high signal intensity. 38 TERS spectra from areas of high-signal intensity in Fig. 3a are plotted in Supplementary Fig. 10. The ratio of D to G band intensity (I_D/I_G) is commonly used to quantify defect density in graphene⁴² and GO¹⁰. The 38 spectra can be broadly classified into two groups. The first group (Supplementary Fig. 10b- e) shows only the D and G bands, whereas the second group (Supplementary Fig. 10f) shows Raman bands from multiple functional groups spectrally overlapping with the D and G bands. Based on the I_D/I_G ratio, the TERS spectra in the first group can be further divided into four types: (1) $I_D/I_G < 1$ (Supplementary Fig. 10b), (2) $I_D/I_G \approx 1$ (Supplementary Fig. 10c), (3) $1 < I_D/I_G <$ 1.5 (Supplementary Fig. 10d), (4) $I_D/I_G > 1.5$ (Supplementary Fig. 10e). Four representative TERS spectra from Supplementary Fig. 10b- e measured at locations A-D in Fig. 3a, are presented in Fig. 3b. In the TERS spectra from A-C, the shape of D and G bands at 1350 cm^{-1} and 1590 cm^{-1} resembles the TERS spectrum at position 2 in Fig. 2e, which can be fitted using a pair of Lorentzian curves. Therefore, spectra measured at A-C can be assigned to GO-COOH areas with increasing density of defects $(A < B < C)^{10}$. The TERS spectrum at location D also shows two Raman bands; however, the Raman band at 1350 cm^{-1} is unusually sharp and more intense than that of Raman band at 1590 cm^{-1} and cannot be fitted using a single Lorentzian curve. This spectrum is similar to the TERS spectrum at position

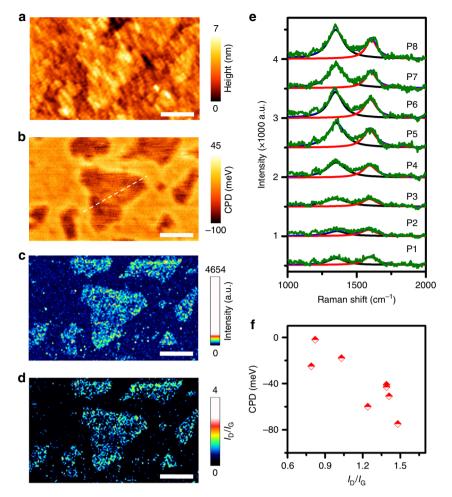


Fig. 4 In situ topographical, chemical and electrical nanoscopy of GO-COOH. **a** Topography map of a few-layer GO-COOH sample obtained whilst TERS mapping. **b** CPD map, **c** TERS map of D band (1350 cm⁻¹) intensity and **d** I_D/I_G ratio map measured in the region shown in (**a**). Area: $2.5 \times 1.7 \,\mu\text{m}^2$. Step size: 16.7 nm. Integration time: 75 ms/pixel. **e** Averaged TERS spectra from the 8 different locations marked as P1-P8 in Supplementary Fig. 15 fitted with two Lorentzian curves. **f** Plot of CPD as a function of I_D/I_G ratio calculated from the averaged TERS spectra shown in (**e**) showing an inverse correlation between local defect density and Fermi level on the GO-COOH surface. All scale bars: 500 nm

3 shown in Fig. 2e, and therefore, can be assigned to a combination of the D band and $\delta_{s, C-CH_3}$ vibrational mode of $C-CH_3^{30,31}$. In contrast to TERS spectra of the first group (showing only D and G band features), TERS spectra of the second group (location E) shown in Supplementary Fig. 10f exhibit a number of different Raman bands. The positions of these Raman bands correlate well with the bands listed in Table 1 indicating that they represent different functional groups present on the GO-COOH surface. It should be noted that the proportions of the two groups of TERS spectra are in first group (positions A–D), whereas only 15% spectra are in the second group (position E) as indicated by the histogram shown in Supplementary Fig. 11. This indicates that the surface of this GO-COOH sample is only partially functionalised.

Interestingly, the TERS spectra measured from pixels in the immediate vicinity of positions A–E were found to be similar to each other. The TERS spectra from five pixels along the line marked across position D in Fig. 3a are plotted in Fig. 3c. In order to eliminate the possibility of artefacts caused by occasional contamination of TERS tip, these lines were chosen normal to the scan direction of TERS mapping. All TERS spectra shown in Fig. 3c exhibit typical broad Raman bands at ~1350 cm⁻¹ and ~1590 cm⁻¹, the positions of D and G bands. I_D/I_G is calculated

to be 1.8 for position D, whereas 20 nm away from position D it decreases to 1.2 (Supplementary Fig. 12). Therefore, the pixels across position D represent a cluster of defects with the highest defect density at the center and decreasing away from it. This reveals heterogeneity in the distribution of defects on the GO-COOH surface within the length-scale of a few nanometres. To the best of our knowledge, such high-resolution chemical information from a GO-COOH surface cannot be obtained non-destructively using any other analytical technique.

In situ topographical chemical and electrical nanoscopy. To study the variation of local electronic properties due to functional groups present on the GO–COOH surface, we carried out in situ topographical, chemical and electrical nanoscopy by performing KPFM & TERS measurements consecutively in the same sample area using the same Au coated TERS tip. Topographic image of a few-layer GO–COOH sample obtained whilst TERS mapping is shown in Fig. 4a. This topography map has a poor spatial resolution because of the relatively large size of the TERS tip-apex, but shows several few-layer GO–COOH flakes (1–2 nm thickness) along with a $\approx 1 \,\mu$ m size GO–COOH flake at the centre. A contact potential difference (CPD) map measured in this region using KPFM is shown in Fig. 4b. In this map, 200–500 nm size areas with a significantly different CPD are observed within the GO-COOH flakes. A line profile across the central GO-COOH flake in Fig. 4b presented in Supplementary Fig. 13 reveals CPD variation of -87 to 34 meV. Since the Fermi level of gold ($E_{f_{Au}}$) is 5.1 eV^{43,44} and the Fermi level of GO-COOH ($E_{f_{GO}}$) is equal to $E_{f_{Au}} + CPD$, $E_{f_{GO}}$ is found to vary from 5.01 to 5.13 eV across the central GO-COOH flake. This Fermi level is similar to the previously reported value of 5 eV of GO⁴⁵, but is higher than the value of 4.5 eV for graphene^{43,44}.

After KPFM measurements, we performed in situ TERS measurement in the same sample area shown in Fig. 4a, b. TERS maps of D and G band intensities at 1350 cm⁻¹ and 1590 cm⁻¹ are shown in Fig. 4c and Supplementary Fig. 14, respectively. The TERS $I_{\rm D}/I_{\rm G}$ ratio image is shown in Fig. 4d. In these TERS maps, a heterogeneous distribution of D and G band intensity is observed across all GO-COOH flakes. The TERS intensity variations indicate local differences in surface chemistry, which is expected to cause variation of surface potential over GO-COOH flakes. To probe correlation of local defect density and surface potential, we calculated the I_D/I_G ratio from the TERS spectra averaged over an area of 0.012 µm² (50 pixels) from 8 different locations on the GO-COOH flake as shown in Supplementary Note 6. The averaged TERS spectra are shown in Fig. 4e. The I_D/I_G ratio calculated from the averaging of TERS spectra over a large area should be similar to the far-field measurement, which is conventionally used to characterise the defect level in graphene⁴² and graphene oxide¹⁰, thereby allowing comparison of the CPD with the local defect density of the GO-COOH flake despite different spatial resolutions of the TERS and CPD maps. The I_D/I_G ratio was calculated by fitting the averaged TERS spectra using two Lorentzian curves at 1350 cm⁻¹ and 1590 cm⁻¹. A plot of the CPD versus I_D/I_G ratio from the eight areas is shown in Fig. 4f indicating an inverse correlation between the two parameters. A high CPD (≈ -20 meV, Fermi level = 5.08 eV) is observed in the areas (P1, P2) with $I_D/I_G < 1$, whereas a low CPD ($\approx -80 \text{ meV}$, Fermi level = 5.02 eV) is observed in the areas (P6, P7, P8) with $I_{\rm D}/I_{\rm G}$ > 1.3. Interestingly, this inverse correlation of Fermi level with I_D/I_G ratio on the GO-COOH surface observed in this work is different to that of a-beam irradiated graphene, where the Fermi level has been reported to increase from 4.5 to 4.9 eV with increasing density of defects⁴⁴. This difference could result from the presence of different kinds of defects on GO-COOH surface. These results demonstrate that the in situ KPFM and TERS method allows a direct correlation of the Fermi level of a GO-COOH flake to the defect density at the nanoscale, thereby significantly improving the accuracy of GO-COOH surface characterisation.

In summary, we have demonstrated nanoscale TERS mapping of structural defects and functional groups present on a GO-COOH surface with an unprecedented spatial resolution of ≈10 nm. Furthermore, we have pushed nanoscale surface characterisation a step forwards by demonstrating in situ topographical, electrical and chemical nanoscopy of a GO-COOH sample by combining TERS with KPFM. This in situ multi-parameter measurement methodology greatly extends the capability of TERS allowing a direct correlation of local topography, chemical composition and electronic properties at the nanoscale not only in 2D materials but on almost any sample surface. In particular, we expect this work to open up the possibility of optimising optoelectronic devices based on novel 2D materials such as graphene, GO, single-layer MoS₂ and others via non-destructive, simultaneous and nanoscale multi-parameter characterisation of their surface properties even under operational conditions.

Methods

TERS measurements. TERS measurements were performed on a NanoRaman system consisting of an atomic force microscope (OmegaScope, formerly AIST-NT, now HORIBA Scientific) combined with a Raman spectrometer (XploRA, HOR-IBA Scientific, France) in side illumination geometry as schematically shown in Fig. 1. A 638 nm excitation laser was focussed on the sample at an angle of 60° using a ×100, 0.7 NA, 20 mm long working distance objective lens (Mitutoyo, Japan). TERS spectra were measured using a spectrometer grating with 600 lines/ mm and an electron-multiplying charged-coupled device detector (Andor, Ireland) with a laser power of 100 μ W at the sample.

For fast and efficient TERS mapping, TERS measurements were conducted in SpecTop^{**} TERS mapping mode in which TERS spectrum at a particular pixel in the TERS map is measured when the tip is in direct contact with the surface, with a typical interaction force of 2–10 nN and integration time of 0.05 s-0.5 s. Transition between the pixels of the TERS map is performed in semi-contact mode, which preserves both the sharpness and plasmonic enhancement of the tip leilminating lateral forces that might otherwise sweep aside or pick up loosely attached contaminants from the sample surface. All TERS measurements were performed using Au coated AFM TERS tips (k = 7 N/m, f = 150 kHz, formerly AIST-NT, now HORIBA Scientific).

Sample preparation. GO-COOH samples for TERS measurements were prepared by spin-coating GO-COOH (ACS Material, USA) on a Au coated glass substrate.

Data availability. The data supporting the findings of this study are available from the corresponding authors upon request.

Received: 22 October 2017 Accepted: 20 June 2018 Published online: 23 July 2018

References

- Liu, S., Shen, B., Niu, Y. & Xu, M. Fabrication of WS₂-nanoflowers@rGO composite as an anode material for enhanced electrode performance in lithium-ion batteries. *J. Colloid Interface Sci.* 488, 20–25 (2017).
- Sanli, L. I., Yarar, B., Bayram, V. & Gursel, S. A. Electrosprayed catalyst layers based on graphene-carbon black hybrids for the next-generation fuel cell electrodes. *J. Mater. Sci.* 52, 2091–2102 (2017).
- Yin, Z. et al. Electrochemical deposition of ZnO nanorods on transparent reduced graphene oxide electrodes for hybrid solar cells. *Small* 6, 307–312 (2010).
- Zhang, L., Xia, J., Zhao, Q., Liu, L. & Zhang, Z. Functional graphene oxide as a nanocarrier for controlled loading and targeted delivery of mixed anticancer drugs. *Small* 6, 537–544 (2010).
- Dreyer, D. R., Park, S., Bielawski, C. W. & Ruoff, R. S. The chemistry of graphene oxide. *Chem. Soc. Rev.* 39, 228–240 (2010).
- Haubner, K. et al. The route to functional graphene oxide. *Chem. Phys. Chem.* 11, 2131–2139 (2010).
- Wang, D.-W., Wu, K.-H., Gentle, I. R. & Lu, G. Q. Anodic chlorine/nitrogen co-doping of reduced graphene oxide films at room temperature. *Carbon* 50, 3333–3341 (2012).
- Ling, K. et al. A self-assembling RNA aptamer-based graphene oxide sensor for the turn-on detection of theophylline in serum. *Biosens. Bioelectron.* 86, 8–13 (2016).
- Wang, H. et al. Graphene oxide-peptide conjugate as an intracellular protease sensor for Caspase-3 activation imaging in live cells. *Ang. Chem. Int. Ed.* 50, 7065–7069 (2011).
- Yang, D. et al. Chemical analysis of graphene oxide films after heat and chemical treatments by X-ray photoelectron and Micro-Raman spectroscopy. *Carbon* 47, 145–152 (2009).
- Erickson, K. et al. Determination of the local chemical structure of graphene oxide and reduced graphene oxide. *Adv. Mater.* 22, 4467–4472 (2010).
- Kumar, N., Stephanidis, B., Zenobi, R., Wain, A. J. & Roy, D. Nanoscale mapping of catalytic activity using tip-enhanced Raman spectroscopy. *Nanoscale* 7, 7133–7137 (2015).
- Stadler, J., Schmid, T. & Zenobi, R. Nanoscale chemical imaging of single-layer graphene. ACS Nano 5, 8442–8448 (2011).
- 14. Su, W. & Roy, D. Visualising graphene edge using tip enhanced Raman spectroscopy. J. Vac. Sci. Technol. B **31**, 041808 (2013).
- Yeo, B. S., Amstad, E., Schmid, T., Stadler, J. & Zenobi, R. Nanoscale probing of a polymer-blend thin film with tip-enhanced Raman spectroscopy. *Small* 5, 952–960 (2009).

- Zhang, R. et al. Chemical mapping of a single molecule by plasmon-enhanced Raman scattering. *Nature* 498, 82–86 (2013).
- Saito, Y., Verma, P., Masui, K., Inouye, Y. & Kawata, S. Nano-scale analysis of graphene layers by tip-enhanced near-field Raman spectroscopy. *J. Raman Spectrosc.* 40, 1434–1440 (2009).
- Snitka, V., Rodrigues, R. D. & Lendraitis, V. Novel gold cantilever for nano-Raman spectroscopy of graphene. *Microelectron. Eng.* 88, 2759–2762 (2011).
- Ghislandi, M., Hoffmann, G. G., Tkalya, E., Xue, L. & De With, G. Tipenhanced Raman spectroscopy and mapping of graphene sheets. *Appl. Spectrosc. Rev.* 47, 371–381 (2012).
- Beams, R., Cancado, L. G., Jorio, A., Vamivakas, A. N. & Novotny, L. Tipenhanced Raman mapping of local strain in graphene. *Nanotechnology* 26, 175702 (2015).
- Pashaee, F., Sharifi, F., Fanchini, G. & Lagugne-Labarthet, F. Tip-enhanced Raman spectroscopy of graphene-like and graphitic platelets on ultraflat gold nanoplates. *Phys. Chem. Chem. Phys.* 17, 21315–21322 (2015).
- Rickman, R. H. & Dunstan, P. R. Enhancement of lattice defect signatures in graphene and ultrathin graphite using tip-enhanced Raman spectroscopy. *J. Raman Spectrosc.* 45, 15–21 (2014).
- Mignuzzi, S. et al. Probing individual point defects in graphene via near-field Raman scattering. Nanoscale 7, 19413–19418 (2015).
- 24. Berweger, S. et al. Optical nanocrystallography with tip-enhanced phonon Raman spectroscopy. *Nat. Nanotechnol.* **4**, 496–499 (2009).
- Ferrari, A. C. et al. Raman spectrum of graphene and graphene layers. *Phys. Rev. Lett.* 97, 187401 (2006).
- Jiang, S. et al. Distinguishing adjacent molecules on a surface using plasmonenhanced Raman scattering. *Nat. Nanotechnol.* 10, 865–869 (2015).
- Jiang, N. et al. Nanoscale chemical imaging of a dynamic molecular phase boundary with ultrahigh vacuum tip-enhanced Raman spectroscopy. *Nano Lett.* 16, 3898–3904 (2016).
- Shao, F., Mueller, V., Zhang, Y., Schluter, A. D. & Zenobi, R. Nanoscale chemical imaging of interfacial monolayers by tip-enhanced Raman spectroscopy. *Ang. Chem. Int. Ed.* 56, 9361–9366 (2017).
- Zhang, K., Kemp, K. C. & Chandra, V. Homogeneous anchoring of TiO₂ nanoparticles on graphene sheets for waste water treatment. *Mater. Lett.* 81, 127–130 (2012).
- Lewis, I. R. & Edwards, H. G. M. Handbook of Raman Spectroscopy. (Marcel Dekker, Inc, New York, 2001).
- Gomez-Navarro, C. et al. Electronic transport properties of individual chemically reduced graphene oxide sheets. *Nano Lett.* 7, 3499–3503 (2007).
- Smith, E. Dent, G. Modern Raman spectroscopy- a practical approach. 1st Ed, (John Wiley & Sons, Chichester, 2005).
- Lerf, A., He, H., Forster, M. & Klinowski, J. Structure of graphite oxide revisited. J. Phys. Chem. B 102, 4477–4482 (1998).
- 34. ACS Material. Carboxyl Graphene. *MATERIALS CATALOG* https://www.acsmaterial.com/carboxyl-graphene-962.html (2018).
- Roy, D., Wang, J. & Williams, C. Novel methodology for estimating the enhancement factor for tip-enhanced Raman spectroscopy. *J. Appl. Phys.* 105, 013530 (2009).
- Yano, T.-a., Verma, P., Saito, Y., Ichimura, T. & Kawata, S. Pressure-assisted tip-enhanced Raman imaging at a resolution of a few nanometres. *Nat. Photonics* 3, 473–477 (2009).
- Kumar, N. et al. Nanoscale chemical imaging of solid-liquid interfaces using tip-enhanced Raman spectroscopy. *Nanoscale* 10, 1815–1824 (2018).
- Su, W., Kumar, N., Dai, N. & Roy, D. Nanoscale mapping of intrinsic defects in single-layer graphene using tip-enhanced Raman spectroscopy. *Chem. Commun.* 52, 8227–8230 (2016).

- Su, W., Kumar, N., Mignuzzi, S., Crain, J. & Roy, D. Nanoscale mapping of excitonic processes in single-layer MoS₂ using tip-enhanced photoluminescence microscopy. *Nanoscale* 8, 10564–10569 (2016).
- Park, K.-D. et al. Hybrid tip-enhanced nanospectroscopy and nanoimaging of monolayer WSe₂ with local strain control. *Nano Lett.* 16, 2621–2627 (2016).
- 41. Muller, V. et al. Structural characterization of a covalent monolayer sheet obtained by two-dimensional polymerization at an air/water interface. *Ang. Chem. Int. Ed.* **56**, 15262–15266 (2017).
- Cancado, L. G. et al. Quantifying defects in graphene via Raman spectroscopy at different excitation energies. *Nano Lett.* 11, 3190–3196 (2011).
- Shi, Y. et al. Work function engineering of graphene electrode via chemical doping. ACS Nano 4, 2689–2694 (2010).
- 44. Kim, J.-H. et al. Work function engineering of single layer graphene by irradiation-induced defects. *Appl. Phys. Lett.* **103**, 171604 (2013).
- Li, S.-S., Tu, K.-H., Lin, C.-C., Chen, C.-W. & Chhowalla, M. Solutionprocessable graphene oxide as an efficient hole transport layer in polymer solar cells. ACS Nano 4, 3169–3174 (2010).

Acknowledgements

W.S. acknowledges financial support from Zhejiang Provincial Natural Science Foundation of China (Grant No. LY18F040006) and Natural Sciences Foundation of China (Grant No. 61306115). Dr Sebastian Wood from the National Physical Laboratory, UK is thanked for proofreading the manuscript.

Author contributions

A.K. and M.C. performed TERS experiments and collected the data. W.S. performed microchemical characterisation of GO–COOH. W.S. and N.K. analysed the data and wrote the manuscript. All authors contributed to the discussion of results.

Additional information

Supplementary Information accompanies this paper at https://doi.org/10.1038/s41467-018-05307-0.

Competing interests: The authors declare no competing interests.

Reprints and permission information is available online at http://npg.nature.com/ reprintsandpermissions/

Publisher's note: Springer Nature remains neutral with regard to jurisdictional claims in published maps and institutional affiliations.

Open Access This article is licensed under a Creative Commons Attribution 4.0 International License, which permits use, sharing, adaptation, distribution and reproduction in any medium or format, as long as you give appropriate credit to the original author(s) and the source, provide a link to the Creative Commons license, and indicate if changes were made. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in a credit line to the material. If material is not included in the article's Creative Commons license and your intended use is not permitted by statutory regulation or exceeds the permitted use, you will need to obtain permission directly from the copyright holder. To view a copy of this license, visit http://creativecommons.org/ licenses/by/4.0/.

© The Author(s) 2018