## ARTICLE

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# Metal-induced n<sup>+</sup>/n homojunction for ultrahigh electron mobility transistors

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#### Abstract

A self-organized n<sup>+</sup>/n homojunction is proposed to achieve ultrahigh performance of thin film transistors (TFTs) based on an amorphous (Zn,Ba)SnO<sub>3</sub> (ZBTO) semiconductor with sufficiently limited scattering centers. A deposited Al layer can induce a highly O-deficient (n<sup>+</sup>) interface layer in the back channel of a-ZBTO without damaging the front channel layer via the formation of a metal-oxide interlayer between the metal and back channel. The n<sup>+</sup> layer can significantly improve the field-effect mobility by providing a relatively high concentration of free electrons in the front n-channel ZBTO, where the scattering of carriers is already controlled. In comparison with a Ti layer, the Al metal layer is superior, as confirmed by first-principles density functional theory (DFT) calculations, due to the stronger metal-O bonds, which make it easier to form a metal oxide AlO<sub>x</sub> interlayer through the removal of oxygen from ZBTO. The field-effect mobility of a-ZBTO with an Al capping layer can reach 153.4 cm<sup>2</sup>/Vs, which is higher than that of the pristine device, i.e., 20.8 cm<sup>2</sup>/Vs. This result paves the way for the realization of a cost-effective method for implementing indium-free ZBTO devices in various applications, such as flat panel displays and large-area electronic circuits.

#### Introduction

Thin-film transistors (TFTs) with amorphous oxide semiconductors (AOSs) have been incorporated into various display applications owing to their many advantages, such as their relatively high mobility ( $\geq 10 \text{ cm}^2/\text{Vs}$ ), large-area uniformity, and low-temperature process capability<sup>1–3</sup>. However, next-generation displays, such as 8 K (8000 × 4000) active-matrix organic light-emitting diode (AM-OLED) displays, demand much higher device performance capabilities, which will allow the integration of gate driver circuitry into these display panels. Typically, a field-effect mobility value of 50 cm<sup>2</sup>/Vs or greater is required<sup>4,5</sup>.

A number of high-mobility semiconductor materials, such as InZnO (IZO), InSnZnO (ITZO), ZnSnO (ZTO),

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understandable because the metal layer acts as a catalyst, easily breaking and rearranging the weakened metaloxygen bonds in the channel during the postannealing process after the deposition of the metal layer. This phenomenon is referred to as the low-temperature MIC of the active layer. The final mechanism stems from an increase in the carrier concentration via electron injection from the metal into the active channel layer driven by the work function difference<sup>20,23</sup>. Based on band transport percolation or mobility edge models<sup>24–26</sup>, the increased mobility can easily be supported. However, despite much effort in this area, it has been difficult to stably implement TFTs based on an ultrahigh-mobility amorphous oxide semiconductor with a field-effect mobility value exceeding 100 cm<sup>2</sup>/Vs.

Fortunately, researchers have recently reported the theoretical design and experimental implementation of an amorphous (Zn,Ba)SnO<sub>3</sub> (a-ZBTO) semiconducting oxide, resulting in high-performance devices with excellent stability<sup>27</sup>. a-ZBTO tends to exhibit decreases in Odeficient states and band tail states as the BaSnO<sub>3</sub> (BTO) fraction increases, meaning that it can only achieve moderately high mobility due to the significantly reduced carrier concentration while maintaining the ultrahigh stability of TFTs. Hence, an ultrahigh field-effect mobility based on a-ZBTO can be expected if the free carrier concentration of a-ZBTO can be increased without compromising the (especially ultrastability-related) existing properties of a-ZBTO. For this reason, it is evident that a reliable strategy is to use a metal-capping approach on the a-ZBTO channel without a postannealing step to increase the carrier concentration without damaging the stability.

In this paper, we propose ultrahigh-mobility a-ZBTObased TFTs achieved via the formation of an  $n^+/n$ homojunction induced by a metal-capping layer considering the collaborative outcomes of experiments and DFT calculations. The  $n^+/n$  homojunction can be selforganized through the formation of a highly O-deficient  $(n^+)$  interface layer in the back channel layer induced by a metal-oxide interlayer even without postannealing, providing a significant increase in the free carrier concentration for the front channel. Most electrons injected from the  $n^+$  layer can be transported without serious scattering due to the low defect density of a-ZBTO. As a result, a device with an Al metal layer exhibits a remarkable field-effect mobility enhancement from 20.8 cm<sup>2</sup>/Vs to 153.4 cm<sup>2</sup>/Vs without significant degradation of other transfer parameters. TFTs with other active materials, in this case IZO and ZTO, were also examined for comparison. The mobility improvement of metal-capped ZBTO devices was far superior to those of IZO and ZTO TFTs. Additionally, to understand the major carrier transport mechanisms, we performed 1/f noise measurements on the ZBTO and Al-assisted ZBTO TFTs. Our results indicate that selecting a semiconducting channel material with few scattering centers, such as ZBTO, is a key factor during the fabrication of metalassisted TFTs capable of ultrahigh performance.

#### Materials and methods

#### **Device fabrication**

ZBTO devices were fabricated on heavily doped p-type Si substrates with a thermally grown 100-nm-thick SiO<sub>x</sub> dielectric material. Next, 20-nm-thick ZBTO active layers were deposited by cosputtering ZnSnO<sub>3</sub> and BaSnO<sub>3</sub> targets with RF powers of 80 and 10 W applied to the ZnSnO<sub>3</sub> and BaSnO<sub>3</sub> targets, respectively. The working pressure was 3 mTorr, while the Ar gas flow rate was fixed at 10 sccm. After depositing the active layer, thermal annealing was conducted at 350 °C for 1 h in air. Next, source-drain electrodes were formed by depositing a 150nm-thick aluminum (Al) film by thermal evaporation using a shadow mask. The resulting channel width (W) and length (L) were 800 and 200 µm, respectively. Finally, 20-nm-thick Al or Ti metal layers were selectively deposited between the source/drain electrodes with dimensions of W/L =  $1300/100 \,\mu m$  without subsequent annealing. In the device patterned by photolithography and a lift-off process, the channel W/L and metal layer W/ L were 100/50 and 800/35  $\mu$ m, respectively.

#### Measurement

Measurements and stability tests to determine the transfer characteristics were conducted using an HP 4156B semiconductor parameter analyzer. The field-effect mobility values were extracted from the saturation region with the drain voltage  $(V_{DS})$  fixed at 10 V. Cross-sectional transmission electron microscopy (TEM) specimens were prepared by in situ lift-out in a dual-beam focused ion beam system (FIB; FEI Helios NanoLabTM). Additionally, X-ray photoelectron spectroscopy (XPS) analyses were performed to study the chemical bonding states of tin cations and oxygen anions using a monochromatic Al Ka X-ray source. Before performing the XPS analyses, the surface of each film was sputtered with a low-energy Ar<sup>+</sup> ion beam (200 eV) for 60 s to eliminate any possible contamination. The peak position was calibrated with respect to the Zn 2p peak, for which the standard binding energy is centered at 1021.7 eV. The threshold voltage  $(V_{th})$  was defined by the gate voltage that induced a drain current of W/L  $\times$  10 nA at a V<sub>DS</sub> of 1 V. Flicker noise measurements were performed using a semiconductor device analyzer (Keysight Technologies B1500A).

#### Density functional theory (DFT) calculations

All DFT calculations were performed using the Vienna ab initio Simulation Package  $(VASP)^{28,29}$  with the

projected augmented wave (PAW) potentials<sup>30</sup>. To trace the diffusion of atoms in the interface region between the a-ZBTO and metal Al/Ti lavers, first-principles molecular dynamics (MD) calculations were employed. The final structure was then obtained using static DFT calculations. The exchange-correlation functional based on PBEsol (Perdew-Burke-Ernzerhof revised for solids)<sup>31</sup> was used, which is known to describe atomic structures well at a low computational cost. The kinetic energy cutoff for the plane-wave basis set was chosen to be 400 (600) eV for MD calculations (volume optimization). In addition, kspace integration was used with a  $\Gamma$ -point of  $(1 \times 1 \times 1)$  to generate an amorphous structure and with a mesh of  $(2 \times$  $2 \times 2$ ) to optimize the crystal structure. All atomic positions were fully relaxed until the Hellmann-Feynman force on each atom was within 0.01 eV  $Å^{-1}$ .

To model the interface structure between ZBTO and the metal, the superlattice-like structure used here consisted of a combined a-ZBTO and metal (a-Al/a-Ti) supercell created via the following sequential processes: (i) a-ZBTO (a-metal) based on a  $(2 \times 2 \times 8)$  a- $Z_{1-x}B_x$ TO with x = 0.312 supercell of 160 atoms (an a-Al/a-Ti supercell of 64 atoms) was independently generated by firstprinciples melt-quenching (MQ) MD calculations. Until the total energy converged, MQ-MD calculations for ZBTO (metal) were repeated several times for melting at 4000 (3000) K and quenching at rates from 1.0 K fs<sup>-1</sup> to 0.5 K fs<sup>-1</sup> to temperatures from approximately 4000 (3000) K to 10 K. (ii) After the a-ZBTO and a-metal structures were combined, simulation annealing was conducted from 300 K to 50 K ( $0.2 \text{ K fs}^{-1}$ ). This process was also repeated until the total energy converged. (iii) To study the diffusion of atoms at the interface while considering the bulk by fixing 43 atoms in the middle region of a-ZBTO, a finite temperature MD simulation was conducted for a long simulation time of 20 ns. The temperature was set at a higher temperature (750 K) than the experimental temperature (approximately room T) to accelerate the atomic diffusion in both the ZBTO and metal layers.

#### **Results and discussion**

#### Effect of the Al metal layer on the ZBTO thin-film layer

Figure 1a, b presents cross-sectional transmission electron microscopy (TEM) images of a-ZBTO films with Ti and Al metal layers, respectively. When a Ti-capping layer is used, the Ti layer is entirely oxidized (Fig. 1a), whereas with the Al-capping layer, a relatively thin oxide  $(AlO_x)$  interlayer forms between the Al and a-ZBTO (Fig. 1b). The X-ray photoelectron spectroscopy (XPS) depth



profiles of the atomic elements present in the stacked structures in Fig. 1c, d also show that the entire and thin oxidized regions in the Ti and Al metal layers, respectively, are consistent with the aforementioned TEM observations. This difference in oxidation can be easily understood. Because Ti readily oxidizes even upon simple exposure to  $air^{32}$ , the formation of  $TiO_x$  on top of a-ZBTO is likely to be spontaneously caused by both oxidation from the ambient environment and consumption of oxygen from the underlying ZBTO. On the other hand, an Al layer is known to form only a very thin and dense oxide on its surface after exposure to air, which prevents further internal oxidation<sup>33</sup>. It is thus reasonable to conclude that the AlO<sub>x</sub> interlayer between ZBTO and Al is formed by the consumption of oxygen from the underlying ZBTO.

Figure 2 shows the O 1 s and Sn 3d XPS spectra of the reference ZBTO and metal-assisted ZBTO films. While the spectra from the ZBTO films without a metal layer were collected from the surface of the films, the spectra from the ZBTO films with the metal layer were collected near the metal/semiconductor interface during depth

profiling. All the subpeaks are shown in Fig. S1. The O 1 s peaks were resolved into three subpeaks centered at 530.2 (O1), 531.2 (O2), and 532.7 eV (O3). The O1, O2 and O3 peaks represent the signals originating from metal-oxygen bonds, O deficiency, and hydrogen-related bonds, respectively<sup>34</sup>. Figure 2a, b shows that the portions of the O1 and O2 peaks for the reference ZBTO films and ZBTO films with a Ti layer are nearly identical. In contrast, the O2 peak portion increases when an Al metal layer is applied, from 21.5 to 30.3 %, as shown in Fig. 2b. The Gibbs free energies of formation ( $\Delta G_f$ ) of ZnO, SnO<sub>2</sub>, BaO, TiO<sub>2</sub> and Al<sub>2</sub>O<sub>3</sub> are -318.4, -520.5, -520.4, -889.1, and -1582.3 kJ/mol, respectively<sup>35</sup>. The lowest  $\Delta G_f$  of Al<sub>2</sub>O<sub>3</sub> means that Al has the strongest oxidation power compared to ZnO, SnO<sub>2</sub>, BaO, and TiO<sub>2</sub>. Therefore, an AlO<sub>x</sub> film most likely forms at the Al/ZBTO interface due to the removal of oxygen from the ZBTO layer, which in turn significantly induces O deficiency in the back channel region. On the other hand, the Sn 3d peaks were also resolved into three subpeaks, as shown in Fig. 2c, d. The subpeaks at 487.0, 486.3, and 484.8 eV can be assigned to signals originating from  $SnO_2$  ( $Sn^{4+}$ ), SnO



spectra of the reference ZBTO films. **d** Comparison of the subpeak area percentages for SnO<sub>2</sub> and SnO of the reference and metal-assisted ZBTO films from the Sn  $3d_{5/2}$  spectra.

 
 Table 1
 Hall effect measurement results of the reference and Al-capped ZBTO thin-films.

TFT	Bulk concentration	Sheet resistance
ZBTO	-7.87 E + 15	104 MΩ/□
ZBTO – with Al layer	-1.91 E + 16	5.2 MΩ/□

 $(\text{Sn}^{2+})$ , and metallic Sn  $(\text{Sn}^{0})$  bonds<sup>36</sup>. The increased portions of the Sn<sup>2+</sup> subpeak in the ZBTO films with the metal layer suggest that the transition of SnO<sub>2</sub> to SnO  $(\text{Sn}^{4+} + 2e^- \rightarrow \text{Sn}^{2+})$  may occur due to the electrons injected from the metal layer. Because the work functions of Al and Ti are similar at 4.28 and 4.33 eV<sup>37</sup>, respectively, the effects of the electron injection by the two metal layers would be nearly identical. This could be confirmed by the result that the increase degrees of Sn<sup>2+</sup>-related subpeaks in both films with the metal layers are nearly identical.

Note that the oxygen and tin subpeak portions in the bulk regions, as shown in Supplementary Fig. S2, are nearly identical regardless of the presence or absence of a metal layer. The above results indicate that the a-ZBTO channel with an Al metal layer consists of two regions, a conductive region on top with a relatively high concentration of free carriers induced by the O deficiency ( $n^+$  layer) and a relatively unaltered ZBTO bulk region (n layer) on the bottom, which can itself form a metal-induced  $n^+/n$  homojunction within the semiconductor layer.

Furthermore, the Hall effect measurement results showed that the carrier concentration of ZBTO thin-films was increased approximately 2 times by the Al metal layer, as listed in Table 1. This result clearly confirms the effect of the Al metal layer on the ZBTO thin-films. Additionally, the increased carrier concentration due to the Al metal layer ( $\sim 10^{16}$ /cm<sup>3</sup>) is not sufficient to form a degenerate channel, which is reported to occur through the injection of a large number of free-electron carriers<sup>38</sup>. Therefore, a degenerate channel would not be formed in the ZBTO channel.

#### **First-principles DFT calculations**

Although shown in the XPS analysis above, it is fundamental to theoretically verify whether, among Al and Ti, the Al metal layer better removes oxygen atoms from the back channel at the interface between a-ZBTO and the metal-capping layer. To do this, first-principles DFT calculations were applied to study the interface of a-ZBTO with a metal-capping layer. Here, the interface structure was modeled by a superlattice-like structure in which a  $(2 \times 2 \times 8)$  ZBTO supercell of 160 atoms and a metal (Al or Ti) supercell of 64 atoms were combined (see Fig. 3a, b). For comparison of the intermixing of atoms for the two cases of Al and Ti, a finite-temperature molecular dynamics (MD) simulation was conducted with a long



simulation time of 20 ns at a temperature of 750 K, which is higher than the experimental temperature, to accelerate the atomic diffusion in the ZBTO-metal structure (see the DFT section of Experimental Methods). Based on the initial (I) and final (F) structures before and after the finite-temperature MD simulation, the relative trends for the Al and Ti metals were compared by analyzing the integrated pair distribution function (*IPDF*) data before and after the finite-temperature MD simulation. Figure 3c–e shows the differences ( $\Delta IPDFs$ ) in the average data per atom for the IPDF between cations (Zn, Ba, Sn and metals such as Al and Ti) and O for the two cases calculated according to  $\Delta IPDF = IPDF(F) - IPDF(I)$ . Interestingly, the coordination numbers for all of the Zn, Ba, and Sn cations with O atoms were calculated and found to

cases is plotted according to  $\triangle IPDF = IPDFs(F) - IPDFs(I)$ . In **c**-**e**, the

vertical dotted lines are the cation-O bond lengths in the ZnO, BaO,

SnO<sub>2</sub>, Al<sub>2</sub>O<sub>3</sub>, and TiO<sub>2</sub> crystals.

be significantly reduced more by the Al metal-capping layer than by the Ti metal-capping layer, depending on the type of cation, similar to the experimental XPS results. In particular, the  $\triangle IPDFs$  vary slightly between the Al and Ti cases for Zn and Sn but dramatically for Ba (Fig. 3c-e). This means that the highly O-deficient n<sup>+</sup> layer at the interface of the Al-ZBTO back channel forms because the Al atoms strongly pull the O atoms from the ZBTO layer to the Al metal layer. The superiority of the Al metal can be easily understood to be due to the greater heat of formation of Al<sub>2</sub>O<sub>3</sub> than that of TiO<sub>2</sub>. Therefore, the role of the metal-capping layer in the formation of the O-deficient region of the a-ZBTO channel is well characterized in terms of the diffusion of atoms during the production of the experimental TFT device.

#### Electrical characteristics of the thin-film transistors

Figure 4a, b shows the device structure and an optical micrograph of the ZBTO TFTs with a metal layer. The metal layer was deposited onto the ZBTO film at half the length of the channel between the source and drain electrodes. The transfer curves of TFT devices based on the reference and metal-assisted ZBTO are shown in Fig. 4c-e, and the extracted transfer parameters are listed in Table 2. The extracted saturation mobility values with respect to the gate voltage are shown in Fig. 4f. As anticipated, the application of a Ti metal layer results in an increase in the field-effect mobility by approximately twofold ( $\sim 53.7 \text{ cm}^2/\text{Vs}$ ). This improved mobility of Tiassisted TFTs originates from the electrons injected from the metal layer, as confirmed by the increase in the Sn<sup>2+</sup>-related peak in the Sn3d XPS spectra of both metalassisted ZBTO films. Unlike the Ti capping case, which corresponds to the typical field-effect mobility increase due to the metal-capping method, a further enhancement is observed in the presence of an Al capping layer with a remarkable peak mobility outcome of 153.4 cm<sup>2</sup>/Vs. This indicates that in addition to the electron injection effect, the main cause of the significantly improved mobility of the Al-assisted TFTs is the increase in the effective carriers stemming from the highly conductive n<sup>+</sup> layer formed within the ZBTO layer.

For comparison, other conventional oxide semiconductor materials were evaluated using an Al metal layer. Figure 5a depicts a schematic diagram of the electron injection mechanism from the conduction region to the bulk region in the active layer. Generally, conventional oxide semiconductors, such as IZO and ZTO, contain numerous oxygen-related and disorder defects that act as electron scattering/trapping centers<sup>39</sup>. Therefore, for conventional oxide semiconductors, most of the injected electrons from the n<sup>+</sup> layer could be trapped by various defects in the n layer (front channel layer), thereby limiting the mobility improvement effect. In contrast, for ZBTO, oxygen-related and disorder defects are effectively reduced by the incorporation of  $BTO^{27}$ . It has been reported that incorporation of BTO significantly decreases the band tail states or oxygen-related defects. As a result of the decreased defect density in ZBTO, most of the generated or injected electrons from the n<sup>+</sup> layer can be transported without serious scattering/trapping<sup>27</sup>. Therefore, an ultrahigh field-effect mobility can be achieved when using ZBTO as an active layer. Figure 5b indeed indicates a lower mobility improvement due to the Al metal layer for the conventional oxide IZO and ZTO TFTs compared to the ZBTO TFTs. Individual transfer curves of metal-assisted conventional oxide TFTs are shown in Supplementary Fig. S3, and the extracted transfer parameters are listed in Supplementary Table S1. For comparison, Fig. 5c shows the field-effect mobility outcomes of all TFTs.

# Low-frequency noise (LFN) characterization of the thin-film transistors

To probe the prevailing carrier transport mechanisms, we measured the 1/f noise of the ZBTO and Al-assisted ZBTO TFTs at  $V_{DS} = 1.0$  V and plotted the current noise spectral densities  $(S_{ID})$  versus the frequency, as seen in the inset of Fig. 6a. The slopes for both devices under all bias conditions obey the typical 1/f noise theory<sup>40</sup>. Figure 6a shows the normalized current noise spectral density  $(S_{ID}/$  $I_D^2$ ) as a function of  $|V_{GS} - V_{TH}|$  at 80 Hz. The slope value of -1.5 for the ZBTO TFTs indicates that the electrons are scattered by coherent trap sites composed of bulk traps and interface traps, as illustrated in Fig. 6c. In contrast, the slope value of almost -2 for the Al-assisted ZBTO TFTs indicates that the electrons are scattered mainly by interface traps, as depicted in Fig.  $6d^{40-44}$ . As a result of the reduced bulk trap sites, the Al-assisted ZBTO TFTs show fairly outstanding transfer characteristics compared to the ZBTO TFTs. The improved charge transport may originate from the excess electrons induced by the Al metal layer, which could passivate the bulk defect states in the ZBTO channel through the Al assisting layer. For further confirmation of transport mechanisms, the slopes of input-referred noise  $(S_{VG})$  were compared (see Fig. 6b).

$$S_{VG} = S_{VFB} \left[ 1 + \alpha \mu_{eff} C_{ox} (V_{GS} - V_{TH}) \right]^2$$
(1)

Here,  $S_{VFB}$  is the flat-band voltage noise density,  $\alpha$  is the Coulomb scattering coefficient,  $\mu_{eff}$  is the effective mobility, and  $C_{ox}$  is the dielectric capacitance per unit area. The slope of the ZBTO TFTs shows a tendency similar to that



in Eq. (1), suggesting that the carrier transport is influenced by the mechanism of bulk trap scattering  $^{42}$ .

$$S_{VG} = \frac{q^2 \cdot \mathbf{k} \cdot \mathbf{T} \cdot \mathbf{N}_T}{C_{ox}^2 \cdot W \cdot L \cdot \gamma \cdot f}$$
(2)

In this equation, q is the electronic charge; kT is the thermal energy;  $N_T$  is the oxide trap density; W and L are the channel width and length, respectively;  $\gamma$  is the attenuation coefficient; and f is the frequency<sup>45</sup>. In contrast, the slope of the Al-assisted TFTs is zero, coinciding with Eq. (2), which corresponds to the interface-trap-dominant model.

#### Reliabilities of the thin-film transistors

Figure 7a shows the device reliability test results for the reference and metal-assisted ZBTO TFTs under negative bias stress (NBS) and positive bias stress (PBS) conducted with a fixed gate voltage  $V_{GS} = \pm 20$  V and a drain voltage  $V_{DS} = 0.1$  V for 3600 s. Additionally, the amounts of threshold voltage shift under NBS and PBS are listed in Table 3. Refer to the detailed transfer curves of ZBTO TFTs and metal-assisted ZBTO TFTs in Fig. S4. The threshold voltage shifts during the bias stress test generally depend on the amount of trapped charge in the front channel region near the semiconductor/gate dielectric interface<sup>46,47</sup>. Despite the conductive region

Table 2Extracted transfer parameters of the reference,Ti-assisted, and Al-assisted ZBTO TFTs.

ZBTO TFT	$\mu_{FE}~(cm^2/Vs)$	S.S. (V/dec)	V <sub>th</sub> (V)	I <sub>on</sub> /I <sub>off</sub>
reference	20.8 ± 2.70	$0.32 \pm 0.06$	$-0.05 \pm 0.29$	~109
with Ti layer	53.7 ± 8.60	$0.44 \pm 0.08$	$-1.85 \pm 0.42$	~109
with Al layer	153 ± 16.6	$0.37\pm0.07$	$-0.77 \pm 0.35$	~10 <sup>9</sup>

induced in the back channel due to the formation of the  $AlO_x$  interlayer, the metal-assisted ZBTO TFTs exhibit relatively stable behavior, comparable to the reference devices. This suggests that the front channel is relatively less affected by the presence of the metal layer, as predicted from the aforementioned thin-film characterization results (e.g., XPS).

The long-term stability of the Al-assisted ZBTO TFTs was evaluated, as shown in Fig. 7b. Figure 7c shows the extracted transfer parameters as a function of the air exposure time. The electrical characteristics are well preserved over a period of 90 days under ambient air. The device performance enhanced by the application of the metal layer does not significantly degrade, which suggests that stable operation of the TFTs is possible with good long-term operational reliability.

For application to high-resolution displays and highly integrated circuits, the corresponding TFT size should be reduced. Accordingly, the possibility of downscaling the Al-assisted ZBTO TFTs using photolithography and a lift-off process was investigated. The channel layer width/ length of the downscaled (DS) device was reduced from 800/200 to  $100/50 \,\mu$ m. Supplementary Fig. S5 exhibits the transfer curves of DS-ZBTO TFTs with and without an Al







metal layer. It was found that the field-effect mobility of the Al-assisted DS-ZBTO TFTs was improved to more than ~120 cm<sup>2</sup>/Vs without large changes in the other transfer parameters. Although TFT fabrication by the lithography process requires further optimization, such as of the area ratio of the metal layer to the channel, the aforementioned effect of an improved mobility due to the metal layer was clearly confirmed for TFTs with a small feature size.

Finally, the extracted field-effect mobility and channel length (device scale) in previously reported works and this work are depicted in Fig. 8. Compared to other works, this work represents remarkably improved mobility ( $\sim 153 \text{ cm}^2/\text{Vs}$ ) without a stabilization or postannealing process. Moreover, most of the previous reports that adopted the metal capping process for enhancing the field-effect mobility were limited to TFTs with relatively long channel lengths. In the case of scaled-down TFTs, only a few

papers have reported the metal capping process, with relatively low field-effect mobilities (>50 cm<sup>2</sup>/Vs). However, our work shows greatly enhanced field-effect mobility (~120 cm<sup>2</sup>/Vs) even for scaled-down TFTs (L = 50  $\mu$ m). Consequently, we suggest that high-performance a-ZBTO TFTs with Al capping layers could be good candidates for future high-resolution display applications.

#### Conclusions

In summary, we investigated the effects of metalcapping layers on ZBTO films and the associated thinfilm device characteristics. The application of a highly oxidizing metal, such as Al, to ZBTO results in the creation of an AlO<sub>x</sub> interlayer, accompanied by the formation of an oxygen-deficient interfacial (n<sup>+</sup>) layer in the n-channel ZBTO film underneath, leading to an n<sup>+</sup>/n homojunction structure. This results in only an increase in the free carrier concentration in the front ZBTO



Table 3  $\Delta V_{TH}$  values of the reference, Ti-assisted, and Al-assisted ZBTO TFTs after NBS and PBS tests.

NBS			
TFT	ZBTO	ZBTO with Ti layer	ZBTO with Al layer
$\Delta$ V <sub>th</sub> (V)	-0.12	-0.46	-0.13
PBS			
TFT	ZBTO	ZBTO with Ti layer	ZBTO with Al layer
$\Delta$ V <sub>th</sub> (V)	0.05	0.36	0.09

channel layer with sufficiently limited scattering centers, remarkably enhancing the field-effect mobility of the corresponding TFTs. In addition, first-principles DFT calculations were utilized to determine the oxidation tendency of  $Al_2O_3$  relative to that of TiO<sub>2</sub>, which is the driving force behind the formation of the highly conductive n<sup>+</sup> layer in the ZBTO layer. As a result, the Alassisted ZBTO TFTs achieved a noticeably enhanced field-effect mobility of 153.4 cm<sup>2</sup>/Vs, far superior to those of other Al-assisted conventional oxide TFTs. From 1/f noise measurements and characterization of ZBTO and Al-assisted ZBTO TFTs, we found that the carrier transport of ZBTO TFTs with the Al assisting layer is enhanced due to the reduced bulk trap sites. The effect of the mobility improvement was also confirmed in smallsized ZBTO TFTs with an Al layer fabricated by lithography. Moreover, reliability tests indicate that the front channel of the metal-assisted devices is relatively less affected. These results suggest that the formation of a metal-induced oxide  $n^+/n$  homojunction for ultrahigh



respectively.

#### electron mobility transistors could be a cost-effective and useful route to replace current low-temperature poly-Si (LTPS) TFT technology.

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#### **Conflict of interest**

The authors declare that they have no conflict of interest.

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