#### NOTES

# Comments on Paper by Nemoto *et al.* Concerning Self Diffusion Coefficient, $D_s$ , and Zero Shear Viscosity, $\eta_0$ of Polystyrene in Dibutylphthalate

## Raymond F. BOYER

Michigan Molecular Institute, 1910 West St., Andrews Road, Midland, Michigan 48640, U.S.A.

(Received August 14, 1989)

KEY WORDS Self Diffusion / Steady Viscosity / Dibutylphthalate / Atactic
Polystyrene / Solutions / Regression Analysis / Liquid–Liquid Transition /
Arrhenius Plots / T<sub>u</sub> /

Nemoto *et al.*<sup>1</sup> obtained data on two narrow distribution at-polystyrenes,  $\bar{M}_n = 43300$  and  $\bar{M}_n = 335000$  at concentration of 40 and 50% by weight in dibutylphthalate at temperatures from 0 to 120°C. Data were presented in both tabular and graphical form. They concluded correctly that these data did not follow a simple Arrhenius plot for either  $\log D_s - T^{-1}$  or  $\log \eta_0 - T^{-1}$ . Then they fitted these data by the WLF procedure.<sup>2</sup> They concluded that this free volume approach yielded a correct description of their results.

Visual inspection of their plotted data, their Figures 3a and 3b specifically, with assistance from a straight edge, suggested to us an opposite conclusion, namely, that all experimental data for  $D_s$  and  $\eta_0$  followed a double or possibly a triple Arrhenius plot whose intersections implied liquid state transitions, one of which is  $T_{ll}$  lowered from its bulk state value by diluent.  $T_{ll}$  is the intramolecular liquid– liquid transition normally found in atactic polymers and copolymers at  $(1.20\pm0.05)$  $T_g(K)$ ; the other is an intramolecular event designated  $T_{ld}$ .<sup>6</sup>

As we have noted on several occasions,<sup>3,5-7</sup> Frenkel and his colleagues postulate that  $T_{ll}$ arises from intermolecular segment-segment associations which are thermally reversible and shear reversible. The source of these associations is the same intermolecular attraction that holds simple liquids together. These associations lead to a three dimensional physical network which inhibits but does not prevent flow. We suggest<sup>6</sup> that the sharp decrease in enthalpy of activation for  $\eta_0$  across  $T_{u}$  results from the disappearance of these physical networks.

Figure 1 is a plot of  $\log D_s$  and  $\log \eta_0$  against 1000/T(K) for 40% by weight of at-PS, 43900 in dibutylphthalate. Slope changes are indicated by vertical lines which occur at about the same values of 1000/T by both methods. A similar plot was prepared for at-PS, 335000 at 50% concentration, using  $D_s$  data, with one sharp slope change at 362 K (not shown).

### EXPERIMENTAL BACKGROUND

Since 1966, when it was first suggested that  $\log \eta_0 - T^{-1}$  plots followed a double Arrhenius pattern with a slope change at  $T_{ll}$ ,<sup>8</sup> we have examined various bodies of such published data as a function of molecular weight for a given polymer<sup>9,10</sup>; polymer types<sup>6,7</sup>; and more rarely polymer concentrations in polymer–diluent system.<sup>11</sup>

The number of intersections found in  $\eta_0$ 





**Figure 1.** Self diffusion,  $D_s$ , and melt viscosity,  $\eta_0$ , triple Arrhenius plots for 40% by weight at -PS,  $\overline{M}_n = 43,300$  in dibutylphthalate from data of Nemoto *et al.*<sup>1</sup> We have added the  $T_{ll}$  and  $T_{l\phi}$  designations.

data for the liquid state depends on the polymer, the range in temperature, the number of data points, and the accuracy of the data. At-PS under ideal conditions indicates three transitions:  $T_{ll}$ ,  $T_{l\rho}$ , and  $T_{l\phi}$  in order of ascending temperature.  $T_{ll}$  increases with  $\bar{M}_n$ , finally leveling off at  $T_{ll}(\infty)$ .  $T_{l\rho}$  is independent of  $\bar{M}_n$ at a constant value =  $T_{ll}(\infty)$ .  $T_{l\phi}$  is also independent of  $\bar{M}_n$  with a value of 190 ± 5 K in PS, but more generally at  $T_{ll}$  + 30—50 K.  $T_{l\rho}$  is most easily found at  $\bar{M}_n < 2000$  and hence should not be located in any data discussed in the present study. All three of these loci are found in Figure 1 of ref 6. An earlier plot<sup>5</sup> designated  $T_{l\phi}$  as  $T_{l\rho}$ .

Data were of 3 types: tabulated, graphical and both as in the present case.<sup>1</sup> Tabulated data have been examined by regression analysis with residuals,<sup>9,10</sup> by point to point first derivatives<sup>9</sup>; and by the Solc automatic intersection method which determines the number of linear segments and their coordinates in a given body of data which will minimize the standard deviation.<sup>12</sup> In the present instance, the number of data point -9 for three lines with FA-40 and FA-50 is borderline for use of computerized intersection method discussed in ref 12.

Actually, visual inspection of plotted data can be quite accurate, especially with high  $T_{ul}$ polymers. The change in enthalpy of activation across  $T_{ul}$  in bulk at-PS drops from *ca*. 75 kcal mol<sup>-1</sup> at  $T < T_{ul}$  to *ca*. 43 kcal mol<sup>-1</sup> just above  $T_{ul}$  with no marked dependence on molecular weight.<sup>9</sup> This same study covering values of  $\overline{M}_n$  from 3500 to 390000 showed no evidence for a change in the behavior of  $T_{ul}$ based on intersections on going through the entanglement molecular weight,  $M_c$  (Figure 11 of ref 9). This is consistent with the findings for  $D_s$  and  $\eta_0$  in dibutylphthalate solutions.<sup>1</sup>

## SELF-DIFFUSION DATA FOR PS

Later, with the advent of self-diffusion data in the polymer literature, such as that of Bachus and Kimmich,<sup>13</sup> we followed this type of work as a possible source of  $T_{ll}$  and other liquid state events. The cited work<sup>13</sup> represented their  $\log D_0 - T^{-1}$  data on PS as a double Arrhenius plot, with an intersection at  $T_{l\phi}$ . Later we examined in considerable detail the published data of Sillescu *et al.*<sup>14</sup> on selfdiffusion of end tagged PS for a range of molecular weights present in high dilution in a host PS of constant molecular weight. By replotting their data we located two mol. wt. independent transitions, one at 162°C which we assigned as  $T_{ll}$  of the host polymer; the second at 180°C which we ascribed to  $T_{l\phi}$  of the host or the tagged polymer. This analysis is found in Figure 8 of ref 8.

We had concluded from our various studies of  $\log \eta_0 - T^{-1}$  and  $\log D_s - T^{-1}$  data that, depending on the temperature range covered, any given data set was correctly represented by a double or triple Arrhenius plot with the intersection(s) indicating liquid state transitions:  $T_{ll}$  alone;  $T_{ll}$  and  $T_{l\phi}$ ; or  $T_{l\phi}$  alone. While any data set could be fitted by a WLF equation with values of  $C_1$  and  $C_2$  in the usual range, WLF was only a convenient approximation to a real situation.<sup>10</sup> Limited experience with self-diffusion data suggested a similar conclusion. Hence we automatically subjected the data under study<sup>1</sup> to similar scrutiny.

We recognize that multiple Arrhenius lines require a slope and intercept for each line and a set of coordinates defining each intersection. The WLF representation is mathematically much simpler. Our concern, however, is with the physical nature of the liquid state. We believe that the intersections under discussion are caused by liquid state transitions and indicate (or locate) transition temperatures, as already stated.

#### OTHER LITERATURE DATA

We have considered two other literature studies in relation to the present paper: the older one covers  $\eta_0$  measurements on poly(*n*butylmethacrylate), PnBMA, in diethylphthalate over the entire composition range and for a tempeature span from 273–393 K. These

Composition	From $\log \eta_0 - T^{-1a}$	From $\log D_{\rm s} - T^{-1}$ a	TBA <sup>b, a</sup>
40 wt% PS	323	320	319
50 wt% PS	330, 328	338	332

**Table I.**  $T_{ll}$ , K for at-PS-diluent systems

<sup>a</sup> PS MW 43906. Diluent dibutyl phthalate.<sup>1</sup>

<sup>b</sup> PS,  $\bar{M}_n = 37000$  plus diluent.<sup>19,20</sup>

<sup>e</sup> Torsional braid apparatus.<sup>19</sup>

**NOTES**: Brief remarks are in order concerning differences in  $\overline{M}_n$ , diluent and effective frequency of the several methods as they affect the determined values of  $T_n$ .

F4 has a slightly higher  $\bar{M}_n$  which should increase  $T_{ll}$  at a given  $W_2$  value. The  $\bar{M}_n$  effect is further illustrated by the  $T_{ll}$  value of F40—50, namely 362 K by  $D_s$ , with  $\bar{M}_n$ = 335000, 50% concentration.

TBA has a frequency of about 0.3 Hz as compared to estimated effective frequencies of  $10^{-2}$ — $10^{-4}$  for  $D_s$  and  $\eta_0$  so that TBA should give higher  $T_{11}$ 's at fixed  $W_2$ .

Lacking detailed studies, it is difficult to predict the relative effectiveness of the two diluents in lowering  $T_{u}$  since they differ both in molecular weight and presumably in solubility parameters.

data are available in graphical form only.<sup>16</sup> The second involves at-PS of  $\bar{M}_n = 37000$  in a diphenylether type diluent measured by torsional braid analysis to locate  $T_{ll}$  and  $T_{l\phi}$ . Graphical and tabulated data are available.

## PS DILUENT SYSTEMS BY DYNAMIC MECHANICAL LOSS

Since the work under study involves at-PS in dibutylphthalate, the most pertinent comparable work known to us is a study of Gillham *et*  $al.^{17}$  of at-PS,  $\bar{M}_n = 37000$  in an aromatic-ether diluent (C<sub>6</sub>H<sub>5</sub>OC<sub>6</sub>H<sub>4</sub>O)<sub>2</sub>C<sub>6</sub>H<sub>4</sub>, MW 447, using the torsional braid analysis technique to detect  $T_g$  and two  $T > T_g$  transitions,  $T_{ll}$  and  $T_{l\phi}$ , over the entire composition range. The original study shows graphical data only.<sup>17</sup> A review paper contains both graphical and tabulated results.<sup>18</sup>

Table I lists intersection temperatures estimated by us from ref 1  $D_s$  and  $\eta_0$  data and values of  $T_{u}$  from ref 18. It seems likely that these former intersections represent  $T_{u}$ . The same conclusion holds for the 50% PS data. The higher temperature intersections at 343 and 348 K in Figure 1 may be  $T_{l\phi}$  but clearly lower than that shown for TBA data in ref 18, namely 370 K at 0.19 Hz.

## SUMMARY AND CONCLUSIONS

1. The  $D_s$  and  $\eta_0$  data of Nemoto *et al.*<sup>1</sup> can be and should be represented by multiple Arrhenius plots with the low temperature intersection corresponding to the  $T_{ll}$  transition temperature for PS of  $\overline{M}_n \sim 40000$ .

2. A 1:1 correspondence in  $T_{ll}$  values obtained from  $D_s$ ,  $\eta_0$ , and TBA was shown for at-PS with  $\overline{M}_n \sim 40000$ .

3. All results in (1) and (2) above are consistent with the Frenkel view of  $T_{ll}$  as caused by intermolecular segment–segment associations which "melt out" on heating and are forced to lower temperature by diluents.

4.  $D_0 - T^{-1}$  data in solution (1) and in bulk<sup>13,14</sup> provide an absolute method for obtaining  $T_{ll}$  and  $T_{l\phi}$  at very low frequencies, just as  $\eta_0 - T^{-1}$  data do.

5. Even if the existence of the  $T_{ll}$  and/or  $T_{l\phi}$  transitions is rejected, the fact remains that double or triple Arrhenious plots are superior to WLF, leading to the need for an explanation of the slope change(s).

Acknowledgement. The assistance of Mr. Stanton Dent of MMI on computer analysis of the data is appreciated.

#### **REFERENCES AND NOTES**

1. N. Nemoto, T. Kojima, T. Inoue, and M. Kurata,

Polym. J., 20, 875 (1988).

- J. D. Ferry, "Viscoelastic Properties of Polymers," 3rd ed, John Wiley and Sons, New York, N.Y., 1980.
- 3. R. F. Boyer, J. Macromol. Sci., Phy., B18, 461 (1980).
- J. K. Gillham and R. F. Boyer, J. Macromol. Sci., Phys., B13, 497 (1977).
- R. F. Boyer, in "Polymer Yearbook," Vol. 2, R. A. Pethrick, Ed., Harwood Academic Publishers, New York, N.Y., 1985, pp 233-234.
- 6. R. F. Boyer " $T_u$  and Related Liquid State Transitions and Relaxations," Vol. 17, in Ency. Polymer Science Engineering, Jacqueline Kroschwitz, Ex. Ed., John Wiley and Sons, Clopedia of 1989, in press.
- R. F. Boyer, in "Structure in the Liquid State of Amorphous Polymers," S. E. Keinath, R. L. Miller and J. K. Rieke, Ed., Plenum Publishing Corp., New York, N.Y., 1987, pp 135–185.
- 8. R. F. Boyer, J. Polym. Sci., C, 14, 267 (1966).
- 9. R. F. Boyer, Eur. Polym. J., 17, 661 (1981).
- R. F. Boyer, J. Polym. Sci., Polym. Phys. Ed., 23, 21 (1985). See especially pp 37—38.
- 11. R. F. Boyer, unpublished study of  $T_{ii}$  in poly(*n*-butylmethacrylate).
- 12. K. Solc, S. E. Keinath, and R. F. Boyer, *Macro-molecules*, **16**, 1645 (1983). This paper describes the computer procedure but does not apply it to  $\log \eta_o T^{-1}$  data.
- 13. R. Bachus and R. Kimmich, Polymer, 24, 964 (1983).
- 14. M. Antonietti, J. Coutandin, and H. Sillescu, Makromol. Chem., Rapid Commun., 5, 525 (1984). A companion paper by Sillescu (p 591) discusses the relation of inter- and self-diffusion in polymers. Our analysis of the diffusion data is given on pages 150— 2 of ref 7.
- 15. A. Teramoto, R. Okada, and H. Fujita, J. Phys. Chem., 67, 1228 (1965).
- Several reasonably accurate techniques can be used for transcribing graphical data to numerical values, as covered in ref 11.
- J. K. Gillham, J. A. Benci, and R. F. Boyer, *Polym. Eng. Sci.*, **16**, 357 (1976). Plotted data only.
- 18. See Figure 9 and Table 2 of ref 4 for both plotted and tabulated data, respectively.