

Spectral Distribution of the Light Scattered from Rodlike Macromolecules in Solution.

II. The Effect of Optical Anisotropy

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ABSTRACT: The theory of the spectrum of the light scattered from rodlike macromolecules presented earlier (H. Maeda and N. Saitô, *J. Phys. Soc. Japan*, **27**, 984 (1969).) is extended to optically anisotropic rodlike molecules. The spectrum is shown to be the sum of an infinite number of Lorentzians. The forward scattering is shown to be free from the effect of translational diffusion. The numerical calculations for TMV solution are presented.

KEY WORDS Spectral Intensity / Light Scattering / Optical Anisotropy / Rodlike Macromolecule / TMV /

The Brownian motions of macromolecules in solution can be investigated through the spectrum of the scattered light. Thus many theoretical¹⁻⁵ as well as experimental⁶⁻¹⁰ studies on polymer solutions have been carried out. The previous paper⁴ presented a theory for optically isotropic rodlike macromolecules and showed that the coupling between rotatory and translational Brownian motions in rodlike macromolecules is important especially for large scattering angles when the length of the macromolecules is large compared with the wave length of the incident light. This fact was confirmed experimentally by Fujime.¹⁰ In this paper a further calculation for rodlike macromolecules is provided taking into account optical anisotropy.

tered and incident light respectively, μ is the induced electric dipole moment of the scatterer and R_0 indicates the position of the observer.

Now if one takes a coordinate system (x, y, z) fixed in space having the origin 0 in the solution. The z -axis is taken parallel to the electric vector E_0 of the incident light, and the angles between x -axis and the direction K_0 and between z -axis and K_f are denoted as Ψ_0 and Θ_0 respectively. The scattered light is observed in the xz -plane in the direction K_f with an analyser which makes an angle Φ_0 between its polarizing direction and the xz -plane, as shown in Figure 1. Then the component X of the electric vector E_s parallel to the direction of the analyser is given by

COORDINATE SYSTEMS

The amplitude of the electromagnetic wave scattered from a scatterer at the origin is proportional to

$$E_s = -\frac{\exp[-iKR_0]}{4\pi R_0 \epsilon} [K_f \times [K_f \times \mu]] \quad (2.1)$$

$$K = K_f - K_0$$

where K_f , K_0 are the wave vectors of the scat-

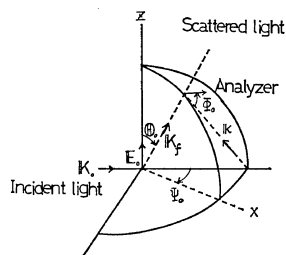


Figure 1. Coordinate systems.

$$X = -[k_f^2 \exp(-i\mathbf{K}\mathbf{R}_0)/4\pi\mathbf{R}_0\epsilon] \otimes [\cos\Phi_0[\mu_x \cos\Theta_0 - \mu_z \sin\Theta_0] + \mu_y \sin\Phi_0] \quad (2.2)$$

where μ_x, μ_y, μ_z are respectively the components of $\boldsymbol{\mu}$ in xyz -coordinate system. The polarizability tensor $\boldsymbol{\alpha}$ of the scatterer in the coordinate system (ξ, η, ζ) fixed on the scatterer is written as

$$\boldsymbol{\alpha} = \begin{pmatrix} \alpha_{\xi\xi} & 0 & 0 \\ 0 & \alpha_{\eta\eta} & 0 \\ 0 & 0 & \alpha_{\zeta\zeta} \end{pmatrix} = (\alpha_0 - \alpha_2/3)\mathbf{1} + \alpha_2 \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \quad (2.3)$$

where it is assumed

$$\alpha_{\xi\xi} = \alpha_{\eta\eta}$$

and put

$$\begin{aligned} \alpha_0 &= \frac{1}{3}(\alpha_{\xi\xi} + \alpha_{\eta\eta} + \alpha_{\zeta\zeta}) \\ \alpha_2 &= (\alpha_{\zeta\zeta} - \alpha_{\xi\xi}) \end{aligned} \quad (2.4)$$

Let T be the transformation matrix between the xyz -system and the coordinate system (ξ, η, ζ) fixed on the scatterer, then one has

$$\begin{pmatrix} \mu_y \\ \mu_x \\ \mu_z \end{pmatrix} = T\boldsymbol{\alpha}T^{-1} \begin{pmatrix} 0 \\ 0 \\ E_0 \end{pmatrix} \quad (2.5)$$

Here in order to use the distribution function for the position and orientation of the molecule obtained in the previous paper, one introduces another coordinate system (x', y', z') with the z' -axis parallel to \mathbf{K} and x' -axis in the plane containing the direction of \mathbf{K} and the z -axis with the origin 0. The orientation of the molecule is described by the polar angle θ, φ (or simply Ω) in the coordinate system (x', y', z') . Hence the transformation matrix is rewritten as,

$$T = T_1 T_2$$

$$T_1 = \begin{pmatrix} \cos\gamma \cos\phi & -\cos\gamma \sin\phi & -\sin\gamma \\ \sin\phi & \cos\phi & 0 \\ \sin\gamma \cos\phi & -\sin\gamma \sin\phi & \cos\gamma \end{pmatrix}$$

$$T_2 = \begin{pmatrix} \cos\theta \cos\varphi & \sin\varphi \cos\theta & -\sin\theta \\ -\sin\varphi & \cos\varphi & 0 \\ \sin\theta \cos\varphi & \sin\theta \sin\varphi & \cos\theta \end{pmatrix} \quad (2.6)$$

where T_1 and T_2 are transformation matrices between (x, y, z) and (x', y', z') and between (x', y', z') and (ξ, η, ζ) systems respectively, and γ and ϕ are the angles between the direction of \mathbf{K} and the z -axis, and between the x -axis and the $z'x'$ -plane respectively. Substituting eq 2.5 and 2.6 into eq 2.2, one has

$$X = K\tilde{X} \quad (2.7)$$

$$K = -k_f^2 \exp(-i\mathbf{K}\mathbf{R}_0)/4\pi\mathbf{R}_0\epsilon \quad (2.8)$$

$$\begin{aligned} \tilde{X} &= -\alpha_2 E_0 \{A \sin\theta \cos\theta \cos\varphi + B \cos^2\theta \\ &+ C \sin^2\theta \cos^2\varphi + H \sin^2\theta \sin\varphi \cos\varphi \\ &+ M \sin\theta \cos\theta \sin\varphi\} \\ &+ (\alpha_0 - \frac{1}{3}\alpha_2) D E_0 \end{aligned} \quad (2.9)$$

where the coefficients $A, B, C, H, M,$ and D are given in Appendix A. They are functions of $\Theta_0, \Psi_0,$ and Φ_0 and the polarizabilities α_0 and α_2 .

The rodlike molecule considered here is represented by an array of N identical beads (or scatterers) connected in a straight line. The beads are anisotropic scatterers with the anisotropic axis in the direction of the molecule. The position of the i th bead and the center of gravity at time t are designated as $\mathbf{R}_i(t)$ and $\mathbf{r}(t)$ respectively in the $x'y'z'$ -system. Put

$$\mathbf{r}_i(t) = \mathbf{R}_i(t) - \mathbf{r}(t) \quad (2.10)$$

and let the orientation of the molecule be Ω in polar coordinates as mentioned above. The distribution function $f(\mathbf{r}, \Omega, t)$ has already been obtained,⁴ especially the function $f(\mathbf{r}, \Omega, t|\Omega_0)$ of the molecule at time t under the initial condition that its center of gravity lies at the origin with its orientation $\Omega_0(\theta_0, \varphi_0)$. In the following, use will be made of the functions \bar{g} and g defined as

$$\begin{aligned} \bar{g} &= e^{k^2 D_t t} g(\mathbf{K}, \Omega, t|\Omega_0) \\ &= e^{k^2 D_t t} \int d\mathbf{r} e^{i\mathbf{r}\cdot\mathbf{K}} f(\mathbf{r}, \Omega, t|\Omega_0) \end{aligned} \quad (2.11)$$

where D_t is the average translational diffusion constant $kT(B_1 + 2B_2)/3$ as given in ref 4.

SPECTRAL INTENSITY OF SCATTERED LIGHT

The spectral intensity I of the light scattered

from the rod-like macromolecule as described in ref 2 is given by²

$$\begin{aligned}
 & I(\Delta\omega, \Phi_0, \Theta_0, \Psi_0)K^{-2} \\
 &= \int_{-\infty}^{+\infty} \langle \sum_{l,m} \tilde{X}(t, m) \tilde{X}^*(0, l) \\
 & \quad \times \exp [i\mathbf{K}(\mathbf{R}_l(0) - \mathbf{R}_m)] \rangle \exp(-i\Delta\omega t) dt \\
 &= \int_{-\infty}^{+\infty} dt \exp[-(i\Delta\omega - \mathbf{K}^2 D_i)t] \sum_{l,m} \\
 & \quad \times \int d\Omega d\Omega_0 \frac{1}{4\pi V} \tilde{X}(t, l) \tilde{X}^*(0, m) \\
 & \quad \times \exp[-i\mathbf{K}(\mathbf{r}_l(0) - \mathbf{r}_m)] \bar{g}(t, \mathbf{K}, \Omega | \Omega_0) \\
 &= \int_{-\infty}^{+\infty} dt \exp[(-i\Delta\omega - D_i \mathbf{K}^2)t] \sum_{l,m} \\
 & \quad \times \int d\Omega d\Omega_0 \frac{1}{4\pi V} \bar{g}(t, \mathbf{K}, \Omega | \Omega_0) \tilde{X}(t, l) \tilde{X}^*(0, m) \\
 & \quad \times \sum_v i^v \left(v + \frac{1}{2}\right) \left(\frac{2\pi}{kr_l}\right)^{1/2} J_{v+1/2}(kr_l) P_v(\cos \theta_0) \\
 & \quad \times \sum_u (-i)^u \left(u + \frac{1}{2}\right) \left(\frac{2\pi}{kr_m}\right)^{1/2} J_{u+1/2}(kr_m) \\
 & \quad \times P_u(\cos \theta) \quad (3.3)
 \end{aligned}$$

The function g can be written as⁴

$$\bar{g} = \mathcal{L}^{-1} q(\mathbf{P}) = \frac{1}{2\pi i} \int_{-\infty}^{+\infty} e^{Pt} q(\mathbf{P}) d\mathbf{P} \quad (3.4)$$

$$\begin{aligned}
 q(\mathbf{P}) &= \sum \varepsilon^k (a_{nm}^{(k)}(\mathbf{P}) \cos m\varphi \\
 & \quad + b_{nm}^{(k)}(\mathbf{P}) \sin m\varphi) P_n^m(\cos \theta) \\
 \varepsilon &= 3(B_3 - B_1)/(2B_1 + B_3) \quad (3.5)
 \end{aligned}$$

where B_1 and B_3 are the mobilities along the ξ and ζ directions of the molecule, and the coefficients $a_{nm}^{(k)}(p)$ and $b_{nm}^{(k)}(p)$ are given in Appendix B. Substitution eq 3.4 and 3.5 into eq 3.3 gives finally:

$$\begin{aligned}
 & I(\Delta\omega, \Phi_0, \Theta_0, \Psi_0)K^{-2} \\
 &= \pi^2 E_0^2 \sum_k (\varepsilon \mathbf{K}^2 D_i)^k \sum_{n, \text{even}} L_f \mathcal{L}^{-1} [\alpha_2^2 (A^2 \\
 & \quad + M^2) X_n [X_n]_k + \frac{1}{4} \alpha_2^2 (C^2 + H^2) Y_n [Y_n]_k \\
 & \quad + \{\tilde{A} [J_n] - \tilde{B} [K_n]\} \{\tilde{A} J_n - \tilde{B} K_n\}] \quad (3.6)
 \end{aligned}$$

L_f is the operator defined by

$$L_f u(t) = \int_{-\infty}^{+\infty} u(|t|) \exp[-i\Delta\omega t + k^2 D_i |t|] dt \quad (3.7)$$

The functions X_n , Y_n , $[X_n]_k$, $[Y_n]_k$, J_n , K_n and $[J_n]$, $[K_n]$ are listed in Appendix C. Further the following definitions are introduced

$$L_f \mathcal{L}^{-1} [W_n]_l = [[W_n]_l] \quad (3.8)$$

where $W_n = X_n$, Y_n , J_n , and K_n . Using these quantities one can rewrite eq 3.6 as

$$\begin{aligned}
 & I(\Delta\omega, \Phi_0, \Theta_0, \Psi_0)K^{-2} \\
 &= \pi E_0^2 \sum_l (\varepsilon \mathbf{K}^2 D_i)^l \sum_{n, \text{even}} [\alpha_2^2 (A^2 + M^2) X_n [[X_n]_l] \\
 & \quad + \frac{1}{4} \alpha_2^2 (C^2 + H^2) Y_n [[Y_n]_l] \\
 & \quad + \{\tilde{A} [[J_n]_l] - \tilde{B} [[K_n]_l]\} \{\tilde{A} J_n - \tilde{B} K_n\}] \quad (3.9)
 \end{aligned}$$

Explicit expressions of $[[W_n]_l]$ are given in Appendix C. Equation 3.9 gives the spectral intensity of the scattered light in any arbitrary direction $(\Theta_0, \Psi_0, \Phi_0)$.

RESULTS AND DISCUSSION

(i) The spectral intensity of the scattered light consists of terms proportional to α_0^2 and α_2^2 as well as $\alpha_0 \alpha_2$. The last term does not exist in the anisotropic flexible chain, and further in rodlike molecules of the length less than the wave length of the incident light as treated by Pecora. This fact was already pointed out by Tagami.¹¹

(ii) Pecora² showed that the isotropic part is a single Lorentzian and the anisotropic part consists of five Lorentzian terms in the case of rods whose length is short compared with the wave length of the light. In larger rodlike molecules, however, both the isotropic and the anisotropic parts consist of infinite numbers of Lorentzians.

(iii) The coupling between translational and rotational Brownian motions gives rise to an effect similar to the spectral intensity for both the optically anisotropic and isotropic rodlike molecules.

For numerical calculations one uses the values of tobacco mosaic virus, *i.e.*, 3000 Å in length, 150 Å in diameter, 390 sec⁻¹ in the rotational diffusion constant and 0.3×10^{-7} cm²/sec in the translational diffusion constant. The results are given in Figures 2—6, where the ordinate scales are arbitrary. In Figure 2, one takes $\Theta_0 = 160.4^\circ$, $\Phi_0 = 0$ and $\alpha_2/\alpha_0 = 0$ and 0.4, and $D_i = 1 \times 10^{-7}$ cm²/sec. The full and dotted lines show the spectra with and without the coupling between the translational and the rotational Brownian motions respectively. It is clear that, in the case of Figure 2, the effect of the coupling is by no

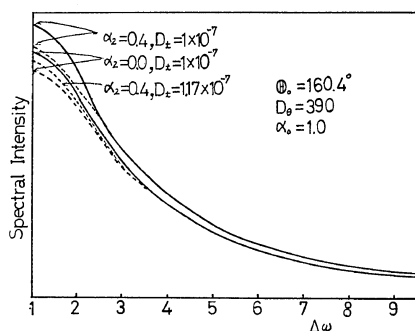


Figure 2. Spectra of the scattered lights from anisotropic rodlike macromolecules compared with those from isotropic ones. Full lines: Results taking into account the coupling of the rotatory and translational Brownian motions, Broken lines: Results without taking the coupling into account.

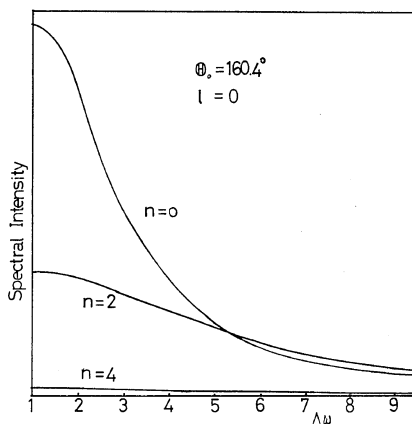


Figure 4. Contribution of the relaxation modes $l=0$ and $n=0, 2, 4$ at $\theta_0=160.4^\circ$ to the total spectral intensity.

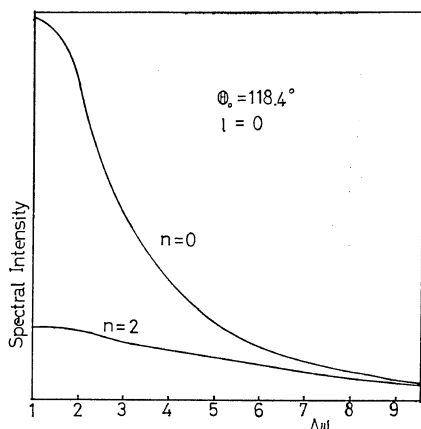


Figure 3. Contribution of the relaxation modes $l=0$ and $n=0, 2$ at $\theta_0=118.4^\circ$ to the total spectral intensity.

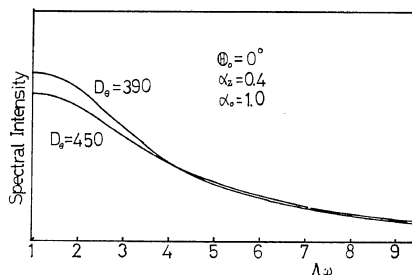


Figure 5. Spectral intensity of the molecules with different rotational diffusion constants in the forward scattering.

means negligible and also the optical anisotropy gives rise to an increase in the spectral intensity. One sees further that the higher the value of the diffusion constant D_i , the smaller the intensity.

In Figure 3, one takes scattering angle $\theta_0=118.4^\circ$ and $l=0$. The term $l=0$ corresponds to the case free from coupling between the translational and rotatory diffusions. The number refers to the relaxation mode number, which is related to the relaxation time as $[D_i k^2 + D_0 n(n+1)]^{-1}$. The term for $n \geq 4$ is negligibly small. Figure 4 is the similar one for $\theta_0=160.4^\circ$ and $l=0$. One sees that the larger the scattering

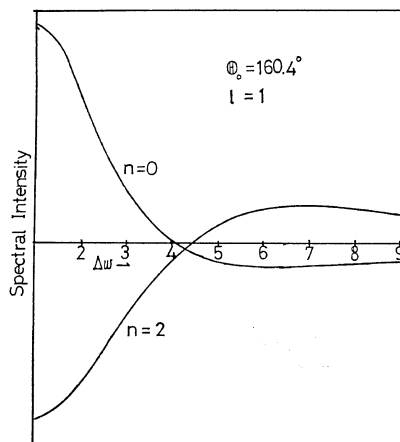


Figure 6. Contribution of the relaxation modes $l=1$ and $n=0, 2$ to the total spectral intensity. The curve for $n=2$ is drawn ten times larger in ordinate than the one for $n=0$.

angle the larger the contribution of the rotational relaxation to the total intensity. In Figure 5 one takes $\theta_0=0^\circ$, $\phi_0=90^\circ$ (forward scattering with crossed polarizer and analyser), $\alpha_2/\alpha_0=0.40$ and $D_\theta=390/\text{sec}$ or $450/\text{sec}$. In this case the rotational Brownian motion only is observed in the spectrum.⁹ In general the spectral intensity is composed of terms of various sets of n and l . In Figure 6, terms for $l=1$ and $n=0, 2$, are given. Figure 2 is the sum of all (n, l) terms.

APPENDIX A

$$\tilde{A}=2B-C$$

$$\begin{aligned} &= \frac{1}{2(1-\sin\theta_0)\cos\psi_0} \{ (3-\sin^2\theta_0)\cos\psi_0\cos\phi_0 \\ &\quad + 3\cos\theta_0\sin\phi_0\sin\psi_0 - 2\sin\theta_0\cos\phi_0 \} \\ \tilde{B} &= -2D(\alpha_0+\alpha_2/3) - C\alpha_2 \\ &= \{ 2(\alpha_0+\alpha_2/3)\sin\theta_0\cos\phi_0 \\ &\quad - \frac{\alpha_2}{2(1-\sin\phi_0\sin\psi_0)^2} \{ 2\sin\theta_0\cos\phi_0 \\ &\quad - (1+\sin^2\theta_0)\cos\phi_0\cos\psi_0 \\ &\quad - \cos\theta_0\sin\psi_0\sin\phi_0 \} \} \end{aligned}$$

$$M^2 + A^2$$

$$\begin{aligned} &= \frac{(\sin\phi_0\cos\psi_0 - \sin\phi_0\sin\theta_0 - \sin\psi_0\cos\phi_0\cos\theta_0)^2}{2(1-\sin\theta_0\cos\psi_0)^2(1+\sin^2\theta_0-2\sin\theta_0\cos\psi_0)} \\ &\quad \times \{ \sin^3\theta_0\cos\psi_0 + \sin^2\theta_0(2\sin^2\psi_0-1) \\ &\quad - \sin\theta_0\cos\psi_0 + 1 \} \end{aligned}$$

$$\begin{aligned} C^2 &= \frac{1}{4(1-\sin\theta_0\cos\psi_0)^2} \{ 2\sin\theta_0\cos\phi_0 \\ &\quad - (1+\sin^2\theta_0)\cos\phi_0\cos\psi_0 \\ &\quad - \cos\theta_0\sin\psi_0\sin\phi_0 \}^2 \end{aligned}$$

$$\begin{aligned} H^2 &= \frac{1}{2(1-\sin\theta_0\cos\psi_0)} \{ \sin\psi_0\cos\phi_0\cos\theta_0 \\ &\quad + \sin^2\theta_0\sin\phi_0 - \sin\phi_0\cos\psi_0 \}^2 \end{aligned}$$

APPENDIX B

The explicit expressions of the coefficients $a_{nm}^l(p)$ and $b_{nm}^l(p)$ are given in ref 4, and are written in the following way,

$$a_{nm}^l \equiv \sum_{j=-l}^l A_{nm}^{l,j} P_{n+2j}(\cos\theta_0) \cos m\varphi_0$$

and b_{nm} 's are given by replacing $\cos m\varphi_0$ by \sin

$m\varphi_0$ with same coefficients. Also the coefficients $A_{nm}^{l,j}$ are defined as follows. For $l=0$ and $l=1$, in particular,

$$A_{m,0}^{0,0} = \varepsilon_n \frac{2n+1}{4\pi}$$

$$A_{n,1}^{0,0} = \varepsilon_n \frac{2n+1}{2\pi} \cdot \frac{(n-1)!}{(n+1)!}$$

$$A_{n,2}^{0,0} = \varepsilon_n \frac{2n+1}{2\pi} \frac{(n-2)!}{(n+2)!}$$

$$A_{n+2,0}^{1,1} = -\varepsilon_n \varepsilon_{n+2} \frac{(n+1)(n+1)}{4\pi(2n+3)}$$

$$A_{n,0}^{1,0} = -\varepsilon_n \frac{2}{4\pi} \left\{ \frac{n^2}{(2n-1)} + \frac{(n+1)^2}{(2n+3)} - \frac{1}{3}(2n+1) \right\}$$

$$A_{n+2,1}^{1,1} = -\varepsilon_n \varepsilon_{n+2} \frac{1}{2\pi(2n+3)}$$

$$\begin{aligned} A_{n,1}^{1,0} &= -\varepsilon_n \frac{(n-1)!}{2\pi(n+1)!} \left\{ \frac{(n^2-1)}{(2n-1)} + \frac{n(n+2)}{2n+3} \right. \\ &\quad \left. - \frac{1}{3}(2n+1) \right\} \end{aligned}$$

$$A_{n+2,2}^{1,1} = -\varepsilon_n \varepsilon_{n+2} \frac{1}{2\pi(n+1)(n+2)(n+3)}$$

$$\begin{aligned} A_{n,2}^{1,0} &= -\varepsilon_n \frac{(n-2)!}{2\pi(n+2)!} \left\{ \frac{n^2-4}{2n-1} + \frac{(n-1)(n+3)}{2n+3} \right. \\ &\quad \left. - \frac{1}{3}(2n+1) \right\} \end{aligned}$$

where

$$\varepsilon_n = \frac{1}{P + D_\theta n(n+1)}$$

APPENDIX C

$$\begin{aligned} J_n &= \sum_{s,j} i^s (s+\frac{1}{2}) (2\pi/kr_j)^{1/2} J_{s+1/2}(kr_j) \\ &\quad \times \int \cos^2\theta P_s(\cos\theta) P_n(\cos\theta) \sin\theta d\theta \\ &= \sum_j i^n (2\pi/kr_j)^{1/2} \left\{ -\frac{(n+1)(n+2)}{(2n+1)(2n+3)} J_{n+5/2}(kr_j) \right. \\ &\quad + \frac{1}{2n+1} \left(\frac{n^2}{2n-1} + \frac{(n+1)^2}{2n+3} \right) \\ &\quad \left. \times J_{n+1/2}(kr_j) - \frac{n(n-1)}{(2n-1)(2n+1)} J_{n-3/2}(kr_j) \right\} \end{aligned}$$

$$\begin{aligned} K_n &= \sum_{s,j} i^s (s+\frac{1}{2}) (2\pi/kr_j)^{1/2} J_{s+1/2}(kr_j) \\ &\quad \times \int P_s(\cos\theta) P_n(\cos\theta) \sin\theta d\theta \end{aligned}$$

$$\begin{aligned}
 &= \sum_j i^n (2\pi/kr_j)^{1/2} J_{s+1/2}(kr_j) \\
 X_n &= \sum_{s,j} i^s (s+\frac{1}{2})(2\pi/kr_j)^{1/2} J_{s+1/2}(kr_j) \\
 &\quad \times \int \sin \theta \cos \theta P_s(\cos \theta) P_n^1(\cos \theta) \sin \theta d\theta \\
 &= \sum_j i^n (2\pi/kr_j)^{1/2} \frac{1}{2n+1} \cdot \frac{(n+1)!}{(n-1)!} \\
 &\quad \times \left\{ \frac{2n+3}{n+2} J_{n+5/2}(kr_j) \right. \\
 &\quad \left. + \frac{2n+1}{(2n-1)(2n+3)} J_{n+1/2}(kr_j) \right. \\
 &\quad \left. - \frac{n-1}{2n-1} J_{n-3/2}(kr_j) \right\} \\
 Y_n &= \sum_{s,j} i^s (s+\frac{1}{2})(2\pi/kr_j)^{1/2} J_{s+1/2}(kr_j) \\
 &\quad \times \int \sin^2 \theta P_s(\cos \theta) P_n^2(\cos \theta) \sin \theta d\theta \\
 &= \sum_j i^n (2\pi/kr_j)^{1/2} \frac{1}{2n+1} \frac{(n+2)!}{(n-2)!} \\
 &\quad \times \left\{ -\frac{1}{2n+3} J_{n+5/2}(kr_j) \right. \\
 &\quad \left. - \frac{2(2n+1)}{(2n-1)(2n+3)} J_{n+1/2}(kr_j) \right. \\
 &\quad \left. - \frac{1}{2n-1} J_{n-3/2}(kr_j) \right\} \\
 [W_n]_l &= \sum_{j=0}^l (2-\delta_{j,0}) A_{n+2j,r}^{l,j} W_{n+2j}
 \end{aligned}$$

where $\delta_{j,0}$ is Kronecker's delta function and the suffix r in $A_{n+2j,r}^{l,j}$ is defined as,

$$r = \begin{cases} 0 & \text{for } W_n = J_n = K_n \\ 1 & \text{for } W_n = X_n \\ 2 & \text{for } W_n = Y_n \end{cases}$$

$$[[W_n]]_l \equiv \sum_{j=0}^l u_j^l(w) W_{n+2j}$$

In particular for $l=0$,

$$\begin{aligned}
 [[J_n]]_0 &= \frac{2n+1}{4\pi} \cdot \frac{[D_\theta n(n+1) + D_t k^2]}{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2} J_n \\
 [[K_n]]_0 &= \frac{2n+1}{4\pi} \cdot \frac{2[D_\theta n(n+1) + D_t k^2]}{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2} K_n \\
 [[X_n]]_0 &= \frac{2n+1}{2\pi} \cdot \frac{(n-1)!}{(n+1)!} \\
 &\quad \times \frac{2[D_\theta n(n+1) + D_t k^2]}{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2} X_n
 \end{aligned}$$

$$\begin{aligned}
 [[Y_n]]_0 &= \frac{2n+1}{2\pi} \cdot \frac{(n-2)!}{(n+2)!} \\
 &\quad \times \frac{2[D_\theta n(n+1) + D_t k^2]}{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2} Y_n
 \end{aligned}$$

For $l=1$, the coefficients $U_j^1(w)$ are given as follows,

$$\begin{aligned}
 U_1^1(J) &= -\frac{2(n+1)(n+2)}{4\pi(2n+3)} \omega_3 \\
 U_0^1(J) &= -\frac{1}{4\pi} \left\{ \frac{n^2}{2n-1} + \frac{(n+1)^2}{2n+3} \right. \\
 &\quad \left. - \frac{1}{3}(2n+1) \right\} \omega_2 \\
 U_1^1(X) &= -\frac{2}{2\pi(2n+3)} \omega_2 \\
 U_0^1(X) &= -\frac{1}{2\pi} \frac{(n-1)!}{(n+1)!} \left\{ \frac{(n-1)(n+1)}{2n-1} \right. \\
 &\quad \left. + \frac{n(n+2)}{2n+3} - \frac{1}{3}(2n+1) \right\} \omega_2 \\
 U_1^1(Y) &= -\frac{2}{2\pi(n+1)(n+2)(n+3)} \omega_3 \\
 U_0^1(Y) &= -\frac{1}{2\pi} \cdot \frac{(n-2)!}{(n+2)!} \left\{ \frac{(n-2)(n+2)}{(2n-1)} \right. \\
 &\quad \left. + \frac{(n-1)(n+3)}{(2n+3)} - \frac{1}{3}(2n+1) \right\} \omega_3
 \end{aligned}$$

where

$$\begin{aligned}
 \omega_1 &= \frac{2[D_\theta n(n+1) + D_t k^2]}{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2} \\
 \omega_2 &= \frac{2\{[D_\theta n(n+1) + D_t k^2]^2 - \Delta\omega^2\}}{\{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2\}^2} \\
 \omega_3 &= \frac{1}{(2n+3)D_\theta} \left\{ \frac{D_\theta n(n+1) + D_t k^2}{[D_\theta n(n+1) + D_t k^2]^2 + \Delta\omega^2} \right. \\
 &\quad \left. - \frac{D_\theta(n+2)(n+3) + D_t k^2}{[D_\theta(n+2)(n+3) + D_t k^2]^2 + \Delta\omega^2} \right\}
 \end{aligned}$$

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