Controlled dephasing of electrons by non-gaussian shot noise

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In a ‘controlled dephasing’ experiment^{1–3}, an interferometer loses its coherence owing to entanglement of the interfering electron with a controlled quantum system, which effectively is equivalent to path detection. In previous experiments, only partial dephasing was achieved owing to weak interactions between many detector electrons and the interfering electron, leading to a gaussian-phase randomizing process^{4–11}. Here, we report the opposite extreme, where interference is completely destroyed by a few (that is, one to three) detector electrons, each of which has a strong randomizing effect on the phase. We observe quenching of the interference pattern in a periodic, lobe-type fashion as the detector current is varied, and with a peculiar V-shaped dependence on the detector’s partitioning. We ascribe these features to the non-gaussian nature of the noise, which is also important for qubit decoherence¹². In other words, the interference seems to be highly sensitive to the full counting statistics of the detector’s shot noise^{13–16}.

Our system is based on the previously developed electronic two-path Mach–Zehnder interferometer (MZI)^{17–20}. Here, we use two edge channels in the integer quantum Hall effect regime, at filling factor $\nu = 2$ (see Fig. 1). The inner edge channel was partitioned and served as a which-path detector. Other than that, the device and the measurement technique were similar to those described in refs 17–19. The MZI was fabricated within a high-mobility two-dimensional electron gas. The two paths were formed by splitting the outer edge channel with a quantum-point-contact constriction QPC1. After enclosing a magnetic flux, the two paths joined in QPC2 and interfered. Metallic ohmic contacts served as sources S1, S2 and S3 and drains D1 and D2. Changing the enclosed flux by $\Delta\Phi$ (via the modulation gate, MG) changed the Aharonov–Bohm phase $\varphi = 2\pi\Delta\Phi/\Phi_0$ ($\Phi_0 = h/e$ the flux quantum)²¹, leading to phase-dependent transmission coefficients; for example, from S2 to D2:

$$T_{S2 \rightarrow D2} \equiv T_{MZI} = |t_{QPC1} t_{QPC2} + e^{i\varphi} r_{QPC1} r_{QPC2}|^2 = T_0 + T_\varphi \cos \varphi, \quad (1)$$

with t and r being the corresponding transmission and reflection amplitudes. The measured visibility, defined as $\nu = T_\varphi/T_0$, ranged from 30% to 60% (refs 17–19). We attribute the non-ideal visibility to phase fluctuations owing to external noise¹⁰.

The inner edge channel served as a path detector (see Fig. 1, caption). When QPC0 was tuned to partition the detector channel (which was biased, $V_{\text{det}} = V_{S3}$), electrons in the upper path of the interferometer became entangled with those in

the detector, resulting in a lower visibility. This dephasing process can be considered as ‘path detection’⁹ or, alternatively, as phase scrambling owing to potential fluctuations in the partitioned detector channel²². The interaction between the inner and the outer channels was characterized before the actual dephasing experiment by first fully transmitting and then fully reflecting the biased inner edge channel emanating from S3 (with QPC0). Full transmission ($T_{QPC0} = 1$) did not lead to an observable effect on the Aharonov–Bohm oscillations of the MZI as a function of V_{S3} (Fig. 2a). However, full reflection ($R_{QPC0} = 1$) had a strong effect on the phase of the interference pattern, which varied linearly with V_{S3} (reaching $\sim 2\pi$ for $V_{\text{det}} \sim 19 \mu\text{V}$), but with nearly no effect on the visibility (Fig. 2b). From this, we can estimate that a mere $n = 1$ –3 electrons suffice to quench the interference for $V_{\text{det}} \sim 19 \mu\text{V}$ and an interferometer dwell time $\tau = L/v_g$ (with $L \approx 10 \mu\text{m}$ and $v_g = (3$ – $10) \times 10^6 \text{ cm s}^{-1}$). This strong coupling between the edges sets the present experiment apart from previous ones. Note that total dephasing using an adjacent edge channel was also achieved in a quantum dot²³.

When QPC0 was tuned to partition the inner channel ($0 < T_{QPC0} < 1$), the visibility diminished as V_{det} increased. Figure 3 shows the dependence of the visibility on T_{QPC0} (partitioning) for three different detector voltages. As the bias V_{S3} increased, the visibility turned from a smooth parabolic curve to a sharp, V-shaped dependence, with a minimum at $T_{QPC0} \sim 0.5$. The dispersion among the experimental points at higher bias resulted from resonances in T_{QPC0} (see inset). We argue below that the V-shaped dependence is a signature of the non-gaussian nature of the detector noise.

We first study a simple model where exactly one electron in the detector scrambles the phase of an interfering electron. Detector electrons were injected with a probability $R_{QPC0} = 1 - T_{QPC0}$ into the channel that interacted with the interferometer. Depending on the presence or absence of a detector electron, the extra phase, $\delta\varphi$, acquired by an interfering electron fluctuated between two values: $\delta\varphi = \gamma V_{\text{det}}$ ($\gamma = (2\pi/19) \text{ rad } \mu\text{V}^{-1}$, from Fig. 2b) and $\delta\varphi = 0$, respectively. Averaging the $\cos \varphi$ term in equation (1) over the two possibilities leads to a visibility^{2,14}:

$$\nu = |\langle e^{i\delta\varphi} \rangle| = |T_{QPC0} + R_{QPC0} e^{i\gamma V_{\text{det}}}|. \quad (2)$$

Equation (2) does not have fitting parameters, so it can be compared directly with the experimental results. Moreover, for

Figure 3 The effect of partitioning the detector channel (by QPCO) on the visibility of the interfering signal, at three different detector bias values. Dotted lines show the best fitted gaussian model; solid lines show the ‘microscopic’ non-gaussian predictions. As V_{det} increases, the dependence of the visibility on T_{QPCO} turns from a smooth one to a sharp V-shape (at $V_{\text{det}} = 15 \mu\text{V}$). **a**, Data points and theoretical predictions for $V_{\text{det}} = 4 \mu\text{V}$ and $9.5 \mu\text{V}$. The gaussian model (dotted lines) is adequate at low bias ($4 \mu\text{V}$) and gives a reasonable fit at intermediate bias ($9.5 \mu\text{V}$). The dashed line is the V-shaped prediction of a single-detector-electron model (equation (2)) for $V_{\text{det}} = 9.5 \mu\text{V}$. Inset: The conductance of QPCO as a function of gate voltage shows sharp resonances. This explains the lack of visibility measurements in the range $0.1 < T_{\text{QPCO}} < 0.4$ and its dispersion at large detector bias (owing to the dependence of the resonances on bias). **b**, Data points and theoretical predictions for $V_{\text{det}} = 15 \mu\text{V}$ (data from ref. 19). The V shape is observed experimentally, as predicted by the ‘microscopic’ non-gaussian model. The gaussian model fails for the arbitrary fitting parameter; the best fit (upper curve) strongly overestimates the visibility at $T_{\text{QPCO}} = 0.5$ and underestimates the visibility for $0.7 < T_{\text{QPCO}} < 0.9$.

Figure 4 The evolution of the interference pattern as a function of detector voltage, V_{det} , for a partitioned detector channel. **a**, Visibility evolution. **b**, Phase evolution (modulo 2π). The non-monotonic behaviour of the visibility is a clear sign of dephasing by non-gaussian noise (data: solid black line and filled circles). Dashed lines: Prediction of the improved theoretical model, which fully accounts for the effects of binomial shot noise, see text, equation (4). The discrepancy at negative V_{det} results from slight nonlinearity of QPCO, leading to a non-accurate R_{QPCO} ($T_{\text{QPCO}} \sim 0.5$, $V_{\text{QPCO}} = -0.0272 \text{ V}$; see inset of Fig. 3a).

To overcome the quantitative shortcomings of equation (2), we now outline a more microscopic approach that predicts the main observed features. A full description of this approach will be provided in a subsequent publication²⁵. We assume that every interfering electron accumulates a random phase, $\delta\varphi$, as it traverses the upper arm of the MZI, owing to the coupling with the

fluctuating electron density in the detector channel. Treating the detector density, ρ_{det} , classically, the phase should be:

$$\delta\varphi = \int_0^\tau \int u(v_g t - x) \rho_{\text{det}}(x, t) dx dt = \int w(x) \rho_{\text{det}}(x) dx, \quad (3)$$

where $u(x)$ is the inter-channel interaction potential, v_g is the electron velocity in the MZI and τ is the traversal time in the upper path. The electron density, propagating with velocity, v_{det} , obeys: $\rho_{\text{det}}(x, t) = \rho_{\text{det}}(x - v_{\text{det}} t, 0) \equiv \rho_{\text{det}}(x - v_{\text{det}} t)$, which yields $w(x) = \int_0^\tau u[(v_g - v_{\text{det}}) t' - x] dt'$. It can be shown that equation (3) can be used even in the quantum case ($\delta\varphi \mapsto \delta\hat{\varphi}$ and $\rho_{\text{det}} \mapsto \hat{\rho}_{\text{det}}$) to calculate the visibility $\nu = |\langle e^{i\delta\hat{\varphi}} \rangle|$ as long as the interfering electron is treated in a single-particle picture. This approach neglects Pauli blocking⁸, which has to be taken into account in a phenomenological way when evaluating the visibility. We find (see the Methods section) that the visibility is a product of factors, each in the form of the single-particle expression of equation (2):

$$\nu = |\langle e^{i\delta\hat{\varphi}} \rangle| = \prod_j |T_{\text{QPCO}} + R_{\text{QPCO}} e^{i\delta\varphi_j}|. \quad (4)$$

The phases $\delta\varphi_j$ are the eigenvalues of the matrix $w_{k'k}$ (the Fourier transform of $w(x)$), which has been restricted to transitions between plane-wave states k', k within the voltage window. They depend on the detector voltage, V_{det} , and obey a ‘sum rule’: $\sum_j \delta\varphi_j = \gamma V_{\text{det}}$. In the limit $V_{\text{det}} \rightarrow 0$, it can be shown that only one non-zero eigenvalue remains and the result reduces to equation (2) (ref. 25).

Choosing the Fourier transform of $w(x)$ as a lorentzian, with its full-width at half-maximum as the single fitting parameter ($\Delta = 12.4 \mu\text{eV}$; the height being deduced from the observed value of ν), we plotted in Figs 3 and 4 the calculated visibility from equation (4). The plot reproduced the phase slip and zero visibility at $V_{\text{det}} = 15 \mu\text{V}$, the second lobe and the eventual decay at higher detector voltages. As more than one detector electron participated

in the dephasing process, the largest eigenvalue, $\delta\phi_1$, becomes smaller than γV_{det} (because of the ‘sum rule’), such that the zero visibility (when $\delta\phi_1 = \pi$) is reached at a higher V_{det} than predicted by equation (2). Although the quality of the fit is good, it may be improved further if $w(x)$, determined by the microscopic physics of the edge channels, were known more precisely. We may conclude that for $V_{\text{det}} < 6 \mu\text{V}$ a single detecting electron dephases the MZI, whereas at $V_{\text{det}} \sim 30 \mu\text{V}$ the number is at most three.

In summary, we presented a unique behaviour of an electron interferometer coupled to a which-path detector. Very strong interactions between electrons in both systems led to dephasing by the characteristic binomial, non-gaussian, shot noise in the detector. The dephased visibility had a linear, V-shaped dependence on the partitioning of the detector’s current, and non-monotonic, periodic, lobe pattern decay as a function of the detector current. This entanglement between nearly single pairs of electrons may be exploited (in future experiments) to test Bell’s inequalities in a system where the detector channel is replaced by another two-path interferometer^{26–28}.

METHODS

SAMPLE AND MEASUREMENTS

The edges of the sample are defined by plasma etching of a GaAs–AlGaAs heterostructure, embedding a high-mobility two-dimensional electron gas, 80 nm below the surface. Two edge channels are formed by applying a perpendicular magnetic field of ~ 3 T, leading to a filling factor of 2 in the bulk (electron temperature ~ 15 mK). Transmission of the outer channel to D2 is measured by applying $\sim 1 \mu\text{V}$ at ~ 1 MHz at S2. The signal at D2 (see Fig. 1), filtered by a cold resonant circuit tuned to 1 MHz with a bandwidth of 30 kHz, is amplified by a low-noise preamplifier at 4.2 K. Note that the inner, small, ohmic contact ($3 \times 3 \mu\text{m}^2$) serves as both D1 and S1. The inner contact was grounded using an air bridge, which had no influence on the transport of the two edge channels underneath.

EVALUATION OF VISIBILITY

Being a many-particle quantum device, the detector’s density cannot be expressed either as a classical function or in a single-particle language. Hence, equation (3) should be rewritten in terms of detector electron operators, $\hat{d}_k$, and the matrix elements of $w(x)$ with respect to a plane-wave basis:

$$\delta\hat{\phi} = \sum_{k,k'} w_{k'k} \hat{d}_{k'}^+ \hat{d}_k.$$

The occupation of each k state fluctuates independently with $n_{k'k} \equiv \langle \hat{d}_k^+ \hat{d}_k \rangle = n_k \delta_{k'k}$, with $n_k = 1/R_{\text{QPC0}}/0$ (at zero temperature) for k below/within/above the detector voltage window $E_F < E(k) < E_F + eV_{\text{det}}$. It now becomes possible to express the expectation value of the many-body operator, $e^{i\delta\hat{\phi}}$, in terms of a determinant involving the matrices $w_{k'k}$ and $n_{k'k}$ (ref. 16):

$$\langle e^{i\delta\hat{\phi}} \rangle = \det[1 + (e^{i\omega} - 1)n]. \quad (5)$$

Equation (5) can be evaluated numerically. It converges in the limit of a large normalization volume and large upper or lower cutoffs in k . However, this expression leads to a suppressed visibility even at zero temperature and $V_{\text{det}} = 0$, which is an artefact of neglecting the Pauli blocking (which prevents the interfering electrons from scattering into occupied states below E_F). This can be cured approximately either by rescaling the visibility by a factor

independent of V_{det} (setting it to 1 at $V_{\text{det}} = 0$), or by restricting the matrix elements $w_{k'k}$ to transitions only within the detector voltage window. The latter approach provides a further simplification of equation (5) and yields equation (4) of the main text.

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Competing financial interests

The authors declare no competing financial interests.

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