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Molecular magnetic switch for a metallofullerene

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The endohedral fullerenes lead to well-protected internal species by the fullerene cages, and even highly reactive radicals can be stabilized. However, the manipulation of the magnetic properties of these radicals from outside remains challenging. Here we report a system of a paramagnetic metallofullerene $\text{Sc}_3\text{C}_2@\text{C}_{80}$ connected to a nitroxide radical, to achieve the remote control of the magnetic properties of the metallofullerene. The remote nitroxide group serves as a magnetic switch for the electronic spin resonance (ESR) signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ via spin-spin interactions. Briefly, the nitroxide radical group can 'switch off' the ESR signals of the $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety. Moreover, the strength of spin-spin interactions between $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and the nitroxide group can be manipulated by changing the distance between these two spin centres. In addition, the ESR signals of the $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety can be switched on at low temperatures through weakened spin-lattice interactions.

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Endohedral fullerenes are constructed by putting atoms or clusters inside fullerene cages, which isolates the internal species with environments, so even those high-reactive species can be well stabilized inside the fullerene cages^{1–6}. For example, paramagnetic endohedral fullerenes such as N@C₆₀ (refs 7–10), Sc@C₈₂ (refs 11,12), Y@C₈₂ (refs 13,14), Sc₃C₂@I_h-C₈₀ (refs 15–17) and so on, encapsulating radicals inside the fullerene cages, show also remarkable high stability, and they can be kept in air under room temperature for a long time, especially for endohedral metallofullerenes (EMFs). Considering the paramagnetic endohedral fullerenes usually show long electron spin relaxation and coherence times, they are expected to have potential applications in many fields such as spin labelling, spintronics, quantum computing and so on¹⁸.

For paramagnetic EMFs, the electronic spin resonance (ESR) technique is a powerful tool to detect the spin distributions and spin–nucleus couplings on internal species^{19–21}. By means of ESR, it was revealed that the magnetic property of the internal species can be roughly manipulated by changing the dynamic movement of internal species. For example, under room temperature the internal Y₂ cluster in Y₂@C₇₉N has a free rotation that leads to a symmetric ESR pattern, but along with the temperature decreasing, the free motion of Y₂ is hindered, leading to spin anisotropy and an asymmetric ESR pattern²². In addition, the spin characters and couplings in Sc₃C₂@C₈₀ was observed to change largely upon chemical modification of the fullerene cage due to the restricted Sc₃C₂ cluster¹⁷.

For better applying the paramagnetic EMFs in quantum information process and molecular devices, however, it is still a challenge to finely manipulate their magnetic property. Recently, Turro *et al.* chemically modified the H₂@C₆₀ with a nitroxide radical, and observed an indirect but strong magnetic communication between the electron spin of nitroxide paramagnet and the nuclear spin of encaged H₂ (refs 23–28). This finding provides a valuable clue for us on manipulating the magnetic property of paramagnetic EMFs via a foreign paramagnet²⁹. Since a strong spin–spin interaction between the paramagnetic fullerene molecule and the paramagnet is expected, thus the magnetic property of paramagnetic EMFs may be controlled by the attached paramagnet.

Herein, we report detailed studies on the fine manipulation for paramagnetic Sc₃C₂@C₈₀ by connecting it with a paramagnet of nitroxide radical. The target system FSc₃C₂@C₈₀PNO• contains two kinds of spins localizing on Sc₃C₂@C₈₀ and nitroxide radical, respectively. The remote nitroxide group serves as a magnetic switch for the ESR signals of Sc₃C₂@C₈₀ through spin–spin interactions. The paramagnetic properties of the metallofullerene Sc₃C₂@C₈₀ can be delicately adjusted by changing the temperature, varying the distance between the two spin centres, or simply quenching the nitroxide radical.

Results

Preparation of metallofullerene and its derivatives. Metallofullerene Sc₃C₂@C₈₀ was synthesized by the Kräschmer–Huffman arc-discharging method³⁰ and isolated by multi-stage high-performance liquid chromatography. Two Sc₃C₂@C₈₀ derivatives, FSc₃C₂@C₈₀PNOH and FSc₃C₂@C₈₀PNO•, were first synthesized through a Prato reaction³¹, respectively, as shown in Fig. 1a,b. The structures and spin density distributions of FSc₃C₂@C₈₀PNOH and FSc₃C₂@C₈₀PNO• were calculated as well, as shown in Fig. 1c,d. The FSc₃C₂@C₈₀PNOH has one spin centre that is localized on Sc₃C₂@C₈₀, whereas the FSc₃C₂@C₈₀PNO• has two unpaired spins localizing on the Sc₃C₂@C₈₀ moiety and nitroxide radical, respectively.

The ESR analysis of metallofullerene with a nitroxide radical. Sc₃C₂@C₈₀ is a typical paramagnetic endohedral fullerene. As reported previously, the ESR spectrum of the pristine Sc₃C₂@C₈₀ shows a symmetric pattern with 21 resonant lines, however, those of Sc₃C₂@C₈₀ derivatives showed a distorted pattern with a greatly increased amount of resonant lines¹⁷. Therefore, ESR spectroscopy was first employed to reveal the electron spin characters of FSc₃C₂@C₈₀PNOH and FSc₃C₂@C₈₀PNO•.

As shown in Fig. 1e, the ESR spectrum of FSc₃C₂@C₈₀PNOH was measured and analysed. Because the I_h symmetry of Sc₃C₂@C₈₀ is broken down after chemical modification, the original three equivalent scandium nuclei (I_{Sc} = 7/2) are classified into two groups in FSc₃C₂@C₈₀PNOH, in which one group contains a single Sc nucleus (g = 1.9948, hyperfine coupling constants (hfcc) = 8.5 G), and the other group contains two equivalent Sc nuclei (hfcc = 5.0 G). In comparison, the previously studied Sc₃C₂@C₈₀ fulleropyrrolidine shows a similar ESR pattern with hfcc of 8.6 G (one Sc nucleus) and 4.8 G (two Sc nuclei), respectively¹⁷. These ESR results reveal that the FSc₃C₂@C₈₀PNOH has a same reaction site with that of Sc₃C₂@C₈₀ fulleropyrrolidine¹⁷.

However, the ESR study of FSc₃C₂@C₈₀PNO• showed only three resonant lines (g = 2.0026, a = 15.5 G) that are derived from nitroxide radical (I_N = 1), and the ESR signals of Sc₃C₂@C₈₀ moiety were not observed, as illustrated in Fig. 1f. The current results reveal that the ESR signals of Sc₃C₂@C₈₀ can be switched off by paramagnetic nitroxide radical through spin–spin interaction. Thus it is interesting that the nitroxide radical group can serve as a remote controller for the ESR signals of Sc₃C₂@C₈₀ moiety. Vividly, if the ESR signals of Sc₃C₂@C₈₀ moiety are regarded as an indicating lamp, the nitroxide radical group can switch it off.

To reveal the mechanism of how the ESR signals of Sc₃C₂@C₈₀ moiety are switched off by the nitroxide radical, we synthesized FSc₃N@C₈₀PNO• in a same way for comparison. Sc₃N@C₈₀ is a diamagnetic molecule, and no spin–spin interaction is expected for FSc₃N@C₈₀PNO•. The signal intensity of FSc₃N@C₈₀PNO• was observed to be stronger than that of FSc₃C₂@C₈₀PNO• at the same concentration (0.151 p.p.m.). That is to say, the spin–spin interactions weaken both of the ESR signals of Sc₃C₂@C₈₀ and nitroxide radical (Supplementary Fig. 1).

For FSc₃C₂@C₈₀PNO•, the spin–spin interactions can be expressed as below:

$$E_{\text{dip}} = \frac{\mu_0}{4\pi} g^2 \beta_e^2 \left[\frac{\mathbf{S}_1 \times \mathbf{S}_2}{r^3} - \frac{3(\mathbf{S}_1 \times \mathbf{r})(\mathbf{S}_2 \times \mathbf{r})}{r^5} \right] \quad (1)$$

Where E_{dip} is the energy of dipolar coupling, and r is the dipole–dipole distance.

In fact, the spin–spin interactions between Sc₃C₂@C₈₀ and nitroxide radical broaden the resonance lines and lower the line intensity in the meantime, so the line width (ΔH) is adopted to represent the ESR line intensity and interpret the spin–spin interactions^{32,33}.

In general, for paramagnetic molecules the ESR line width is inversely proportional to the relaxation time (T), including the spin–lattice relaxation time (T_1) and the spin–spin relaxation time (T_2):

$$\Delta H = \frac{h}{2\pi g \beta} \left(\frac{1}{T_1} \right) + \frac{h}{2\pi g \beta} \left(\frac{1}{T_2} \right) \quad (2)$$

Therefore, the strong dipole–dipole interactions between nitroxide radical and Sc₃C₂@C₈₀ in FSc₃C₂@C₈₀PNO• reduced the spin–spin relaxation time (T_2), resulting in decreased ESR signals of both Sc₃C₂@C₈₀ and nitroxide radical moieties.

Figure 1 | Magnetic switch for the ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$. (a) The structure of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$. (b) The structure of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$. (c) The calculated structure and spin density distributions of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$. (d) The calculated structure and spin density distributions of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$. (e) The ESR spectrum of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$ at 293 K in toluene. (f) The ESR spectrum of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ at 293 K in toluene. The lamps in e and f show the 'on' and 'off' states of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ ESR signals, respectively.

Note that the transformation between $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$ is reversible, that is, the nitroxide radical in $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ turns into the corresponding hydroxylamine derivative ($\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$) using *p*-toluenesulfonylhydrazide, and the $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$ can be back to $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ by means of oxidation with copper acetate (Supplementary Fig. 2). Therefore, the magnetic property of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ can be easily manipulated by a chemical method, and the remote nitroxide group serves as a switch in this process.

The distance-dependent ESR signals. On the basis of equation (1), the spin–spin interactions between $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and nitroxide radical moieties can be efficiently reduced by elongating their distance, thus two other $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and nitroxide radical derivatives, that is, $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-2}$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-3}$, were synthesized with longer distances between these two spin centres. The pulsed ESR measurements on $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-2}$ revealed that the T_2 of nitroxide radical becomes longer when the distance of these two spins increases (Supplementary Fig. 3). As shown in Fig. 2, it is obvious that the ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ also gradually boost up along with the distance increasing, and this process is like lighting a lamp and making it brighter.

However, the increased chain length between $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and nitroxide radical would result in a strengthened spin–lattice interaction, which will bring another line-broadening effect for both nitroxide radical and $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety. The spin–lattice relaxation time (T_1) can be expressed as below:

$$\frac{1}{T_1} = \frac{2\pi g\beta_e}{h} \left(B_x^2 + B_y^2 \right) \frac{\tau_c}{1 + \tau_c^2 \omega_s^2} \quad (3)$$

Where B_x^2 and B_y^2 are mean square amplitudes of the fluctuating fields along the *x*- and *y*-directions, and τ_c is the correlation time of the motion that causes the fluctuation. For molecules with a

spherical shape, τ_c in liquid solution corresponds to the rotational correlation time τ_r , which can be approximated by the Stokes–Einstein relation:

$$\tau_r = \frac{4\pi\eta a^3}{3k_B T} \quad (4)$$

Where *a* is the rotationally effective radius of the molecule, and η is the viscosity of the solvent. In liquids with a low viscosity ($\tau_c^2 \omega_s^2 \approx 1$), the T_1 is dependent of a^3 , and T_1 decreases with increasing the rotationally effective radius of the molecule. Therefore, increasement of the molecular size would shorten the T_1 , leading to weaker ESR signals.

Since $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$, $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-2}$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-3}$ are all rigid structural molecules, the intramolecular dipolar coupling strength (*D*) can be estimated, in which the coupling strength of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ with $r = 1.38$ nm was estimated to be about 24.3 MHz following the classical point dipole approximation, and those of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-2}$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-3}$ were estimated to be about 11.2 and 5.48 MHz, respectively.

The temperature-dependent ESR signals. It is known that the spin–lattice relaxation time (T_1) of unpaired spin is tightly related to the temperature. As expressed in equation (3), the decrease of temperature would reduce the B_x and B_y , and then increase the T_1 and lead to higher ESR line intensity. Therefore, the temperature-dependent ESR studies of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ were performed, as shown in Fig. 3. It can be observed that no ESR signal of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ was observed at 293 K, but since 253 K, the ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ appeared together with three strong resonant lines of nitroxide radical, and continuously increased along with the temperature further decreasing. Finally at 213 K, the ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety can be clearly observed. The electrostatic spin–phonon

Figure 2 | Distance-dependent ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ derivatives. (a–c) Structures of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$, $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-2}$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-3}$. (d–f) ESR signals of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$, $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-2}$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\text{-3}$ at 293 K. The lamps in d–f show the strengthened ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety along with the enlarged distance from nitroxide radical.

Figure 3 | Temperature-dependent ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ derivatives. The ESR spectra of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ at variable temperatures in toluene solution. The lamps represent the strengthened ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety along with the decreased temperatures.

interaction was also analysed for these temperature-dependent ESR spectra³⁴. As the lowest temperature in our system is 213 K, under this condition the $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ toluene solution is still in liquid state, so the spin–phonon interaction is rather small and negligible. From the temperature-dependent ESR spectra, it can be seen that the temperature also can light the signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety in $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ and make it brighter.

Table 1 | ESR data.

Temperature (K)	Line width of nitroxide (G)	Line width of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ (G)
293	3.17	—
253	3.08	2.63
233	2.47	2.37
213	2.35	2.13

ESR, electronic spin resonance.
The ESR spectra line width of nitroxide and $\text{Sc}_3\text{C}_2@\text{C}_{80}$ in $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ at variable temperatures.

Moreover, it should be noted that the ESR signals of the nitroxide were also enhanced. Therefore, along with the temperature decreasing, the prolonged T_1 enhances not only the ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety, but also those of the nitroxide radical (Supplementary Fig. 4). The line widths of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety and nitroxide radical are listed in Table 1.

Discussion

Through connecting the paramagnetic metallofullerene $\text{Sc}_3\text{C}_2@\text{C}_{80}$ with a nitroxide radical, we have realized the manipulation of ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$. The remote nitroxide group serves as a magnetic switch for ESR signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$, that is, the paramagnetic nitroxide group can ‘switch off’ the ESR

signals of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ moiety. It was revealed that the spin-spin interactions between $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and nitroxide radical play a key role in realizing this kind of magnetic switch. Moreover, through increasing the distance between $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and nitroxide radical, or decreasing the temperature, we can finely adjust the paramagnetic property of $\text{Sc}_3\text{C}_2@\text{C}_{80}$.

Such controllable paramagnetism and switchable ESR signals have potential applications in quantum information processing and molecular devices. For example, these magnetic molecules can be fabricated to a single molecule membrane for data storage considering their transferrable two electron spin states (0/1), which can be written and read out by means of scanning tunnelling microscope. In addition, these magnetic molecules can be utilized as a probe for the reaction transition state considering the susceptible magnetic switch of the nitroxide group.

Methods

Synthesis of $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and $\text{Sc}_3\text{N}@\text{C}_{80}$ derivatives. $\text{Sc}_3\text{C}_2@\text{C}_{80}$ and $\text{Sc}_3\text{N}@\text{C}_{80}$ were heated with *N*-ethylglycine and 2,2,6,6-tetramethylpiperidine-1-oxyl 4-formylbenzoate (1a) (Supplementary Fig. 5), which were synthesized as literature methods²⁴ at 120 °C to give corresponding fulleropyrrolidines with yields of nearly 50% in toluene solution for 15 and 50 min, respectively. Pure $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ and $\text{FSc}_3\text{N}@\text{C}_{80}\text{PNO}^\bullet$ were isolated by HPLC using Buckyprep column (Supplementary Fig. 6). $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\cdot 2$ and $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet\cdot 3$ were synthesized according to same procedure used for compound $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$, except 2,2,6,6-tetramethylpiperidine-1-oxyl 4'-formylbiphenyl-4-carboxylate (2a) and 2,2,6,6-tetramethylpiperidine-1-oxyl 4'-p-terphenyl-4-carboxylate (3a) (Supplementary Fig. 7) were used.

Synthesis of $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNOH}$. To a solution of ~0.5 mg of the nitroxide derivative $\text{FSc}_3\text{C}_2@\text{C}_{80}\text{PNO}^\bullet$ in ~2 ml toluene was added ~1 mg *p*-toluenesulfonohydrazide, and stirred under air for about 15 min.

Characterization of metallofullerene derivatives. Ultraviolet/visible-near-infrared spectra of purified metallofullerene derivatives (Supplementary Fig. 8) were collected on Lambda 950 UV/Vis/NIR Spectrometer (PerkinElmer Instruments). ¹H NMR spectra of $\text{FSc}_3\text{N}@\text{C}_{80}\text{PNOH}$ was measured in chloroform-*d* on a Bruker 600 MHz spectrometer (Supplementary Fig. 9).

ESR measurements of metallofullerene derivatives. ESR spectra were measured on a JEOL JEF FA200 X-band spectrometer (Supplementary Fig. 10). The samples were degassed and the oxygen was removed from the solutions. All of the samples are dissolved in toluene solution at the same concentration.

Calculations on $\text{Sc}_3\text{C}_2@\text{C}_{80}$ derivatives. Density functional theory calculations were investigated by Perdew, Burke and Ernzerhof/double numerical plus polarization using the DMol3 code in Accelrys Materials Studio^{35,36}.

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Author contributions

T.W. conceived and designed the experiments. T.W. and C.W. wrote the paper. Experiments were carried out by B.W. The ESR data were analysed by T.W. and B.W. Calculations were carried out by Z.Z. All authors discussed the results and contributed to manuscript preparation.

Additional information

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