

REVIEW ARTICLE

Transporting droplets through surface anisotropy

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This review article examines digital microfluidic systems that manipulate droplets through surface anisotropy. These systems are categorized as surface tension driven or contact line driven. Surface tension driven systems include electrowetting on dielectric, Marangoni flow on microheater arrays, and chemical gradient surfaces, whereas contact line driven systems include anisotropic ratchet conveyors, nanostructured Parylene ratchets, and tilted pillar arrays. This article describes the operating principles and outlines the fabrication procedures for each system. We also present new equations that unify several previous models of contact line driven systems. The strengths and weaknesses of each system are compared, with a focus on their ability to perform the generation, switching, fusion, and fission of droplets. Finally, we discuss current and potential future applications of these systems.

Keywords: droplets; digital microfluidics; electrowetting; thermocapillary forces; contact line oscillation

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INTRODUCTION

Droplet microfluidics is an interdisciplinary field focused on the transport of fluids in small discrete volumes rather than through continuous flow. Like conventional microfluidics, droplet microfluidics arose from advances in microelectromechanical system (MEMS). Early droplet microfluidic systems generated micro- or nano-scale droplets of one liquid by applying a shear flow with a second immiscible liquid, essentially creating an emulsion (typically water in oil)^{1–3}. The segmented flow of droplet microfluidics significantly reduces the volume and reactant quantity requirements compared to continuous flow systems, and it circumvents the issues of Taylor dispersion, solute–surface interactions and cross-contamination^{4–7}. For these reasons, droplet microfluidic platforms are ideal for many biological and chemical applications, such as high throughput screening^{8–10}, enzymatic and cell assays^{11–13}, and combinatorial chemistry^{14–16}. However, droplets in segmented flow systems are not individually addressed. More recently, an emerging area of microfluidics, described as “digital microfluidics” (DMF), has made it possible to manipulate individual droplets without microchannels or encapsulating liquids^{6,7}. Although current DMF systems operate with larger droplets at lower transport rates^{6,7}, these systems have the potential to meet a broad range of applications and fulfill a unique niche in the field of microfluidics. Many DMF systems can control droplets through anisotropies on the underlying substrate; these systems will be the focus of this review.

DMF SYSTEMS

Each droplet on a DMF system can be thought of as a unit or “bit” of information, and DMF systems can manipulate droplets independently. This is a significant advantage over segmented flow systems that control droplets in series through the flow of an immiscible liquid in microchannels^{6,7}. DMF systems covered in this review can continuously transport droplets by using surface anisotropy. This anisotropy can be actively created (e.g., by applying a voltage or shining a light source) or passive (e.g., by way of chemical gradients or surface texture). These anisotropies manipulate droplets either by modulating the surface tension or

by asymmetrically guiding the contact line (“footprint”) of the droplet.

Surface tension driven DMFs

The term “surface tension driven” was first used by Lee *et al.*¹⁷ in 2002 to describe DMF devices that transport liquid droplets by creating surface tension gradients within the droplet. These gradients induce flow patterns that move the droplet along a substrate. DMF systems have implemented this principle by using electrical or thermal energy, as well as chemical gradient surfaces.

Electrowetting

Electrowetting on dielectric (EWOD) was the first DMF platform and is one of the best-developed methods of manipulating droplets^{7,18}. EWOD controls droplets by using voltages applied at electrodes just below the substrate surface (Figure 1a). The applied voltage creates a charge at the interface between the droplet and the electrode. This charge causes a local reduction in surface tension (also referred to as interfacial energy) at the solid–liquid interface (γ_{SL}), which reduces the contact angle in this region^{17–19}. The surface tension is unchanged in the area where no voltage is applied, and the contact angle is larger in this region. The effect of voltage on surface tension is described by Lippmann’s equation (Equation (1))^{17–19}, in which ϵ and d are the permittivity and thickness of the insulator between the electrodes and the droplet, respectively, and V is the applied voltage.

$$\gamma_{SL}(V) = \gamma_{SL}(0) - \frac{\epsilon}{2d} V^2 \quad (1)$$

This local change in solid–liquid surface tension induces flow in the direction of lower surface tension, effectively wetting the surface where the voltage is applied. This wetting effect is described by a combination of Lippmann’s and Young’s Equations (Equation (2))^{17–19}, where θ_0 is the contact angle with no applied voltage, c is the capacitance per unit area of the electrode, and γ_{LG} is the surface tension at the liquid–gas interface.

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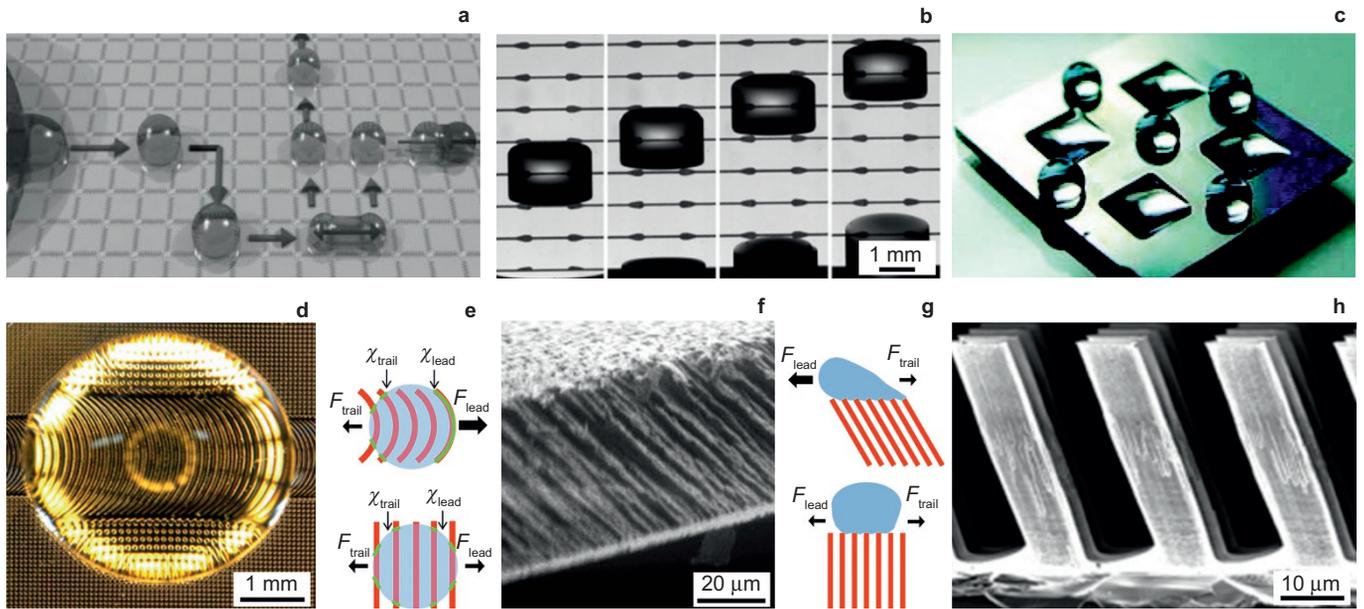


Figure 1 DMF systems transport droplets through surface anisotropy. These systems include (a) EWOD (adapted with permission from Ref 7 Copyright 2008 John Wiley & Sons), (b) microheater array (adapted with permission from Ref 20 Copyright 2003 IEEE), (c) chemical gradient (adapted with permission from Ref 21 Copyright 2006 American Chemical Society), (d) ARC (adapted with permission from Ref 22 Copyright 2012 John Wiley & Sons) with (e) schematic of operating principle, (f) nanostructured Parylene (adapted with permission from Ref 23 Copyright 2008 Elsevier) with (g) schematic of operating principle, and (h) TPA platforms (adapted with permission from Ref 24 Copyright 2014 John Wiley & Sons).

$$\cos \theta = \cos \theta_0 + \frac{1}{2\gamma_{LG}} cV^2 \quad (2)$$

This equation shows that the contact angle (θ) will decrease with increasing applied voltage (V). This makes the electrode surface more hydrophilic and “pulls” the droplet onto the surface over the electrode^{17–19}.

EWOD devices are traditionally fabricated by patterning metal electrodes (typically gold or chromium) on a glass or quartz wafer, and coating the electrodes with a hydrophobic, dielectric material, such as Teflon^{17–19}. A second wafer, typically etched or coated with a dielectric, is often placed on top to enclose the droplet^{17–19}.

Optoelectrowetting (OEW) devices rely on the same principles of EWOD, but the voltage is applied by focusing light on a photoconductive material beneath the droplet. OEW devices were first fabricated by patterning a layer of indium-tin-oxide (ITO) on a transparent glass slide, and connecting ITO electrodes to the biasing lines with strips of amorphous silicon (a-Si:H)²⁵. This layer is separated from the glass by a thin layer of silicon dioxide, and coated with Teflon, a hydrophobic dielectric material. A voltage is initially applied to the biasing lines, and illumination of the a-Si:H connection triggers the electrowetting effect at the ITO electrode²⁵.

Microheater arrays

Microheater arrays transport droplets through thermocapillary forces (Figure 1b). Because surface tension is a function of temperature, applying a thermal gradient to a droplet creates a surface tension gradient within the droplet, which results in Marangoni flow^{26,27}. The Marangoni flow opposes the viscous or “restoring” force caused by the surface that acts on the droplet, resulting in a net velocity (V). This relationship is described by Equation (3)^{26,27}:

$$V = V^* h_0 \left[\underbrace{\frac{1}{\gamma_{LG}} \frac{dS}{dT}}_{\text{Viscous force}} + \frac{3}{2} \underbrace{\frac{1}{\gamma_{LG}} \frac{d\gamma_{LG}}{dT}}_{\text{Marangoni effect}} \right] \frac{dT}{dx} \quad (3)$$

where V^* is the characteristic velocity (liquid–gas surface tension, γ_{LG} , divided by viscosity), h_0 is the height of the center of mass of the droplet, T is the prescribed temperature, and S , the spreading coefficient, is the relationship between solid–gas, γ_{SG} , solid–liquid, γ_{SL} , and liquid–gas surface tension ($S = \gamma_{SG} - (\gamma_{SL} + \gamma_{LG})$) observed on the substrate surface. Interestingly, the spreading coefficient shows little response to temperature changes when S is small ($dS/dT \approx 0$)^{26,27}. In this case, the Marangoni flow dominates and the droplet moves in the direction of higher surface tension (i.e., towards the cold side) because $d\gamma_{LG}/dT$ is negative^{26,27}. Microheater arrays make use of this phenomenon with periodically placed resistive heaters that essentially “push” droplets towards a cooler region when the substrate is heated²⁰.

Microheater arrays have been fabricated by patterning resistive heaters composed of titanium or chromium onto a rigid substrate (e.g., glass). Gold is used to form the leads to the heaters, and self-assembled monolayer patterns are used to prevent the droplets from spreading outside the path of resistive heaters²⁰.

Chemical gradients

The surface tension at the solid–liquid interface of a droplet is related to the contact angle (Figure 1c) through Young’s equation (Equation (4))²⁸.

$$\gamma_{SL} = \gamma_{SG} - \gamma_{LG} \cos \theta \quad (4)$$

The contact angle of a liquid droplet measures wettability, which is determined by the chemical composition of a surface²⁸. As a result of this relationship, a chemical gradient surface will cause the

spreading coefficient ($S = \gamma_{SG} - (\gamma_{SL} + \gamma_{LG})$) to vary spatially²⁶. In this case, Marangoni flow is not present, so the droplet will move in the direction of lower surface tension (i.e., towards the hydrophilic region)²⁶. This is similar to the effect described by Equation (3), but Equation (5) more accurately models the velocity of the droplet in this case²⁶.

$$V = \frac{h_0}{3\eta \ln(l_{\max}/l_{\min})} \frac{dS}{dx} \quad (5)$$

The log defines the limits of the flow region, where l_{\max} is the macroscopic cutoff, l_{\min} is a molecular size, and η is the viscosity of the droplet²⁶. The velocity of the droplet is directly proportional to the surface tension gradient (described by the change in the spreading coefficient along the surface, dS/dx) and, by extension, the wettability gradient. This concept was first used to show that a chemical gradient could transport a droplet uphill²⁹. However, the gradients used were static and could move droplets only a limited distance. Since then, droplets have been moved on wettability gradients created by selective exposure of photosensitive surfaces³⁰.

Fabrication of a photosensitive surface for inducing wettability gradients was first performed by coating a glass substrate with macrocycle amphiphile (O-carboxymethylated calix[4]resorcinarene)³⁰. More recently, photoreversible or “rewritable” photosensitive surfaces have been fabricated through layer-by-layer deposition of a polyelectrolyte (poly[allylamine hydrochloride]) and SiO₂ nanoparticles on negatively charged silicon wafers²¹, as well as by atomic layer deposition of zinc oxide films on glass and quartz wafers³¹. These systems can continuously drive droplet motion by maintaining surface energy gradients through active control of light beams.

Contact line driven DMFs

When vertical vibrations are applied to a droplet resting on a substrate, axisymmetric waves will form along the surface of the droplet³². Once they achieve sufficient amplitude, the vibrations cause the contact line (the perimeter of the droplet in contact with the substrate) to oscillate³². Each oscillation cycle is composed of two phases: an expansion phase and a contraction phase. A droplet oscillating on a homogeneous substrate enters the expansion phase when the contact line is at its smallest circumference. The contact line then advances in all directions throughout this phase. During the contraction phase, the contact line starts at its largest circumference and recedes until the next expansion phase begins. Contact line driven DMFs use surfaces that introduce an asymmetry to this oscillation cycle. Unlike surface tension driven DMFs, these systems move droplets through an imbalance of pinning forces on the edges of the contact line.

Anisotropic ratchet conveyors

Anisotropic ratchet conveyors (ARCs) use two sources of asymmetry: a heterogeneous surface pattern and a difference in pinning forces during dewetting compared to wetting^{22,33,34}. The ARC surface pattern consists of a path of arcs or curved “rungs” surrounded by a hydrophobic border (Figure 1d) and it can be texture or chemical based. In texture ARCs, the rungs are patterned mesas surrounded by a trench, and the hydrophobic region is created by an array of pillars^{22,33}. This topography induces a superhydrophobic state, or Cassie–Baxter state³⁵, as air is trapped under the droplet in the trenches between rungs. Chemical ARCs, by contrast, are composed of alternating hydrophobic and hydrophilic rungs on a flat surface³⁴. Applied vertical vibrations cause the contact line of a droplet to oscillate along the path of rungs. The contact line expands relatively equally in both directions, as the wetting process is less sensitive to ARC surfaces^{22,34}. However, during the dewetting process, when the

contact line recedes, the leading edge of the contact line experiences higher pinning forces than the receding edge due to the conforming shape of the rungs²². This phenomenon is described by Equation (6)³⁴, where χ_{lead} and χ_{trail} are the line fractions of the leading and trailing edges, respectively, that pin to the rungs, w is the width of the drop projected orthogonally to the pinning direction; and θ_1 and θ_2 are the equilibrium contact angles of the relatively hydrophilic rungs and relatively hydrophobic region between rungs, respectively.

$$F_{\text{Anisotropy}} = (\chi_{\text{lead}} - \chi_{\text{trail}})w\gamma_{\text{SL}}(\cos\theta_1 - \cos\theta_2) \quad (6)$$

The difference in pinning between the two edges results in a net force, especially during the contracting phase of oscillation^{22,34}. Equation (6) shows the difference between the forces on the leading edge (F_{lead}) and the trailing edge (F_{trail}) of the droplet. If we instead look at the ratio ($F_{\text{trail}}/F_{\text{lead}}$) of these opposing forces, we obtain Equation (7), which shows the four components that account for the anisotropy in ARC systems: the two contact angles θ_1 and θ_2 and the line fractions χ_{lead} and χ_{trail} on each edge of the droplet. The rung pattern causes the line fraction to be larger on the leading edge of the droplet, without which the forces would be equal on both edges of the droplet, and no net motion would occur (Figure 1e).

$$\frac{F_{\text{trail}}}{F_{\text{lead}}} = \frac{\chi_{\text{trail}}(\cos\theta_1 - \cos\theta_2)}{\chi_{\text{lead}}(\cos\theta_1 - \cos\theta_2)} = \frac{\chi_{\text{trail}}}{\chi_{\text{lead}}} \quad (7)$$

Both texture and chemical ARCs are fabricated by a single mask process. For texture ARCs, silicon or glass wafers are patterned with photoresist and etched with a deep reactive-ion etching (RIE) system to create the trenches and rungs. The ARCs are then coated with a hydrophobic silane to prevent the droplets from infiltrating the trenches^{33,34}. Texture ARCs can also be fabricated out of polydimethylsiloxane (PDMS) by using a silicon or glass wafer etched with the pattern negative as a mold³⁶. Chemical ARCs are fabricated using a similar process, but the wafers are first coated with a relatively hydrophilic layer of trimethylsilanol (TMS). The surface pattern is formed with photoresist and by etching exposed regions of TMS with oxygen plasma. The hydrophobic regions are then created with a silane (perfluorooctyltrichlorosilane) or with gold and dodecanethiol³⁴.

Nanostructured Parylene

Nanostructured Parylene ratchets are composed of a Parylene film made of tilted nanorods (Figure 1f)³⁷. Unlike ARCs, the asymmetry of the Parylene ratchets does not come from a surface pattern, but rather the angle of the nanorods. This geometry provides an anisotropic wetting effect on the droplet. The contact angles on the advancing (θ_A) and receding (θ_R) edges of a droplet on these nanofilms exhibit a directional dependence, which results in a net force³⁷. The ratio of forces on opposing edges of the droplet is shown in Equation (8)³⁷.

$$\begin{aligned} \frac{F_{\text{trail}}}{F_{\text{lead}}} &= \frac{\chi_{\text{trail}} \cos\theta_{R, \text{trail}} - \cos\theta_{A, \text{trail}}}{\chi_{\text{lead}} \cos\theta_{R, \text{lead}} - \cos\theta_{A, \text{lead}}} \\ &= \frac{\cos\theta_{R, \text{trail}} - \cos\theta_{A, \text{trail}}}{\cos\theta_{R, \text{lead}} - \cos\theta_{A, \text{lead}}} \text{ with } \frac{\chi_{\text{trail}}}{\chi_{\text{lead}}} = 1 \end{aligned} \quad (8\ddagger)$$

Droplet transport is induced by vertical vibrations, as in ARCs, which cause the contact line to oscillate³⁷. Interestingly, this

†The original subscripts PIN and REL³⁷ were changed to trail and lead, respectively, to maintain consistency between equations.

model is similar to the model for ARC systems shown in Equation (7). However, the line fraction is the same on both edges of the droplet, and the four components that account for the anisotropy in this system are all contact angles, two on each edge of the droplet. In this case, the net force results from the difference in the contact angles induced by the angle of the nanorods. Removing the tilt of the nanorods results in no net force, as the respective contact angles are the same on each edge of the droplet (Figure 1g). The similarity between the models for ARC and nanostructured Parylene systems presents the interesting possibility of combining the key features of each platform. For example, a tilt could be introduced to the rungs in the ARC system, or a nanostructured Parylene device could be fabricated with an ARC surface pattern. Such a combination could potentially provide dramatic improvements in the transport efficiency of these systems.

Nanostructured Parylene ratchets are fabricated using a technique called vapor deposition polymerization (VDP)²³. This deposition process is different from the commonly used process of physical vapor deposition, because the reactive monomers can only bind to the end of the polymer chain in VDP²³. This results in high aspect ratio Parylene rods with a nano-scale diameter. To create the angular structure of the Parylene ratchets, the substrate is tipped at an angle with respect to the incident vapor flux^{23,37} during the VDP.

Tilted pillar arrays

Tilted pillar arrays (TPAs) use silicon pillars arranged in an angular nano- or micro-structure to control the contact line pinning of a vibrated droplet (Figure 1h)²⁴. Much like nanostructured Parylene ratchets, the tilted pillars create differences between the contact angles on the leading and trailing edges, resulting in a net force that moves the droplet through cycles of contact line oscillation²⁴. Interestingly, droplets move with the direction of the tilt on nanostructured pillars, but move against the direction of the tilt on microstructured pillars²⁴. This observation suggests that it would be useful to combine tilted structures and surface patterns in contact line driven DMF systems. It would be interesting to understand how the direction of the tilt competes with the surface pattern. For example, if a tilted microstructure is in the same direction as the surface pattern, the net forces from the two anisotropies may cancel out and transport would not be possible. However, if the structures are combined so the net forces are in the same direction, then the transport efficiency may be increased, as described in the previous section.

Fabrication of TPAs begins with a silicon wafer substrate, the pattern for nanostructured TPAs is formed by using a metal, thin-film dewetting process^{24,38}, whereas the pattern for microstructured TPAs is formed by conventional photolithography. In both cases, the silicon wafer is etched using “glancing-angle” RIE, wherein the substrate is tilted at a 70° angle inside an RIE chamber^{24,38}. Once the tilted pillars have been etched into the substrate, the surface is then functionalized with a silane coating^{24,38}.

FUNCTIONS WITH DROPLETS

In addition to moving the droplets, DMF systems are also capable of performing specific functions on the droplets. Key functions that DMF systems perform are generation, switching, fusion, and fission.

Generation

Droplet generation describes the process in which individual droplets of a reproducible size are pulled from a large reservoir. This function allows for a large sample to be divided into many

smaller parts, which increases the number of reactions or tests that can be performed with a single sample. DMF systems also require a sample to be divided into small droplets. Without this function, a sample has to be distributed into smaller droplets through manual pipetting or with a separate apparatus. Thus far, droplet generation has been reported in electrowetting (EWOD³⁹ and OEWD²⁵) and ARC³⁶ systems.

EWOD systems generate droplets by first activating a pair of electrodes near the meniscus (edge) of a reservoir (Figure 2a). The meniscus is then pulled further from the reservoir by sequentially activating electrode pairs in line with the first pair³⁹. The location where the droplet separates from the reservoir determines the size of the generated droplet, but it is dependent on the initial shape of the meniscus, which is not predictable. However, the shape of the meniscus in the region near the first electrode pair can be stabilized with “side electrodes”. The side electrodes are placed on either side of the path formed by the initial electrode pairs. This design allows for relatively uniform droplets to be produced, with a volume of 2 to 3 μL ³⁹. OEWD systems perform this function based on similar principles, but with a light beam that activates electrodes adjacent to the reservoir, pulling a droplet on the active OEWD region²⁵. The sizes of the generated droplets are similar to the spot size of the beam²⁵.

Droplet generation for ARCs is possible using a pattern of concentric circles alternating between mesas and trenches³⁶. Like smaller droplets, the contact line of this reservoir will oscillate in response to vertical vibrations of sufficient amplitude. As the vibration amplitude is increased over a second threshold, the morphology of the contact line will shift from a nearly perfect circle to a trigonal planar shape^{32,36}. The alternating circle pattern distorts or modifies this contact line, which introduces instabilities at the “vertices” of the contact line³⁶. As the contact line expands across the surface texture, the vertices become pinned on the mesas. When the contact line retracts, a “pinch-off” occurs as the contact line is sheared over the trench (Figure 2b). With this design, the resulting droplet size is dependent on the initial size of the reservoir, and droplets as small as 0.1 μL have been generated³⁶.

Switching

Switching is the process of selectively moving a droplet from one path to another. This function is essential for sorting droplets and increasing the throughput of downstream processes. Switches are relatively simple to implement with DMF systems that move droplets by modulating surface tension because a surface tension gradient can be created in any direction within the droplet. On an EWOD or OEWD device, a droplet will move in the direction of an activated electrode. With an array of individually addressable electrodes, a droplet can be switched by activating an electrode next to the current path of the droplet^{7,25,39}. The droplet will be pulled onto this electrode and can be subsequently directed to a new path in any direction. Microheaters are also organized in individually addressable arrays to move droplets, but unlike EWOD electrodes, individual microheaters can only move droplets along a single axis⁴⁰. However, microheater switches have been achieved through the use of specially designed “intersections”. These intersections can switch a droplet by turning it 90 degrees onto a new track (Figure 2c). The intersection accomplishes this by creating a heated region that pushes the droplet in the direction of the new path while inhibiting movement on the direction of the original path⁴⁰. With photosensitive surfaces, as in electrowetting, droplets can be directed in any direction or pattern depending on the location of the light-activated chemical gradient relative to the droplet^{21,31}.

Implementing switches on contact line driven systems has been proven to be more of a challenge. ARC systems have been

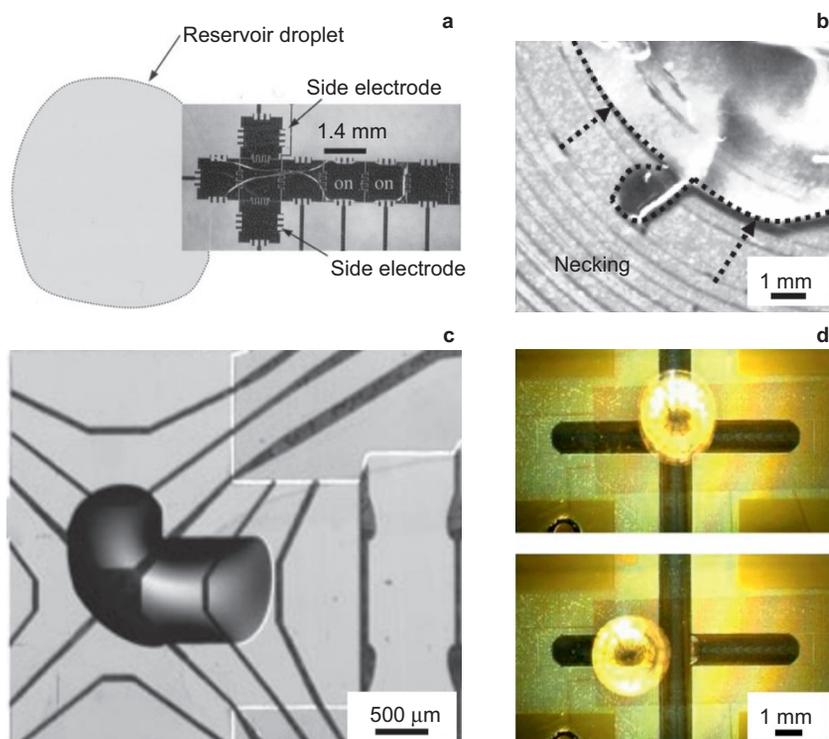


Figure 2 Droplet generation and switching. (a) EWOD (adapted with permission from Ref 39 Copyright 2003 IEEE) and (b) ARC systems can generate droplets by pulling liquid from a large reservoir (adapted with permission from Ref 36 Copyright 2008 IEEE). (c) Microheater intersections can also selectively switch droplets (adapted with permission from Ref 40 Copyright 2003 AIP Publishing LLC). (d) EWOD switches can be integrated in ARC devices (adapted with permission from Ref 41 Copyright 2009 IEEE).

developed that use EWOD electrodes to pull droplets onto a new path of rungs (Figure 2d)⁴¹. However, a switch that selectively moves droplets solely through geometric modulation of the contact line has yet to be developed. Interestingly, ARC systems are able to selectively move specific droplet sizes based on the frequency of applied vibrations⁴². This ability is likely present in nanostructured Parylene ratchets and TPAs as well, as the amplitude threshold for contact line oscillation is based on droplet volume⁴³. As mentioned previously, the direction of droplet motion is also dependent on the feature size of the TPAs. Such phenomena could eventually lead to the production of a switch that relies only on the geometric control of the contact line.

Droplet fusion

The fusion, or merging, of two droplets is a simple function for DMF systems to perform. With no encapsulating liquid, two droplets will spontaneously merge when they come in contact. Surface tension driven systems can accomplish this by creating surface tension gradients in two droplets that are directed towards each other, which causes the droplets to converge (Figure 3a)³⁹. Droplet fusion can be accomplished on contact line driven systems in a similar fashion by transporting droplets along two converging paths²². Fusion can also be performed along a single path with two different sized droplets (Figure 3b). This is accomplished by using a vibration frequency that modulates the contact line of only one droplet. One droplet will be transported along the path while the other remains static. When the moving droplet reaches the static droplet, the two will merge, and the resulting droplet can then be transported using a different vibration frequency²².

Although droplets can merge spontaneously, this does not guarantee that the resulting droplet will be homogeneously mixed. Homogeneity in the resulting droplet will eventually occur through diffusion, but this is a relatively slow process. This process can be hastened on electrowetting devices through a series of “mixing steps” wherein the droplet is moved back and forth between two electrodes or electrode pairs⁴⁴. This mixing step is not necessary on microheater arrays because the thermo-capillary convection currents that move the droplet also mix a fused droplet⁴⁰. Similarly, in contact line driven systems, the applied vertical vibrations drive the droplets with a cyclical motion, which also aids in mixing fused droplets²².

Droplet fission

Once two droplets have been fused and mixed, it may be desirable to divide the resulting droplet for multiple, subsequent uses. Therefore, the ability to perform droplet fission, or split the droplet into two or more smaller droplets, is very useful for DMF systems. Droplet fission has been accomplished with EWOD, OEW, and microheater systems. With EWOD, droplets can be split by activating electrodes on two opposing edges of the droplet, and by turning off the electrode beneath the center of the droplet. This pulls the edges of the droplet apart and causes the center of the droplet to neck, pinching apart the droplet (Figure 3a)³⁹. OEW devices perform droplet fission the same way, using two light beams to activate electrodes on opposing droplet edges²⁵. Droplet fission on microheater arrays occurs in a similar fashion, but microheaters in the center of the droplet are activated. This induces a diverging Marangoni flow, which thins the center of the droplet⁴⁰. Necking occurs as the droplet center thins, leading to eventual pinch-off and the formation of two new droplets (Figure 1c)⁴⁰.

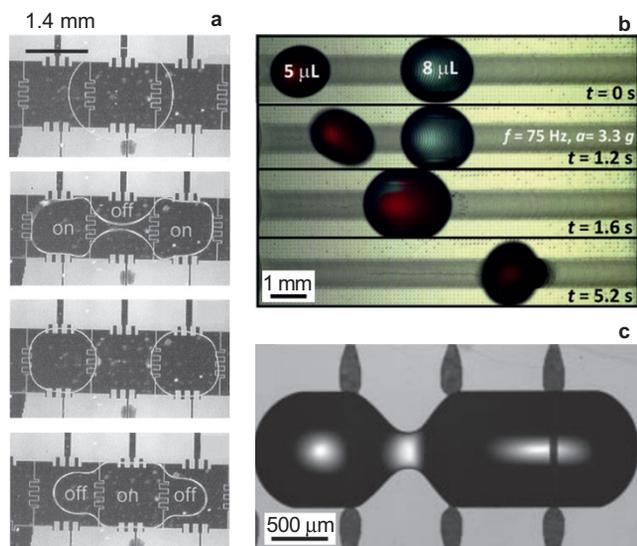


Figure 3 Droplet fission and fusion. (a) The sequence depicts the fission of one large droplet into two smaller droplets, and the recombination of the two droplets (adapted with permission from Ref 39 Copyright 2003 IEEE). (b) Two droplets are merged on an ARC path, and the resulting droplet is then transported (adapted with permission from Ref 22 Copyright 2012 John Wiley & Sons). (c) Microheater arrays can split droplets by creating thermal gradients (adapted with permission from Ref 40 Copyright 2003 AIP Publishing LLC).

Droplet fission has yet to be achieved with contact line driven systems. Transport in contact line driven systems is possible because the force acting on the leading edge of the droplet is larger than the force acting on the trailing edge. Splitting a droplet requires nearly equal forces acting on two opposing edges. These forces must also be large enough to overcome the surface tension of the droplet and persist long enough for the necking region to reach pinch-off⁴⁵. Electrowetting and microheater systems can maintain such a force for an indefinite amount of time, whereas contact line driven systems can only provide this force during one-half of the contact line oscillation cycle (the contracting phase). However, the time required for droplet pinch-off is dependent on the diameter of the necking region⁴⁵. Therefore, contact line driven systems could potentially perform droplet fission if a surface pattern could induce a contact line geometry with a small necked region. Additionally, it may be possible for a surface pattern to create a necking region that induces a resonance with an increasing transient response. In that case, the diameter of the neck would be reduced during each cycle of oscillation until pinch-off.

APPLICATIONS OF DMF SYSTEMS

DMF systems can be used for many of the same applications as segmented flow microfluidic systems. Although DMF systems operate with larger droplets and slower throughput, the ability to selectively control individual droplets makes DMF systems better suited for some applications (e.g., combinatorial chemistry⁷). Furthermore, these DMF systems move droplets through interactions with the substrate surface. Although this raises issues involving solute–surface interactions, it also provides an opportunity to address unique applications. For example, droplets can be moved to temperature sensitive regions on a chip or surface to aid in thermal management^{46,47}. The movement of droplets along a surface could be used to pick up and remove undesired particulates,

serving as a self-cleaning surface⁴⁸ or to transport desirable particulates deposited onto a surface⁴⁹.

EWOD was the first DMF system to be introduced, and it therefore has the most best-developed toolbox. EWOD devices have already been tested in applications such as cell isolation and analysis⁵⁰, bioassays⁵¹, and DNA analysis^{52,53}. Despite this success, a major criticism of EWOD technology is the complex circuitry that is required to address individual electrodes in an array^{6,7}. Driving droplets on microheater arrays and photosensitive chemical gradient surfaces also requires relatively complex circuitry or off-chip control systems. However, a more recent type of OEW system called single-sided continuous optoelectrowetting (SCOEW) uses a continuous layer of photoconductive amorphous silicon instead of a pixelated grid of electrodes to perform electrowetting. This system can also drive droplets with light produced from a liquid crystal display⁵⁴. Additionally, contact line driven systems only require an externally applied vibration to move droplets. Continued development of these more recently conceived systems could lead to a platform that can more efficiently address many applications. Hybrid DMF systems are also possible⁴¹. These would combine features of multiple DMF systems and could potentially lead to the realization of a robust platform for micro-total analysis and lab-on-a-chip systems.

Another field of microfluidics, known as paper or paper-based microfluidics, involves transporting liquids using capillary forces (paper wicking) in channels defined by hydrophobic barriers^{55,56}. These devices can be made from relatively inexpensive materials and do not require external stimuli to function. These properties have attracted attention for applications requiring low-cost analytical platforms (e.g., health care diagnostics, environmental monitoring, food quality control, and forensics)^{55,56}. However, paper microfluidics are limited by poor efficiency; large portions of a sample are often retained in the channel or lost to evaporation during transport⁵⁶. Many paper microfluidic systems also use a colorimetric detection method that usually has a high limit of detection⁵⁶. This approach prevents these systems from accurately measuring solutions with low analyte concentrations. The complexity of most surface tension driven DMF systems may become cost- or resource-prohibitive, but SCOEW shows great potential for low-cost diagnostic and analytic systems. The low power requirements and ability to drive droplets with a liquid crystal display make this system a good candidate for portable diagnostic systems. Additionally, contact line driven DMF systems require only a passive surface architecture and a vertical vibration. These surfaces can be made from many different materials, and the required vibrations are within the parameters of commercially available speakers found in cellular phones^{22,24,49}, which have become virtually ubiquitous and can be easily transported. DMF systems transport droplets as an entire self-contained entity, which avoids issues of sample retention experienced by paper microfluidic systems. Additionally, this property allows for integration with many detection systems, as the entire droplet can be transported onto a desired detection region. Implementing sensitive detectors with DMF systems could enable applications requiring detection limits outside the capabilities of paper microfluidic devices. These DMF systems could serve as a more efficient platforms for low-cost analytical devices, or they could be combined with paper microfluidic detection zones to improve the transport efficiency.

CONCLUSION

As of today, four main categories compose the field of microfluidics: continuous flow, segmented flow, paper, and digital. Continuous flow systems transport liquid through channels in a laminar flow regime. Segmented flow systems reduce reactant quantities by transporting fluid in discrete droplets separated

by an immiscible fluid. Both continuous and segmented flow systems can operate at very high throughputs but require complicated pump systems. Paper microfluidics provide a simpler, inexpensive platform, but are limited by poor transport efficiency. DMF systems fill the gap between segmented flow and paper systems, operating with larger droplets and slower throughputs than segmented flow systems but with higher efficiency than paper systems. DMF systems also have the unparalleled ability to selectively control individual droplets.

DMF systems can transport and manipulate droplets through surface anisotropy. EWOD, microheater arrays, and chemical gradient systems move droplets by modulating surface tension, whereas ARC, nanostructured Parylene, and TPA systems guide the contact line of a droplet. Both surface tension and contact line driven systems can perform droplet generation and fusion. However, only surface tension driven systems can currently perform switching and fission functions. Developing methods to perform these functions on contact line driven systems is still an open problem for researchers in the field.

In the future, the ability of DMF systems to individually control and selectively perform functions on droplets will be essential for building true lab-on-a-chip systems. DMF devices will also serve as platforms for low-cost analytical systems, because their transport efficiency and capacity for digital integration can be useful in applications outside the capabilities of paper microfluidic systems.

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COMPETING INTERESTS

The authors declare no conflict of interest.

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