

ORIGINAL ARTICLE

Highly efficient hybrid warm white organic light-emitting diodes using a blue thermally activated delayed fluorescence emitter: exploiting the external heavy-atom effect

Dongdong Zhang, Lian Duan, Yunge Zhang, Minghan Cai, Deqiang Zhang and Yong Qiu

To attain high efficiencies in hybrid white organic light-emitting diodes (WOLEDs), mutual quenching of the fluorophors and phosphors should be minimized. Efforts have been devoted to reducing the triplet quenching of phosphors; however, the quenching of fluorophors by the external heavy-atom effect (EHA) introduced by the phosphors is often ignored. Here, we observed that conventional fluorophors and fluorophors with thermally activated delayed fluorescence (TADF) behave differently in the presence of EHA perturbors. The efficiencies of the conventional fluorophors suffer greatly from the EHA, whereas the TADF fluorophors exhibit negligible changes, which makes TADF materials ideal fluorophors for hybrid devices. WOLEDs using a blue TADF fluorophor and an orange phosphor achieve a maximum forward viewing external quantum efficiency of 19.6% and a maximum forward viewing power efficiency of 50.2 lm W^{-1} , among the best values for hybrid WOLEDs. This report is the first time that the EHA effect has been considered in hybrid WOLEDs and that a general strategy toward highly efficient hybrid WOLEDs with simple structures is proposed.

Light: Science & Applications (2015) 4, e232; doi:10.1038/lsa.2015.5; published online 2 January 2015

Keywords: external heavy-atom effect; hybrid white OLEDs; thermally activated delayed fluorescence

INTRODUCTION

Recent conceptual advancements have led to many exciting approaches to white organic light-emitting diodes (WOLEDs).^{1–4} For practical use, high efficiency and long lifetime are the most critical parameters. Blue phosphors can obtain high luminous efficiency but suffer from short lifetime, while blue fluorophors are exactly the opposite.^{5,6} Efforts have been devoted to resolve the tradeoff between the high efficiency and long lifetime by combining the advantages of phosphors and fluorophors, that is, by fabricating hybrid WOLEDs with blue fluorophors and green/red or orange phosphors.^{6–9}

To attain high efficiencies in hybrid WOLEDs, mutual quenching of the fluorophors and phosphors should be minimized. It has been long realized that the quenching of the green and red phosphorescence by the low triplet energy fluorophors is one of the main losses of excitation energy in these devices.⁶ Therefore, a ‘triplet harvesting’ strategy by doping a blue fluorophor and red/green phosphors in separated regions of high-triplet-energy host materials has been proposed^{1,10} and is further improved by incorporating blue fluorophors with high triplet energy (T_1) to prevent the undesirable ‘back energy transfer’ from the phosphors to the blue fluorophors.^{2,7,8}

However, the heavy metal in the phosphors introduce an external heavy-atom effect (EHA) in the emissive layer (EML) and may lead to

quenching of the fluorescence due to the increased singlet (S_1)–triplet (T_1) intersystem crossing (ISC) rate (k_{ISC}) constant and the radiative or non-radiative decay rate of triplet excitons (k_p) of the fluorophors.^{11–13} This process could be another possible mechanism of energy loss in hybrid WOLEDs, especially for the WOLED with a single EML, even though this process has been ignored in previous studies.^{1,2,14} The quenching process between the fluorophor and the EHA perturber is rather complicated and not fully understood.¹² Multiple factors would affect the fluorescence quenching, such as the temperature, the distance between the fluorophor and the EHA perturber, and the nature of the fluorophor and perturber, respectively.^{12,15} Berberan-Santos *et al.* reported that even though the intrinsic quenching constants are very similar for phenanthrene-iodide and C70-bromobenzene systems, for phenanthrene-iodide, the quenching effect is high, while for C70-bromobenzene, the quenching effect is low as the competing decay processes are fast.¹⁵ The same group of authors also observed that both k_{ISC} and the rate of reverse ISC (k_{RISC}) of C70 are significantly improved when an EHA perturber is introduced, as the singlet–triplet split (ΔE_{ST}) of C70 is very small.¹¹ Notably, the ratio of the thermally activated delayed fluorescence (TADF) to the prompt fluorescence is higher than the unperturbed C70 at elevated temperatures,¹¹ suggesting that the loss in TADF is less

significant than that in the prompt fluorescence, most likely due to the existence of the competing $T_1 \rightarrow S_1 \rightarrow S_0$ channel even though the radiative decay rate of $C70 S_1$ is low.¹⁶ Therefore, a conventional fluorophor with a large ΔE_{ST} and a TADF fluorophor would behave differently in the presence of an EHA perturber, as shown in Supplementary Fig. S1a and S1b. It is conceivable that recovery of the total photoluminescence (PL) efficiency in the presence of an EHA perturber might be possible for a fluorophor with both a small ΔE_{ST} and a high radiative decay rate of S_1 .

Only recently, Adachi *et al.* developed efficient pure organic TADF materials with small ΔE_{ST} , while maintaining a high radiative decay rate (k_r) by carefully tuning the overlap of the highest occupied molecular orbital and the lowest unoccupied molecular orbital.^{17–19} Therefore, it would be important to understand how a TADF fluorophor would perform in the presence of a phosphorescent EHA perturber.

Here, the effect of phosphorescent EHA perturbers on the PL behaviors of a TADF fluorophor is studied for the first time. It is observed that the TADF fluorophor is much less sensitive to the EHA perturber than the conventional fluorophors due to the almost unchanged k_{ISC} and enhanced k_{RISC} . By doping a yellow phosphor and a blue fluorophor with TADF into a single EML, hybrid WOLEDs with a maximum forward-viewing external quantum efficiency (EQE_{max}) of 19.6% and a maximum forward-viewing power efficiency (PE_{max}) of 50.2 $lm W^{-1}$ are achieved. The total EQE_{max} of 33.3% and total PE_{max} of 85.3 $lm W^{-1}$ were also recorded using an integrating sphere. The high efficiency is achieved by minimizing the possible quenching processes between the phosphor and fluorophor.

MATERIALS AND METHODS

The TADF fluorophor, 1,4-dicyano-2,3,5,6-tetrakis(3,6-diphenylcarbazol-9-yl)benzene (t-Bu4CzTPN), was synthesized according to the method reported in the literature,¹⁹ with 3,6-diphenylcarbazol being replaced by 3,6-di-tert-butylcarbazole. Dipyrzino[2,3-f:2',3'-h]quinoxaline-2,3,6,7,10,11-hexacarbonitrile (HATCN) was purchased from LG Chem China Investment Co., Ltd. (Beijing, China), while *N,N'*-bis(1-naphthalenyl)-*N,N'*-diphenyl-[1,1'-biphenyl]-4,4'-diamine (NPB) was purchased from Beijing Visionox Technology Co. Ltd. (Beijing, China). *N,N,N'*-tris(4-(9-carbazolyl)phenyl)amine (TCTA), 1,3-bis(9*H*-carbazol-9-yl)benzene (mCP) and 3-(1,1'-biphenyl-4-yl)-5-(4-*tert*-butylphenyl)-4-phenyl-4*H*-1,2,4-triazole (TAZ) were purchased from Jilin Optical and Electronic Materials Co. Ltd. (Changchun, Jilin province, China). The 2CzPN was purchased from Xi'an Polymer Light Technology Co. Ltd. (Xi'an, Shanxi province, China).

PL characterization

Ultraviolet-visible absorption spectra were recorded using an Agilent 8453 spectrophotometer. The films were spin coated on the quartz substrate using dichloromethane as the solvent, which was volatilized before the measurement. The concentration of mCP was 10 $mg mL^{-1}$. The thicknesses of the films have little effect on their PL characters. The PL efficiencies and spectra of the films were measured using an absolute photoluminescence quantum yield measurement system (Hamamatsu C11347). The PL transient decay curves of the films were measured using a transient spectrometer (Edinburg FL920P).

Device characterization

The device structure of the hybrid WOLEDs was ITO/HATCN (5 nm)/NPB (40 nm)/TCTA (10 nm)/mCP (11 nm): 6 wt-%, 10 wt-% and 20 wt-% 2CzPN: 0.2 wt-% PO-01/TAZ (40 nm)/LiF (0.5 nm)/Al (150 nm). The OLEDs were fabricated by thermal evaporation under high vacuum ($\sim 7 \times 10^{-4}$ Pa) onto clean ITO-coated glass substrates. The forward-viewing electrical characteristics of the devices were measured

using a Keithley 2400 source meter, while the total ones were measured using an integrating sphere (Bluefly illumia 610). The electroluminescence spectra and luminance of the devices were obtained on a PR650 spectrometer. All the device fabrication and characterization steps were performed at room temperature under ambient laboratory conditions. For measurement of the transient electroluminescence characteristics, short-pulse excitation with a pulse width of 15 μs was generated using an Agilent 8114A pulse generator. The amplitude of the pulse was 10 V, and the baseline was -3 V. The period was 50 μs , the delayed time was 25 μs and the duty cycle was 30%. The decay curves of the devices were detected using an Edinburg FL920P transient spectrometer.

RESULTS AND DISCUSSION

To compare the effect of EHA on the TADF materials and conventional fluorophors, a TADF fluorophor (t-Bu4CzTPN) and two conventional fluorophors, bis(10-hydroxybenzo[h] quinolinato)beryllium complex (Bebq₂) and tris(8-hydroxyquinolino)aluminum (Alq₃), were selected. Here, Alq₃ was selected as a widely used conventional fluorophor, while Bebq₂ was selected because it possesses high triplet energy with a small ΔE_{ST} of 0.2 eV,²⁰ but not sufficient for TADF emission. The emission of the latter was approximately 490 nm, which can represent the high triplet blue fluorophors used in hybrid WOLEDs. Dichloro(1,5-cyclooctadiene)platinum(II) (Pt(COD)Cl₂), iridium(III) [bis(4,6-difluorophenyl)-pyridinato-N,C^{2'}]picolinate (Firpic) and tris(phenylpyrazole)iridium (Ir(ppz)₃) were used as EHA perturbers for the fluorophors. The molecule structures are shown in Figure 1a. To solely determine the effect of the EHA on the PL characteristics of the fluorophors, the excitation wavelength should be carefully selected to ensure that in the experiment, the phosphors cannot be excited and act only as EHA perturbers. Therefore, excitation wavelengths of 490 nm, 430 nm and 430 nm were selected for Firpic, Ir(ppz)₃ and Pt(COD)Cl₂, respectively, as observed in Figure 1b and 1c. To study the EHA in a rigid medium,¹³ mCP was used as a host here. The absorption spectrum of pure mCP shown in Supplementary Fig. S2 indicates that the mCP cannot be excited with an excitation wavelength of 430 nm or 490 nm, which rules out the interaction between the mCP and EHA perturbers. The concentration of the fluorophors is fixed at 5 wt-%, while those of the phosphors are steadily increasing.

The efficiencies of the films of mCP: 5 wt-% fluorophors: *x* wt-% (0; 5; 10; 15; 20; 25) phosphors are measured. As observed in Figure 2a and Supplementary Fig. S3, compared with the sharply reduced efficiencies of Bebq₂ or Alq₃, with increasing concentration of the phosphor, the PL efficiencies of t-Bu4CzTPN are only slightly affected. The Stern–Volmer plots for quenching of these fluorophors by phosphors are presented in Figure 2b and Supplementary Fig. S3. The plots clearly demonstrate that the efficiency of Bebq₂ and Alq₃ are significantly quenched, whereas that of t-Bu4CzTPN is only slightly affected. All the spectra of these films are presented in Supplementary Fig. S4. The spectra of these fluorophors are identical as the concentration of the EHA perturbers increases, indicating that the effect of the polarity changing can be ruled out,¹⁸ and therefore, the change in the PL characteristics of these films is mainly due to the EHA effect.

For conventional fluorescence processes, the triplet excitons formed through ISC cannot be further recycled, and thus, the fluorescence efficiency (ϕ_F) can be expressed as²¹

$$\phi_F = \frac{k_r}{k_r + k_{ISC} + k_{nr}} \quad (1)$$

Here, when the EHA perturber is introduced, k_{ISC} is greatly increased, while k_r and the nonradiative rate of singlet (k_{nr}) are not

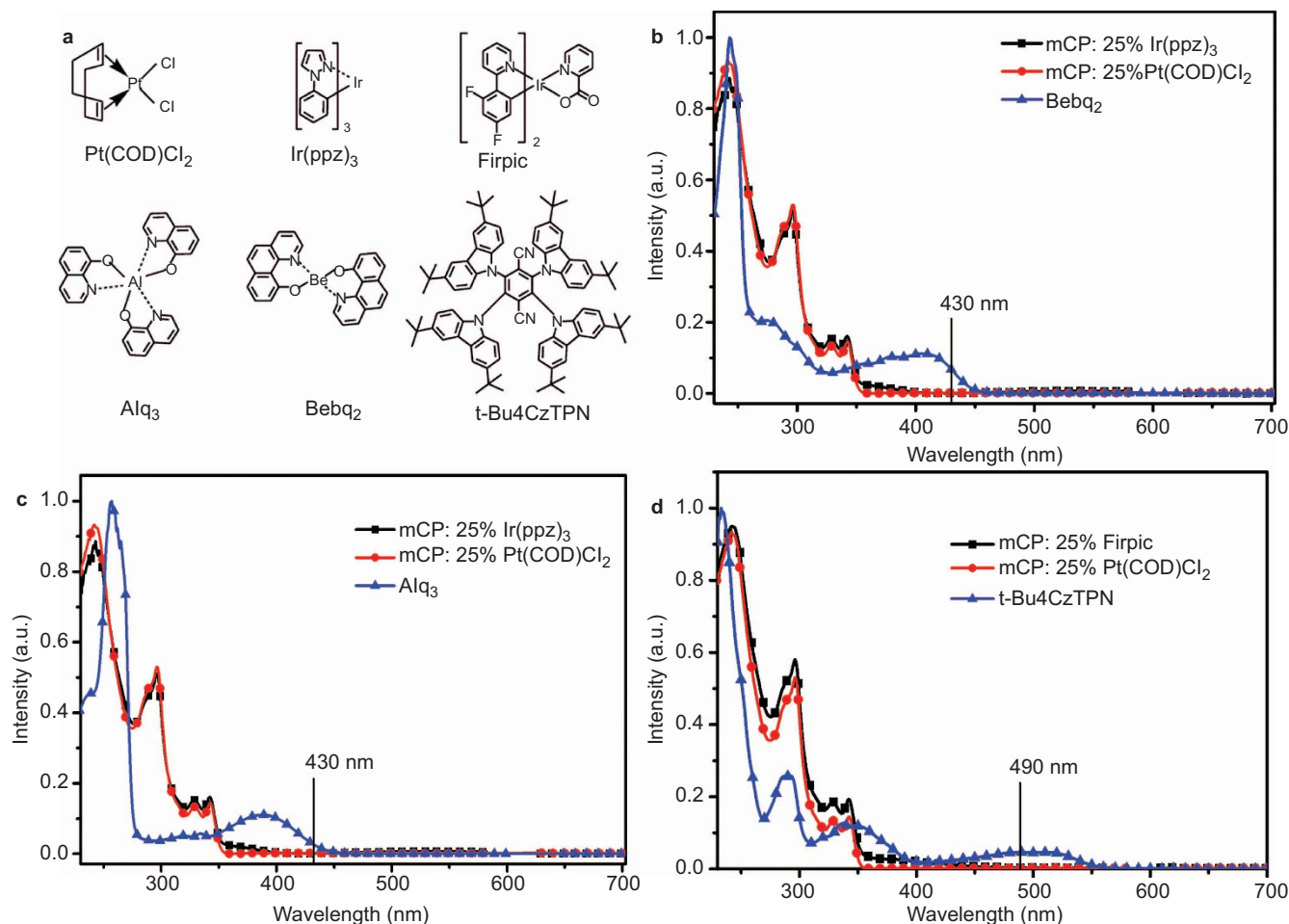


Figure 1 (a) The absorption spectra of Ir(ppz)₃, Pt(COD)Cl₂ in mCP: perturber films and Bebq₂ in dichloromethane solution. (b) The absorption spectra of Ir(ppz)₃, Pt(COD)Cl₂ in mCP: perturber films and Alq₃ in dichloromethane solution. (c) The absorption spectra of Firpic, Pt(COD)Cl₂ in mCP: perturber films and t-Bu4CzTPN in dichloromethane solution.

affected, resulting in significantly reduced ϕ_F . Therefore, the sharply reduced efficiency of Bebq₂ or Alq₃ is due to the increased k_{ISC} with the introduction of the EHA perturber; for a stronger EHA, more excitons are lost, resulting in significant efficiency quenching. As observed in Figure 2c and Supplementary Fig. S5c, with the increasing concentration of Ir(ppz)₃, the lifetime of the decay curves of Bebq₂ and Alq₃ are gradually decreased. The same trend was also observed when Pt(COD)Cl₂ was used as the EHA perturber (Supplementary Fig. S5). The lifetime of the singlet excitons is expressed by²²

$$\tau = \frac{1}{k_r + k_{ISC} + k_{nr}} \quad (2)$$

Because k_r and k_{nr} are not affected, k_{ISC} should be greatly increased, resulting in a reduced lifetime of the singlet excitons and confirming our hypothesis above. The increase of k_{ISC} can be calculated using Equation (2) and is shown in Supplementary Table S1. More triplet excitons will be formed in these fluorophors, and the radiative decay of the triplet excitons will be improved; therefore, such a system has been widely used to measure the triplet energies of fluorophors.²³ However, as observed in Equation (1), if k_r of the fluorophors is much higher than the EHA-enlarged k_{ISC} , the effect of EHA on ϕ_F will be insignificant.

For the TADF fluorophors, as observed in Figure 2d and Supplementary Fig. S5, the lifetimes of the prompt decay curves of t-Bu4CzTPN are almost the same with or without EHA perturbors

(17 ns for mCP: t-Bu4CzTPN: Firpic, see Supplementary Table S1), indicating that k_{ISC} of the fluorophors is less sensitive to the EHA perturbors. The reason may be attributed to the high intrinsic k_{ISC} of the TADF fluorophors without the perturber due to their small ΔE_{ST} . As reported, the EHA will only affect the pre-exponential factor but not the activation energy.¹² The increase of k_{ISC} when an EHA perturber is introduced is rather limited for the $S_1 \leftrightarrow T_1$ transition compared with the intrinsically high k_{ISC} of the TADF fluorophors. To further demonstrate that k_{ISC} is insensitive to the EHA perturber, the transient decay curves of t-Bu4CzTPN in toluene solution without degassing as the concentration of Firpic increased were also measured (Supplementary Fig. S6). The lifetime of the transient decay curves remains the same (10 ns, see Supplementary Table S1) as the concentration of Firpic increases, further demonstrating that the k_{ISC} of the fluorophors is less sensitive to the EHA perturber. The increase of k_{ISC} as observed for C70 derivatives¹¹ with the introduction of the EHA perturber, is the same as for t-Bu4CzTPN with a small ΔE_{ST} because the EHA manifests itself mainly by an increase of the $S_n \rightarrow T_n$ or $S_n \leftarrow T_n$ ISC rate constants. The small reduction of the PL efficiency of t-Bu4CzTPN can be attributed to the increased non-radiative decay rate of the triplet (k_{np}) when the EHA perturber is introduced.¹² In contrast, when Pt(COD)Cl₂ is used, the PL efficiency of t-Bu4CzTPN is slightly increased (Figure 2a) within a certain range of Pt(COD)Cl₂ concentration. The different behaviors can be interpreted

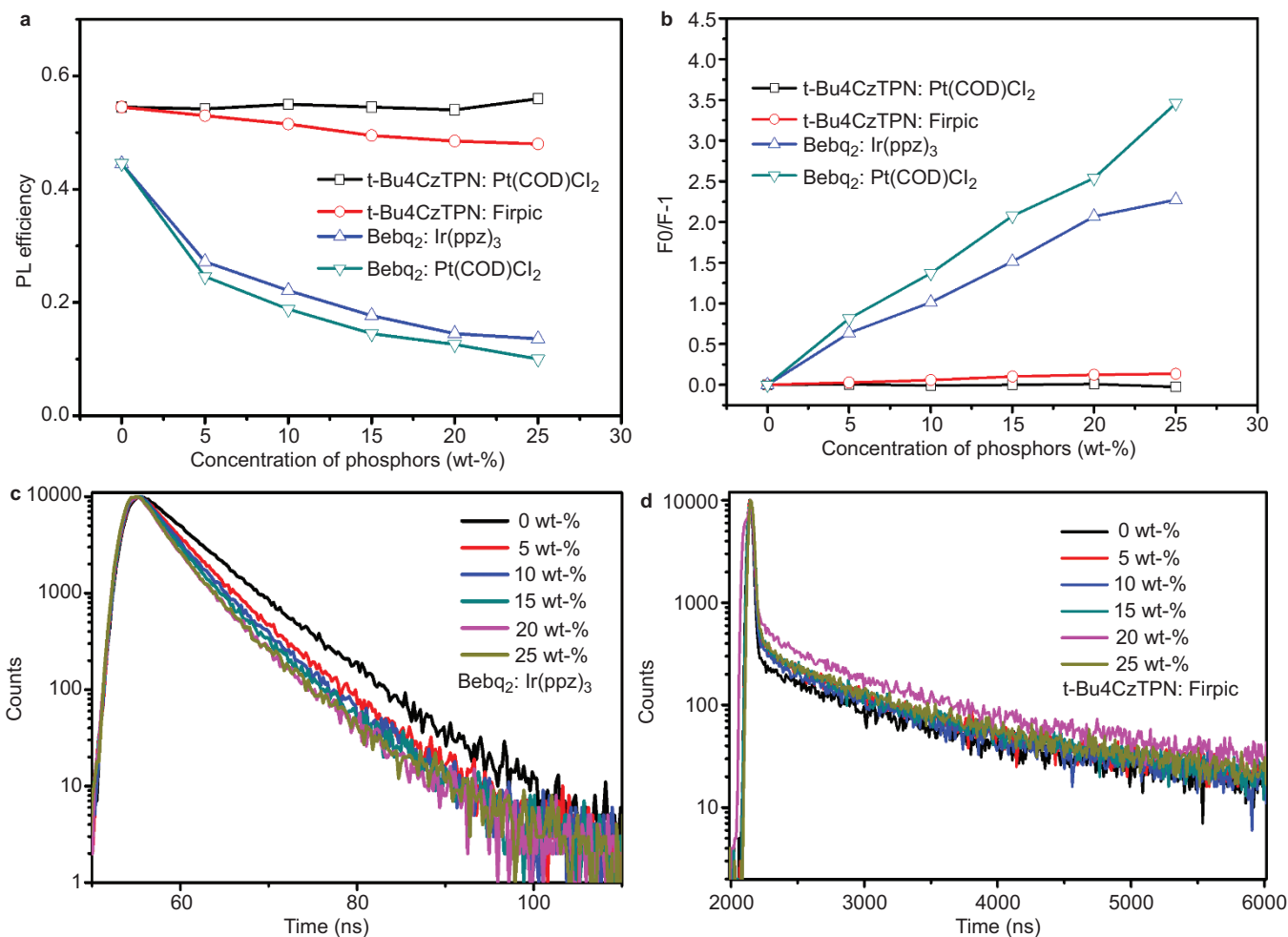


Figure 2 (a) The PL efficiencies of the films of mCP: fluorophors: EHA perturbors. (b) The Stern–Volmer plots for quenching of these fluorophors by EHA perturbors. F₀ stands for the PL efficiency of the fluorophors without EHA perturbors while F is the PL efficiency with EHA perturbors. (c) The transient decay curves of films of mCP: Bebq₂: Ir(ppz)₃ measured at 490 nm with an excitation wavelength of 430 nm. (d) The transient decay curves of films of mCP: t-Bu4CzTPN: Firpic measured at 576 nm with an excitation wavelength of 490 nm. EHA, external heavy-atom effect; PL, photoluminescence.

as the combinational effects of the processes of TADF fluorophors affected by the EHA perturbors. This result demonstrates that under some circumstances, if the k_{np} is not increased significantly when the EHA perturber is added, the PL efficiency of the TADF fluorophors can even be improved. The combinational effects of ISC, RISC and non-radiative decay processes of the TADF fluorophors under the effect of EHA perturbors causes the PL efficiencies of these TADF fluorophors to be almost unchanged or even to have the potential to be promoted under certain circumstances, indicating that TADF fluorophors and phosphors may offer an ideal solution toward high efficiency and stable white OLEDs.

The importance of constructing hybrid WOLEDs is the selection of guest emitters and its host. Here, the TADF blue fluorophor, 2CzPN, is selected due to its high PL efficiency and EQE_{max} of 13.6% in the optimized blue device.²⁴ A commonly used yellow phosphor, PO-01, is selected as the complementary phosphor.²⁵ The efficiencies of TADF materials are sensitive to the triplets of the hosts; hence, mCP with high triplet energy is selected.²⁴ Before the fabrication of devices, the PL efficiencies of films of mCP: 10 wt-% 2CzPN: *x* wt-% PO-01 (where *x*=0, 0.2, 0.4, 0.6, 0.8, 1.0) were measured, as observed in Supplementary Fig. S7b. The efficiencies of mCP: *x* wt-% PO-01 films

were also measured (Supplementary Fig. S7a). The ratio of the blue emission and orange emission can be calculated from the integral of their own spectra, and hence, the PL efficiency of the mCP: 2CzPN: PO-01 without EHA perturber can be calculated (Supplementary Fig. S7). As observed in Supplementary Table S2, the experimental PL efficiencies are higher than the theoretical ones calculated using the spectra of those films. The reason may be attributed to the increased PL efficiency of 2CzPN due to the EHA perturber of PO-01. The EHA effect may increase the ϕ_{RISC} of 2CzPN as we discuss above. For comparison, the interactions between Bebq₂ and PO-01 are also measured because the emission of Bebq₂ is close to 2CzPN. As observed in Supplementary Fig. S8, the experimental PL efficiencies of mCP: 10 wt-% Bebq₂: *x* wt-% PO-01 (where *x*=0, 0.2, 0.4, 0.6, 0.8, 1.0) are lower than the calculated ones, indicating that the EHA effect may decrease the PL efficiency of Bebq₂. Therefore, the positive synergy effects between 2CzPN and PO-01 made them a promising combination to fabricate high-efficiency WOLEDs.

The structure of the device fabricated is shown in Figure 3a. Both the fluorophor and phosphor are codoped into mCP as the EML, which is sandwiched between the electron-transporting layer of the TAZ film and the electron/exciton-blocking layer of the TCTA film. Meanwhile,

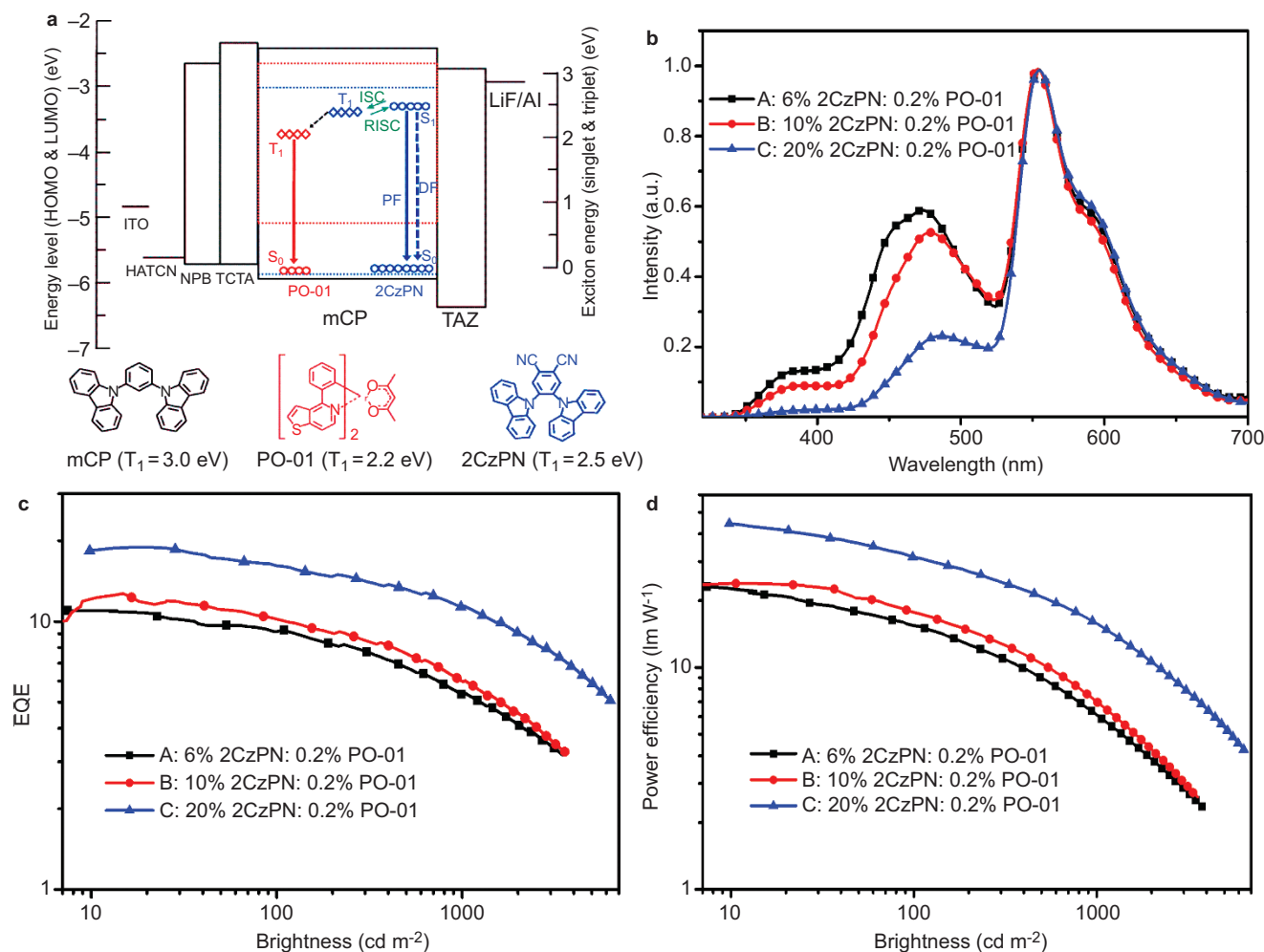


Figure 3 (a) The structure and energy diagram of the WOLEDs. (b) The spectra of the devices measured at 8 V. (c) EQE-brightness characteristics. (d) Power efficiency-brightness characteristics. EQE, external quantum efficiency; ISC, intersystem crossing; RISC, reverse intersystem crossing; WOLED, white organic light-emitting diodes.

NPB is introduced as the hole-transporting layer. Based on structural engineering to enhance the efficiency and spectral quality of the WOLEDs, optimization of the concentration of 2CzPN has been performed with that of PO-01 fixed at 0.2 wt-%. The concentrations of 2CzPN are 6 wt-%, 10 wt-% and 20 wt-% for device A, device B and device C, respectively.

All the devices exhibit warm white emission with the spectra covering the entire range from 400 nm to 700 nm, as observed in Figure 3b. With an increasing concentration of 2CzPN, the intensity of blue emission is gradually reduced due to the more efficient energy transfer from 2CzPN to PO-01. Here, the emission of the mCP host remains visible, indicating that energy transfer from mCP to 2CzPN or PO-01 is not complete. As observed in Supplementary Fig. S9c–9e, the blue emission intensities of the devices are gradually increased with increasing voltage, similar to the previous reported hybrid WOLEDs.⁸ The Commission International de L'Éclairage coordinates of the white devices changed from (0.38, 0.44) at 4 V to (0.33, 0.38) at 10 V for device A, (0.38, 0.46) at 4 V to (0.32, 0.39) at 10 V for device B and (0.43, 0.49) to (0.38, 0.45) for device C. The CRIs range from 53 at 4 V to 62 at 10 V for device A, 53 at 4 V to 62 at 10 V for device B and 53 at 4 V to 62 at 10 V for device C. For light applications, the most important reference light source is the warm white standard illumination A

with Commission International de L'Éclairage coordinates of (0.445, 0.405), exhibiting higher intensity in the orange region.⁶ In addition, candle-light-style OLEDs emitting yellowish orange light have been demonstrated to be a physiologically friendly light at night.^{26,27} Therefore, devices A–C are physiologically friendly light due to their relatively low intensities of blue emission.

The highest maximum forward viewing EQEs are 19.6%, 12.1% and 10.3% for devices with 20% (device C), 10% (device B) and 6% 2CzPN (device A), respectively (Figure 3c). The low efficiency with lower 2CzPN dopant concentration is partly due to the inefficient energy transfer. The device with 20% 2CzPN produces a maximum forward viewing power efficiency of 50.2 lm W⁻¹ (Figure 3d). Because illumination sources are typically characterized by their total emitted power, our WOLED exhibits a total EQE_{max} of 33.3% and a total PE_{max} of 85.3 lm W⁻¹ (measured using an integrating sphere), which roll off to 22.7% and 34.0 lm W⁻¹ at a high brightness of 1000 cd m⁻². To the best of our knowledge, these values are among the best results reported thus far for hybrid white devices (see Table 1) and are even comparable to the values of all phosphor white devices.⁴ More importantly, considering the PL efficiency of mCP: 20 wt-% 2CzPN: 0.2 wt-% PO-01 of 0.752 and assuming an optical outcoupling efficiency of approximately 20%–30%,²⁸ the EQE_{max} of 19.6% approaches the theoretical maximum value

Table 1 Device performances of some hybrid WOLEDs

	EQE ^a _{max}	EQE ^b _{max}	PE ^a _{max} (lm W ⁻¹)	PE ^b _{max} (lm W ⁻¹)	CIE (x,y)
Ref 2	—	20.3% ^c	—	57.6 ^c	(0.45,0.48) ^d
Ref 7	—	21.8%	—	57.3	(0.48,0.44) ^d
Ref 8	—	26.6%	—	67.2	(0.46,0.44) ^d
Ref 9	19.0%	—	41.7	—	(0.43,0.43) ^c
Device A	10.3%	—	23.0	—	(0.36,0.42) ^d
Device B	12.1%	—	24.5	—	(0.36,0.44) ^d
Device C	19.6%	33.3%	50.2	85.3	(0.42,0.48) ^d

Abbreviations: CIE, Commission International de L'Éclairage.

^a Forward viewing efficiency.

^b Total efficiency.

^c Measured at 1000 cd m⁻².

^d Measured at 100 cd m⁻².

(15%–22%), which indicates that the number of triplet excitons wasted is negligible. For the blue TADF material, there are two parallel approaches to minimize the triplet excitons lost, as demonstrated in Figure 3a. On the one hand, because of the higher triplet energy of 2CzPN compared with that of PO-01, energy transfer from the triplet of 2CzPN to that of PO-01 is energetically favorable, while the energy back transfer from PO-01 is prohibited. On the other hand, as discussed above, due to its small ΔE_{ST} , the triplet excitons in 2CzPN can be thermally upconverted into the emissive singlet ones, hence further reducing triplet loss. For hybrid WOLEDs using conventional blue fluorophors with a large ΔE_{ST} , unity IQE can only be achieved when all the triplet excitons are used in phosphors, which results in the restriction for spectra tuning because the ratio of excitons used in blue fluorophors must be lower than 25%. The method demonstrated here has the potential to allow free tuning of the emission spectra from cold white to warm white similar to the all phosphor devices, although only warm white emission is obtained for device A, and may achieve longer lifetimes than the all phosphors devices.²⁹ The mutual quenching of the fluorophors and phosphors in such hybrid devices are minimized, making TADF materials adequate for use in high-efficiency hybrid WOLEDs.

Though the efficiency roll-off of device C is rather high with only 58% of the highest EQE remaining at the brightness of 1000 cd m⁻², it remains less significant than the monochromatic device using 2CzPN as the emitter, as shown in Supplementary Fig. S10. This finding agrees well with former studies that singlet-triplet annihilation and triplet-triplet annihilation dominate efficiency roll-off for

2CzPN-based devices due to the long triplet lifetime of 2CzPN (273 μ s).³⁰ The narrow recombination zone due to the use of a hole-dominant host of mCP may also account for the efficiency roll-off of the device. To demonstrate our idea, we fabricated another WOLED using a bipolar host of 2,6-bis(3-(9H-carbazol-9-yl)phenyl)pyridine (26DczPPy) with the same structure as device C. As observed in Supplementary Fig. S11, the efficiency roll-off is significantly reduced, with 85% of the highest efficiency remaining at a brightness of 1000 cd m⁻², indicating that the performance of the hybrid WOLEDs can be improved by selecting adequate hosts.

To reveal the energy transfer process in the devices, the transient electroluminescence of devices A–C was measured. For comparison, the transient decay curves of devices only doped with 2CzPN and PO-01 were also measured. The transient decay curves for 2CzPN were measured at 480 nm for 2CzPN and at 560 nm for PO-01.

As observed in Figure 4a, in the turn-off region, transient overshoots are observed from the decay curves at 480 nm after the voltage pulse was turned off, evidently due to charge trapping in the blue dopant.³¹ Therefore, direct charge recombination on 2CzPN in the EML should occur. In addition, the recombination should also occur on the host because the emission of mCP is observed from the EL spectra, indicating that the energy of 2CzPN emission is also transferred from the host. Moreover, compared with blue devices using 2CzPN as the emitter, the ratio and lifetime of the delayed part of 2CzPN in the WOLEDs is increased rather than reduced. This result can be interpreted based on the EHA effect introduced by the

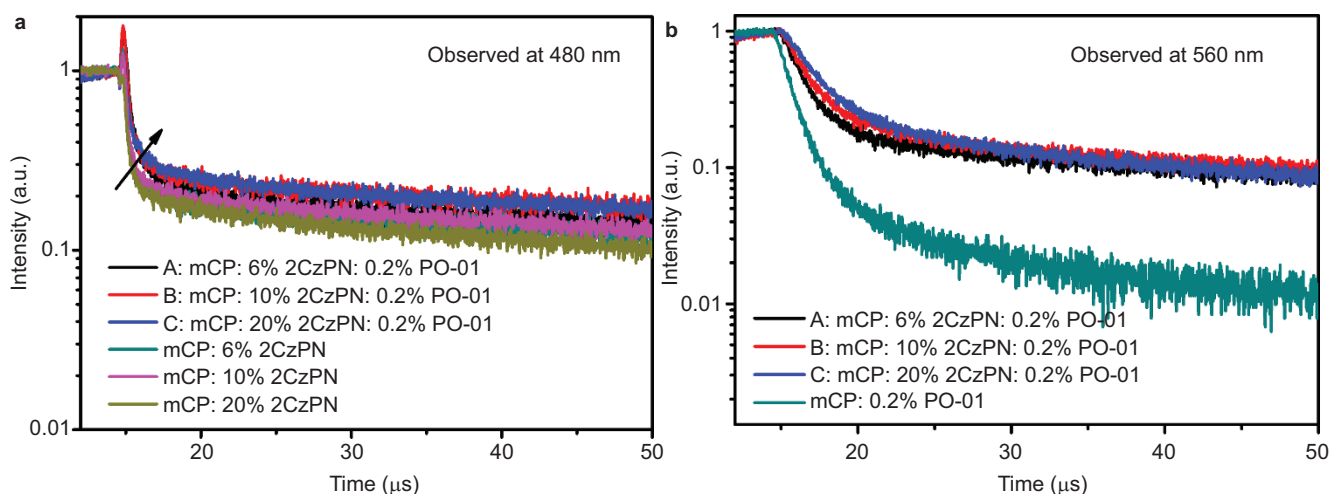


Figure 4 (a) The transient decay curves of devices observed at 480 nm. (b) The transient decay curves of devices observed at 560 nm.

phosphors, as discussed above, which improves the random change between the excited singlets and triplets.

For PO-01, the excitons can be generated in three ways, that is, energy transfer from mCP, energy transfer from 2CzPN, or direct recombination of charges on PO-01. As observed in Figure 4b, no overshoots are observed in the decay curves of PO-01 in the WOLEDs, indicating that direct charge recombination on PO-01 is negligible. The energy transferred from 2CzPN to PO-01 can be determined by the longer lifetime of PO-01 in WOLEDs compared with that of the device with only PO-01 as the emitter (Figure 4b). The long-lived triplet excitons of 2CzPN transferred to PO-01 results in the long lifetime of PO-01 emission. The energy transferred from mCP is also possible because excitons are also formed on the host, especially at higher driving voltages.

The operation lifetime of device C was measured at an initial brightness of 1000 cd m⁻² under constant current. As observed in Supplementary Fig. S12a, the lifetime of device C is approximately 140 min. The lifetimes of the monochromatic devices of 2CzPN or PO-01 in the same device architecture are even shorter. Because PO-01 has been demonstrated to be a stable emitter,²⁵ we believe that the instability of the mCP host should be one of the reasons for the short lifetime of the devices. Moreover, WOLED with 26DCzPPy as the host has a much longer lifetime than device C (Supplementary Fig. S12b), demonstrating that the stability of the hybrid OLEDs can be improved by selecting better host materials.

CONCLUSIONS

In conclusion, we have demonstrated that the PL efficiencies of the conventional fluorophors are effectively quenched by the EHA effect of the phosphors, while those of the TADF materials are less affected and may be slightly increased. Hybrid WOLEDs using a blue TADF fluorophor of 2CzPN and a yellow phosphor of PO-01 achieved an EQE_{max} of 19.6% and a maximum PE_{max} of 50.2 lm W⁻¹, corresponding to a total EQE_{max} of 33.3% and a total PE_{max} of 85.3 lm W⁻¹. Possible pathways of quenching, such as back energy transfer from the fluorophor to the phosphor and the quenching of the fluorescence of the fluorophor due to the presence of EHA, can be effectively blocked in this type of hybrid WOLEDs. This method opens a new avenue for achieving high-performance hybrid WOLEDs with simple structures.

ACKNOWLEDGEMENTS

We would like to thank the National Natural Science Foundation of China (Grant Nos. 51173096, 21161160447 and 61177023) for financial support.

- 1 Sun YR, Giebink NC, Kanno H, Ma BW, Thompson ME *et al*. Management of singlet and triplet excitons for efficient white organic light-emitting devices. *Nature* 2006; **440**: 908–912.
- 2 Schwartz G, Pfeiffer M, Reineke S, Walzer K, Leo K. Harvesting triplet excitons from fluorescent blue emitters in white organic light-emitting diodes. *Adv Mater* 2007; **19**: 3672–3676.
- 3 Su SJ, Gonmori E, Sasabe H, Kido J. Highly efficient organic blue-and white-light-emitting devices having a carrier- and exciton-confining structure for reduced efficiency roll-off. *Adv Mater* 2008; **20**: 4189–4194.
- 4 Lee SY, Yasuda T, Yang YS, Zhang QS, Adachi C. Luminous butterflies: efficient exciton harvesting by benzophenone derivatives for full-color delayed fluorescence OLEDs. *Angew Chem Int Ed* 2014; **53**: 6402–6406.
- 5 Baldo MA, O'Brien DF, You Y, Shoustikov A, Sibley S *et al*. Highly efficient phosphorescent emission from organic electroluminescent devices. *Nature* 1998; **395**: 151–154.
- 6 Schwartz G, Reineke S, Rosenow TC, Walzer K, Leo K. Triplet harvesting in hybrid white organic light-emitting diodes. *Adv Funct Mater* 2009; **19**: 1319–1333.

- 7 Zheng CJ, Wang J, Ye J, Lo MF, Liu XK *et al*. Novel efficient blue fluorophors with small singlet-triplet splitting: hosts for highly efficient fluorescence and phosphorescence hybrid WOLEDs with simplified structure. *Adv Mater* 2013; **25**: 2205–2211.
- 8 Ye J, Zheng CJ, Ou XM, Zhang XH, Fung MK *et al*. Management of singlet and triplet excitons in a single emission layer: a simple approach for a high-efficiency fluorescence/phosphorescence hybrid white organic light-emitting device. *Adv Mater* 2012; **24**: 3410–3414.
- 9 Sun N, Wang Q, Zhao YB, Chen YH, Yang DZ *et al*. High-performance hybrid white organic light-emitting devices without interlayer between fluorescent and phosphorescent emissive regions. *Adv Mater* 2014; **26**: 1617–1621.
- 10 Schwartz G, Fehse K, Pfeiffer M, Walzer K, Leo K. Highly efficient white organic light emitting diodes comprising an interlayer to separate fluorescent and phosphorescent regions. *Appl Phys Lett* 2006; **89**: 083509.
- 11 Baleizão C, Berberan-Santos MN. External heavy-atom effect on the prompt and delayed fluorescence of [70]fullerenes. *Chem Phys Chem* 2010; **11**: 3133–3140.
- 12 Kearvell A, Wilkinson F. Fluorescence quenching and external spin-orbit coupling effects. *Mol Cryst* 1968; **4**: 69–81.
- 13 Berberan-Santos MN. External heavy-atom effect on fluorescence kinetics. *PhysChemComm* 2000; **3**: 18–23.
- 14 Baldo MA, Thompson ME, Forrest SR. High-efficiency fluorescent organic light-emitting devices using a phosphorescent sensitizer. *Nature* 2000; **403**: 750–753.
- 15 Rae M, Fedorov A, Berberan-Santos MN. Fluorescence quenching with exponential distance dependence: application to the external heavy-atom effect. *J Chem Phys* 2003; **119**: 2223–2231.
- 16 Berberan-Santos MN, Garcia JM. Unusually strong delayed fluorescence of C₇₀. *J Am Chem Soc* 1996; **118**: 9391–9394.
- 17 Sato K, Shizu K, Yoshimura K, Kawada A, Miyazaki H *et al*. Organic luminescent molecule with energetically equivalent singlet and triplet excited states for organic light-emitting diodes. *Phys Rev Lett* 2013; **110**: 247401.
- 18 Ishimatsu R, Matsunami S, Shizu K, Adachi C, Nakano K *et al*. Solvent effect on thermally activated delayed fluorescence by 1,2,3,5-Tetrakis(carbazol-9-yl)-4,6-dicyanobenzene. *J Phys Chem A* 2013; **117**: 5607–5612.
- 19 Uoyama H, Goushi K, Shizu K, Nomura H, Adachi C. Highly efficient organic light-emitting diodes from delayed fluorescence. *Nature* 2012; **492**: 234–240.
- 20 Jeon WS, Park TJ, Kim SY, Pode R, Jang J *et al*. Ideal host and guest system in phosphorescent OLEDs. *Org Electron* 2009; **10**: 240–246.
- 21 Goushi K, Yoshida K, Sato K, Adachi C. Organic light-emitting diodes employing efficient reverse intersystem crossing for triplet-to-singlet state conversion. *Nat Photonics* 2012; **6**: 253–258.
- 22 Zhang DD, Duan L, Li C, Li YL, Li HY *et al*. High-efficiency fluorescent organic light-emitting devices using sensitizing hosts with a small singlet-triplet exchange energy. *Adv Mater* 2014; **26**: 5050–5055.
- 23 Tanaka I, Tabata Y, Tokito S. Observation of phosphorescence from tris(8-hydroxyquinoline) aluminum thin films using triplet energy transfer from iridium complexes. *Phys Rev B* 2005; **71**: 205207.
- 24 Masui K, Nakanotani H, Adachi C. Analysis of exciton annihilation in high-efficiency sky-blue organic light-emitting diodes with thermally activated delayed fluorescence. *Org Electron* 2013; **14**: 2721–2726.
- 25 Zhang DD, Duan L, Li YL, Li HY, Bin ZY *et al*. Towards high efficiency and low roll-off orange electrophosphorescent devices by fine tuning singlet and triplet energies of bipolar hosts based on indolocarbazole/1, 3, 5-triazine hybrids. *Adv Funct Mater* 2014; **24**: 3551–3561.
- 26 Jou JH, Hsieh CY, Tseng JR, Peng SH, Jou YC *et al*. Candle light-style organic light-emitting diodes. *Adv Funct Mater* 2013; **23**: 2750–2757.
- 27 Stevens RG, Brainard GC, Blask DE, Lockley SW, Motta ME. Breast cancer and circadian disruption from electric lighting in the modern world. *CA Cancer J Clin* 2014; **64**: 207–218.
- 28 Smith LH, Wasey JA, Barnes WL. Light outcoupling efficiency of top-emitting organic light-emitting diodes. *Appl Phys Lett* 2004; **84**: 2986–2988.
- 29 Nakanotani H, Masui K, Nishide J, Shibata T, Adachi C. Promising operational stability of high-efficiency organic light-emitting diodes based on thermally activated delayed fluorescence. *Sci Rep* 2013; **3**: 2127.
- 30 Masui K, Nakanotani H, Adachi C. Analysis of exciton annihilation in high-efficiency sky-blue organic light-emitting diodes with thermally activated delayed fluorescence. *Org Electron* 2013; **14**: 2721–2726.
- 31 Liu R, Gan ZQ, Shinar R, Shinar J. Transient electroluminescence spikes in small molecular organic light-emitting diodes. *Phys Rev B* 2011; **83**: 245302.



This work is licensed under a Creative Commons Attribution-NonCommercial-NoDerivs 3.0 Unported License. The images or other third party material in this article are included in the article's Creative Commons license, unless indicated otherwise in the credit line; if the material is not included under the Creative Commons license, users will need to obtain permission from the license holder to reproduce the material. To view a copy of this license, visit <http://creativecommons.org/licenses/by-nc-nd/3.0/>