

ORIGINAL ARTICLE

Extremely anisotropic single-crystal growth in nanotwinned copper

Chia-Ling Lu¹, Han-Wen Lin¹, Chien-Min Liu¹, Yi-Sa Huang¹, Tien-Lin Lu¹, Tao-Chi Liu¹, Hsiang-Yao Hsiao¹, Chih Chen¹, Jui-Chao Kuo² and King-Ning Tu³

By electroplating of nearly unidirectionally $\langle 111 \rangle$ -oriented nanotwinned and fine-grained Cu on a Si wafer surface followed by annealing at 400–500 °C for up to 1 h, we grew many extremely large $\langle 100 \rangle$ -oriented single crystals of Cu with sizes ranging from 200 to 400 μm . By patterning and annealing the nanotwinned Cu films, we grew an array of $\langle 100 \rangle$ -oriented single crystals of Cu with sizes ranging from 25 to 100 μm on Si. In comparison, single-crystal nano-wire growth is a one-dimensional anisotropic growth process, in which the growth along the axial direction is much faster than in the radial direction. We report here a bulk-type two-dimensional crystal growth of an array of numerous $\langle 100 \rangle$ -oriented single crystals of Cu on Si. This growth process has the potential for microbump applications in three-dimensional integrated circuit-packaging technology for hand-held consumer electronic products.

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INTRODUCTION

With the popularity of hand-held devices, mobile consumer electronic products are expected to undergo a very rapid growth in the global semiconductor market in the near future. As Si chip technology approaches the limit of Moore's law, the microelectronic industry is turning to packaging technology to extend the limit using three-dimensional integrated circuits (3D ICs), which, in essence, represents the merging of chip technology and packaging technology. On a limited chip area and in a tight device space, the best way to increase the density of high-performance transistor circuits is to stack the chips vertically, that is, in 3D ICs.

For a successful merger, precise fabrication of the microstructure in the packaging technology is required. In the trend of miniaturization of Si chip technology, the precise fabrication of hundreds of millions of transistors and their gates and contacts on a Si chip has been the foundation of Moore's law. One example is the use of lattice strain to increase the carrier mobility in Si channels. In 3D ICs, there are thousands of microbumps on a through-Si-via chip, and precise control of the microstructure of all the microbumps is required, which is a new paradigm in the microelectronic industry. Because the merging of chip technology and packaging technology is needed to extend Moore's law, it is crucial to have precise control of manufacturing in packaging technology.

We consider the diameter of a solder joint below. In flip chip packaging technology, the diameter is $\sim 100 \mu\text{m}$, and in microbumps for 3D IC packaging, the diameter is $\sim 10 \mu\text{m}$. When the diameter changes by 10 times, the volume changes by 1000 times, which indicates that if we assume that the grain size is $10 \mu\text{m}$, the

microbump will consist of one grain; however, the flip chip joint will contain 1000 grains. Clearly, we can assume an isotropic microstructure in the latter, and microstructure control is not a concern; however, in the former, the microstructure is anisotropic, for example, the orientation of a grain can change from one microbump to another microbump. The anisotropy tends to cause early failure of microbumps in the high-performance devices; thus, precise microstructure control is required.

Microstructure control starts from crystal growth, which is one of the most important topics in multi-disciplinary sciences because the crystal growth controls the properties of semiconductors, metals and ceramics in thermal processing.^{1–11} Crystal growth from a melt is known to be anisotropic, which is evidenced by facet formation. In the solid state, crystal growth is manifested by grain growth, which tends to be isotropic, especially in face-centered cubic metals. Many face-centered cubic metal thin films exhibit abnormal grain growth after annealing. Several researchers have reported that giant grains up to several hundred micrometers can be grown in Al,^{12,13} Ag^{14,15} and Cu thin films.^{16,17} However, we report here a new type of abnormal grain growth that results from the extremely anisotropic two-dimensional crystal growth of Cu via nearly unidirectionally oriented nano-twins; the growth in the vertical direction is much slower than that in the lateral directions. In comparison, single-crystal nano-wire growth is a one-dimensional anisotropic growth process in which the growth along the axial direction is much faster than that along the radial direction. Although graphene is grown using two-dimensional anisotropic growth, this process involves single atomic layer growth and hardly any growth in the vertical direction. We applied the extremely

¹Department of Materials Science & Engineering, National Chiao Tung University, Hsinchu, Taiwan, Republic of China; ²Department of Materials Science & Engineering, National Cheng Kung University, Tainan, Taiwan, Republic of China and ³Department of Materials Science and Engineering, University of California at Los Angeles, Los Angeles, CA, USA
Correspondence: Professor C Chen, Department of Materials Science & Engineering, National Chiao Tung University, Hsinchu, Taiwan 30010, Republic of China.
E-mail: chih@mail.nctu.edu.tw

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anisotropic growth to grow an array of numerous (100) single-crystal microbumps of Cu on a Si surface.

Exceptionally large grain growth in polycrystalline Cu is known as abnormal grain growth, in which a bimodal distribution of the grain size is obtained.⁵ The larger grains are the abnormal grains. We note that in our extremely anisotropic growth, we do not have a bimodal distribution; instead, we have a normal grain size distribution of all the {100} grains. When we patterned the $\langle 100 \rangle$ -oriented single crystals of Cu on a Si wafer surface, the grain size exhibited a mono-size distribution. The size was the diameter of the Cu microbump, which varied from 25 to 100 μm . This mode of crystal growth is unique, in which numerous single crystals of the same size are grown at the same time. For the typical mode of abnormal grain growth, it would not be possible to obtain the {100} single crystals from those small grains in the bimodal distribution.

MATERIALS AND METHODS

Highly oriented $\langle 111 \rangle$ Cu grains with densely packed nanotwins were fabricated by pulsed electroplating. A high-purity CuSO_4 solution was adopted as the electrolyte, and a high-purity 99.99% copper sheet was employed as the cathode. Proper surfactants and 40 p.p.m. HCl were added to the electrolyte. A Si wafer was adopted as the substrate, and a TiW layer with a thickness of 200 nm was sputtered as the adhesion layer, followed by sputtering of a 200-nm-thick $\langle 111 \rangle$ Cu as a seed layer using an Oerlikon ClusterLine 300 machine (OC Oerlikon Corporation AG, Pfäffikon, Switzerland). The Si wafer was cut into pieces of $3 \times 1 \text{ cm}^2$ or $2 \times 1 \text{ cm}^2$, and these pieces were immersed in the electrolyte during pulsed electroplating. The stirring rate to grow the oriented nanotwinned Cu (nt-Cu) was 600 r.p.m., and the current density was 50 mA cm^{-2} . The duty cycles were $T_{\text{on}} = 0.02 \text{ s}$ and $T_{\text{off}} = 1.5 \text{ s}$. The deposition rate was 1.2 nm s^{-1} under this electroplating condition. The film thickness was controlled to be $\sim 8 \mu\text{m}$ in this study. For the preparation of the equivalent Cu layer without the nanotwins, we used the same electroplating solution at 200 mA cm^{-2} .

To investigate the grain growth of Cu, the Cu films were subjected to annealing at temperatures ranging from 400 to 500 $^\circ\text{C}$ for extended periods. A

quartz-tube furnace with a vacuum at 5×10^{-7} torr was used for the annealing. After annealing, a focused ion beam (FIB) was employed to examine the grain structure of the Cu films. X-ray diffraction was used to analyze the preferred texture of the Cu films. Electron backscattered diffraction (EBSD) was performed to examine the individual grain orientation in the Cu films. The EBSD measurements were performed with a JEOL 7001 F field-emission scanning electron microscope (JEOL Ltd., Tokyo, Japan) with an EDAX/TSL system operated at 25 kV. The analysis was performed with a step size of 50 nm for the Cu grains. OIM software (TSL, Inc., Draper, UT, USA) was used to analyze the crystallographic orientation maps and preferred orientations based on Kikuchi patterns.

RESULTS AND DISCUSSION

We begin with the results of electroplating $\langle 111 \rangle$ -oriented, nt-Cu on Si wafers. The spacing of the nanotwins ranged from 10 nm to 100 nm.¹⁸ Highly oriented $\langle 111 \rangle$ nt-Cu was fabricated on the seed layer. Figures 1a and b show the plan-view of the inverse pole figure map from EBSD and the preferred orientation for the inverse pole figure on the surface grains fabricated at 50 mA cm^{-2} , respectively. The surface grains are all $\langle 111 \rangle$ oriented, and the average grain size for this condition is $2.38 \pm 0.85 \mu\text{m}$. X-ray diffraction also reveals that the Cu film possesses an extremely high $\langle 111 \rangle$ texture, as illustrated in Figure 1c. Figure 1d plots the statistical data for the misalignment angles of the Cu (111) grains in Figure 1a. The misorientation angle of a grain is defined as the angle difference between the pole of the grain and the normal direction to the Si substrate. The average misorientation was obtained by averaging the misorientation angles for all the grains in the EBSD image. For the as-fabricated $\langle 111 \rangle$ -oriented grains in Figure 1a, the average misorientation was 6.4° . An FIB top view of the nt-Cu microstructure is provided in Supplementary Figure S1a in Supplementary Information. All of the {111} twin planes are parallel to the wafer surface; thus, their normal, the $\langle 111 \rangle$ axis, is perpendicular to the wafer surface. EBSD indicates that nearly all of the grains are $\langle 111 \rangle$ oriented; we observe a very high degree of

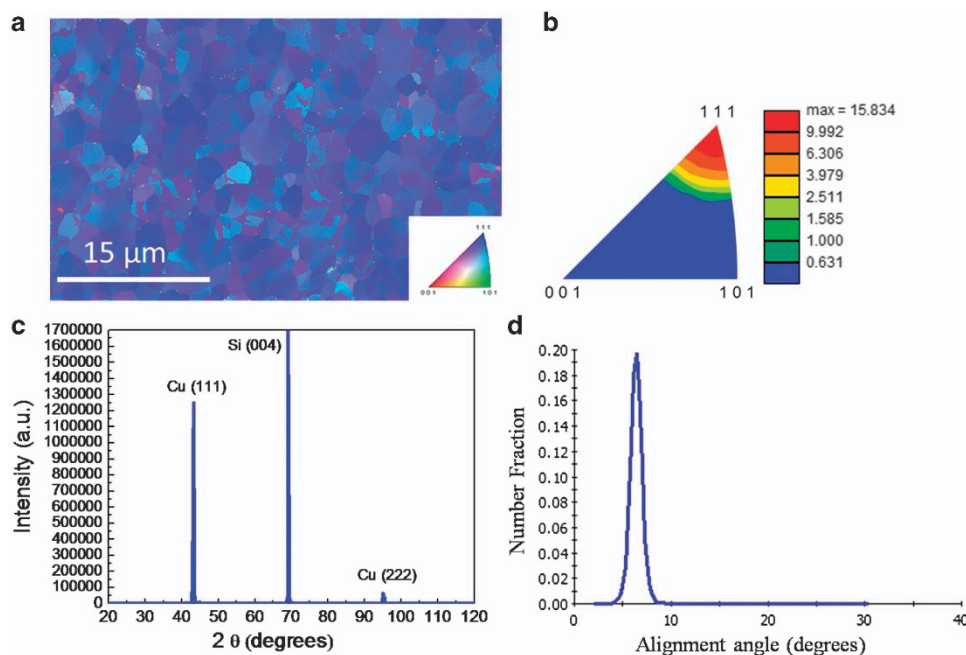


Figure 1 Preferred orientation of the as-deposited Cu film. (a) Plan-view of the inverse pole figure map from for the $\langle 111 \rangle$ -oriented surface grains with color coding of the inverse pole figure. (b) The inverse pole figure of the fabricated Cu film in terms of the out-of-plane direction in a. (c) X-ray diffraction for the as-deposited Cu film. (d) The statistical results showing the numerical fraction of grains as a function of the misalignment angle from the [111] direction.

$\langle 111 \rangle$ texture, over 90%, compared with the value for most published studies on textured Cu films of $\sim 20\%$.¹⁹ By annealing the oriented nt-Cu in the temperature range of 400–500 °C for a few minutes, we obtained some exceptionally large grains that had diameters of ~ 200 – $400 \mu\text{m}$. A top view of one of these grains is presented in Figure 2a; this grain is ~ 65 times larger than the surrounding grains and appears as an abnormal grain. The inverse pole figure map for the grains reveals that the large grain is $\langle 100 \rangle$ oriented. The corresponding FIB image for the grains is presented in Supplementary Figure S1b. Figure 2b shows an enlargement of the white rectangle area in Figure 2a. Figure 2c illustrates the evolution of the X-ray diffraction spectra of a sample annealed at 450 °C for various times, and we observed only two peaks of $\{111\}$ and $\{200\}$ diffraction; the intensity of the latter increases, whereas the former decreases with annealing time.

To convert the abnormal grain growth to solid-state crystal growth, we patterned an array of nt-Cu bumps with diameters ranging from 100 to 25 μm and thicknesses of 10 μm using a lithographic technique. Figure 3a presents the EBSD inverse pole figure map for a 100- μm -diameter circular bump of $\langle 111 \rangle$ -oriented nt-Cu in the as-plated state. The grains are all $\langle 111 \rangle$ oriented. Supplementary Figure S2a presents a corresponding plan-view FIB image of the bump. After annealing at 450 °C for 60 min, Figure 3b presents an image of nine bumps; six of these bumps were completely transformed into a $\langle 100 \rangle$ -oriented single-crystal of Cu, and the other three were very near complete transformation. The corresponding plan-view FIB

image is presented in Supplementary Figure S2b. To verify whether the entire bump was transformed into a single crystal, one of the $\langle 100 \rangle$ bumps was etched by FIB. Figure 3c presents the tilted view of the bump, and the entire bump is a single crystal. We were also able to produce an array of $\langle 100 \rangle$ Cu single crystals with diameters of 75, 50 and 25 μm , as shown in Figure 4. This figure demonstrates the precise control of growth of an array of single-crystal Cu microbumps. The blurry edges in Figure 4f may be attributed to the electron-charging effect in the scanning electron microscope. The 75- μm Cu pads in Figure 4d occupy $\sim 55\%$ of the surface area. However, the 25- μm Cu pads in Figure 4f occupy only $\sim 8\%$ of the surface area. Therefore, the electron-charging effect may occur in the sample with 25- μm Cu pads.

Furthermore, when molten Pb-free solder reacts with a $\langle 100 \rangle$ -oriented single-crystal Cu, an oriented growth of the intermetallic compound of Cu–Sn was obtained.²⁰ In the microbumps for 3D IC interconnects, the amount of solder is very small. Thus, the properties of the Cu–Sn intermetallic compounds, not the solder, will dominate the performance of the microbumps. When a few microns of solder were used, all of the solder transformed to Cu–Sn intermetallic compounds after reflow.^{21,22} Therefore, the $\langle 100 \rangle$ -oriented single-crystal Cu pads can be used to control the electrical and mechanical properties of the microbumps because the growth of oriented intermetallic compound can be controlled.

To determine the mechanism of anisotropic grain growth, we examined cross-sectional views in Figure 5 of samples before

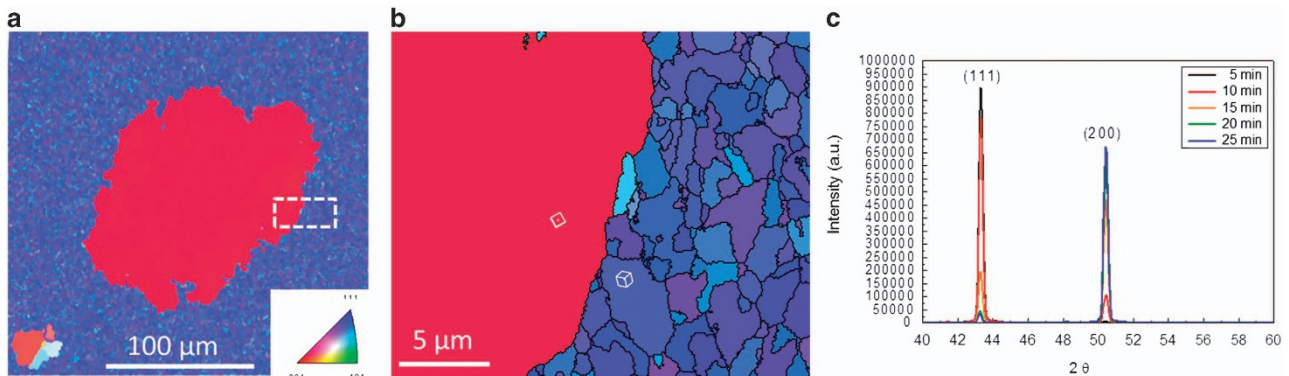


Figure 2 Two-dimensional anisotropic crystal growth in $\langle 111 \rangle$ nt-Cu films. (a) Plan-view inverse pole figure map of the large $\langle 100 \rangle$ grain obtained from electron backscattered diffraction. (b) Enlarged inverse pole figure map for the dashed rectangle in a. (c) Evolution of the X-ray diffraction intensities for the samples annealed at 450 °C for various times.

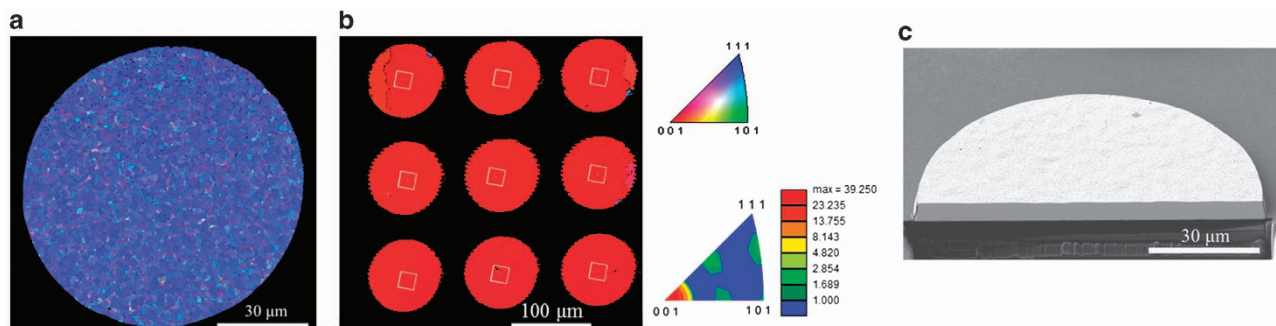


Figure 3 Fabrication of $\langle 100 \rangle$ Cu single crystals after annealing at 450 °C for 60 min. (a) Plan-view electron backscattered diffraction inverse pole figure map for the as-fabricated $\langle 111 \rangle$ Cu crystals. (b) The inverse pole figure map labeled with the crystal lattice for nine $\langle 100 \rangle$ single crystals. (c) Focused ion beam tilt image showing the etched cross-section for a $\langle 100 \rangle$ single crystal.

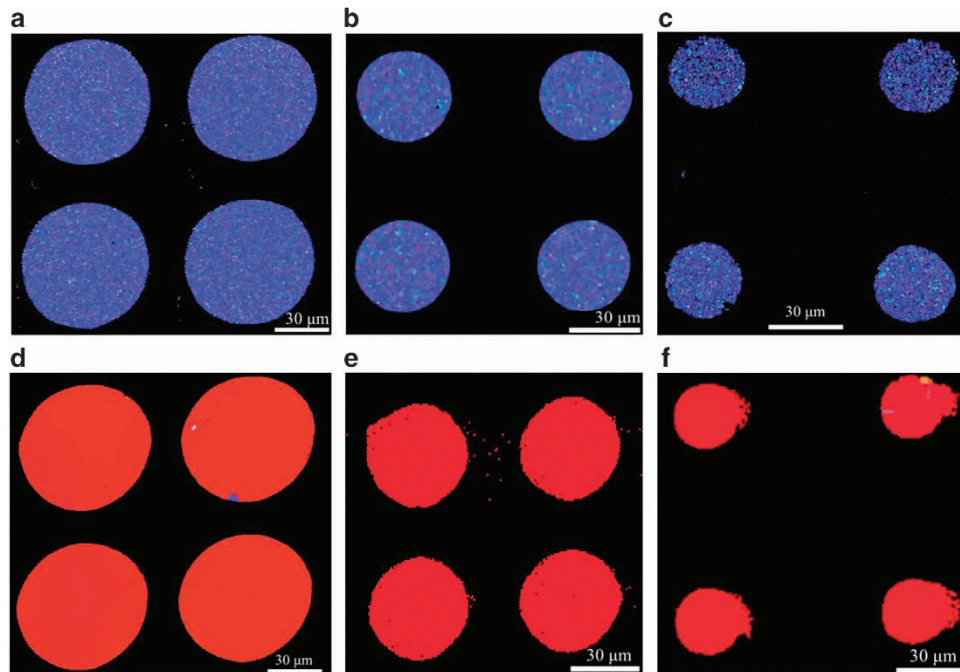


Figure 4 Fabrication of arrays of $\langle 100 \rangle$ Cu single crystals. Plan-view electron backscattered diffraction inverse pole figure map for as-fabricated $\langle 111 \rangle$ Cu crystals: (a) 75- μm diameter, (b) 50- μm diameter and (c) 25- μm diameter. After annealing at 450 $^{\circ}\text{C}$ for 60 min, the $\langle 111 \rangle$ Cu crystals transformed into $\langle 100 \rangle$ single crystals: (d) 75- μm diameter, (e) 50- μm diameter and (f) 25- μm diameter.

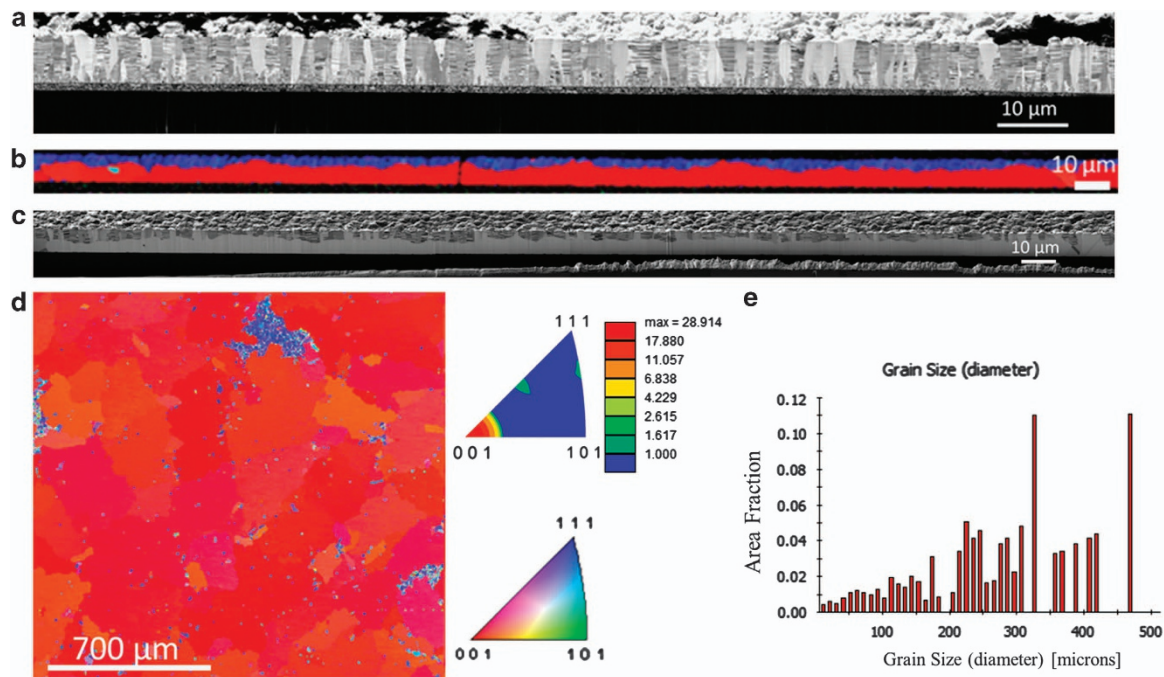


Figure 5 Extremely anisotropic growth of the $\langle 100 \rangle$ Cu crystals. (a) Cross-sectional focused ion beam (FIB) image of the as-fabricated sample. (b) Cross-sectional inverse pole figure map of the large $\langle 100 \rangle$ crystal with a lateral crystal size of 290 μm . (c) Cross-sectional FIB image of the sample in b. (d) Plan-view of the inverse pole figure map from electron backscattered diffraction for the sample annealed at 450 $^{\circ}\text{C}$ for 25 min. (e) Grain size distribution for the grains in d.

(Figure 5a) and after (Figures 5b and c) annealing at 400 $^{\circ}\text{C}$ for 20 min. The $\langle 100 \rangle$ grain on the left-hand side of Figures 5b and c grew laterally to 290 μm . However, its height was only $\sim 5 \mu\text{m}$ in the vertical direction. The anisotropic grain growth has a very slow growth

in the vertical direction, which is at least 50 times slower than the lateral growth. In the $\langle 100 \rangle$ -oriented grains, no nanotwins are detected. Furthermore, the anisotropic grain growth started from the bottom interface rather than from the free surface.

In Figure 5, when we carefully examine the interface at the bottom between the substrate and the columnar nanotwinned grains, we observe a very thin seeding layer of Cu with a high density of $\langle 111 \rangle$ -oriented tiny grains of Cu; however, some randomly oriented

tiny grains of Cu are also observed, including a few of the $\langle 100 \rangle$ -oriented grains, which serve as nuclei of the growth of the anisotropic large grains. The very thin seeding layer was deposited to enhance the growth of the $\langle 111 \rangle$ -oriented nt-Cu layer during electroplating. The

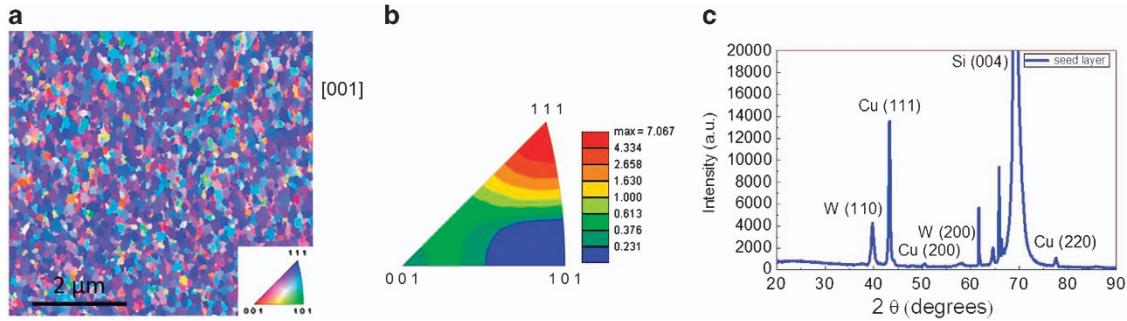


Figure 6 Texture of the 200-nm-thick Cu seed on the TiW adhesion layer. (a) Plan-view of the inverse pole figure map for the $\langle 111 \rangle$ - and $\langle 100 \rangle$ -oriented seed layers with color coding of the inverse pole figure. (b) The inverse pole figure of the Cu seed layer in terms of the out-of-plane direction in a. (c) X-ray diffraction for the Cu seed layer.

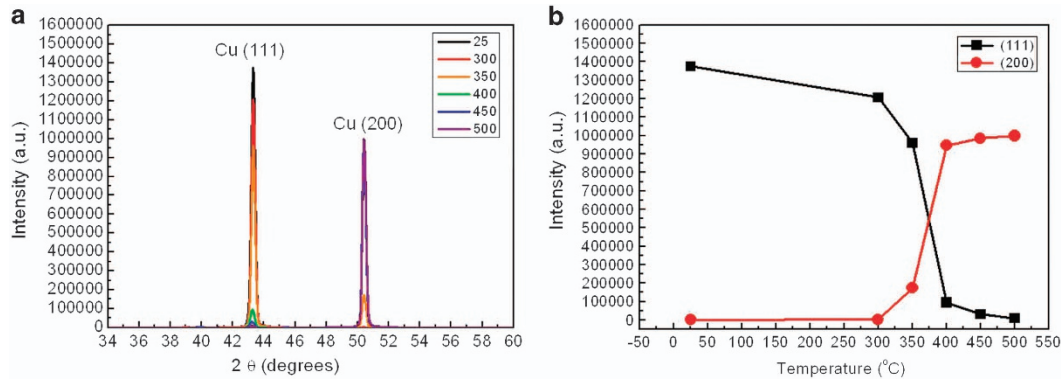


Figure 7 Orientation evolution in the nt-Cu films after annealing at various temperatures. (a) X-ray diffraction of nt-Cu films annealed at temperatures ranging from 300 to 500 °C for 30 min. (b) A plot of the $\{111\}$ and $\{200\}$ intensities versus annealing temperature. The $\langle 100 \rangle$ preferred orientation increases as the annealing temperature increases.

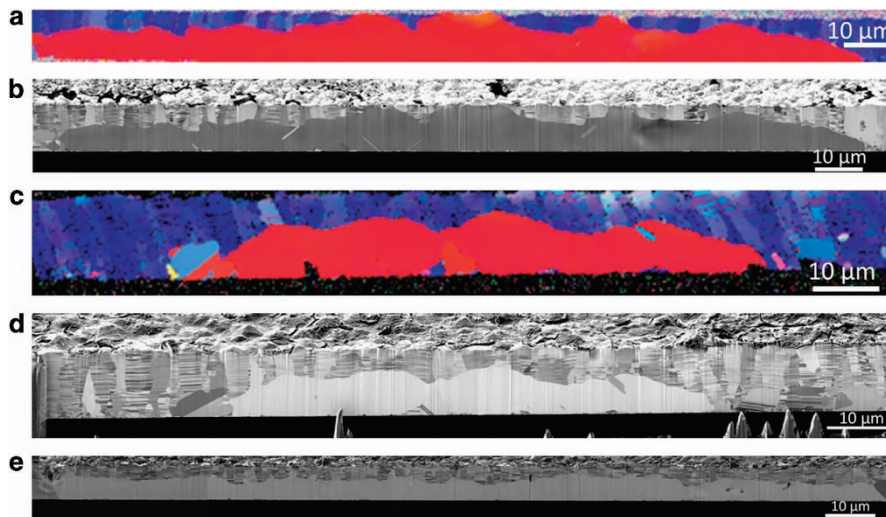


Figure 8 Extremely anisotropic growth of the $\langle 100 \rangle$ Cu grains. (a) Cross-sectional map of the inverse pole figure from electron backscattered diffraction (EBSD) for the sample annealed at 350 °C for 30 min. (b) The corresponding cross-sectional focused ion beam (FIB) image for the sample in a. (c) Cross-sectional map of the inverse pole figure from EBSD for the sample annealed at 400 °C for 10 min. (d) The corresponding cross-sectional FIB image for the sample in c. (e) Cross-sectional FIB images for the sample after annealing at 400 °C for 20 min.

microstructure of the seeding layer is shown in Figure 6. The Cu seed layer has a crucial role in the extremely anisotropic single-crystal growth. Figures 6a and b present the plan-view of the inverse pole figure map from EBSD and the inverse pole figure of the out-of-plane direction, respectively, for the Cu seed layer sputtered on the TiW adhesion layer. The Cu seed film is highly $\langle 111 \rangle$ oriented. However, it is noteworthy that some of the grains were $\langle 100 \rangle$ oriented, as indicated by the inverse pole figure in Figure 6b. Figure 6c shows the X-ray reflections for the Cu seed layer. The diffraction also indicates that the seed layer possesses a strong $\langle 111 \rangle$ texture. However, the intensity of the Cu $\{200\}$ and $\{220\}$ reflections was still detectible. These pre-existing $\langle 100 \rangle$ seed grains will act as nucleation sites for extremely anisotropic grain growth of Cu.

To confirm that the anisotropic single-crystal growth started from the substrate interface, we performed a systematic investigation of the transformation of the $\langle 111 \rangle$ -oriented nt-Cu into the $\langle 100 \rangle$ anisotropic grains between 300 and 500 °C for 30 min. Figure 7 shows the evolution of the X-ray diffraction spectra for the Cu films annealed at 300, 350, 400, 450 and 500 °C for 30 min. Figure 7a presents the

X-ray diffraction spectra of the nt-Cu films annealed at 300, 350, 400, 450 and 500 °C for 30 min. Figure 7b plots the intensity of the Cu $\{111\}$ and $\{200\}$ reflections in Figure 7a. The intensity of the $\{111\}$ reflection decreases as the temperature increases. However, the intensity of the $\{200\}$ reflection increases as the temperature increases. Transformation of $\langle 111 \rangle$ into $\langle 100 \rangle$ grains occurs in the above temperature ranges.

Anisotropic growth of the $\langle 100 \rangle$ grains can be easily observed in the samples annealed at 350 and 400 °C. At temperatures higher than 450 °C, grain growth becomes fast, and the transformation of $\langle 111 \rangle$ grains into $\langle 100 \rangle$ grains is completed in a few minutes. Figures 8a and b present cross-sectional EBSD and FIB images, respectively, for the anisotropic growth of a $\langle 100 \rangle$ grain annealed at 350 °C for 30 min. This $\langle 100 \rangle$ grain has a width of $\sim 160 \mu\text{m}$. Figures 8c and d show a $\langle 100 \rangle$ grain annealed at 400 °C for 10 min. The width of this grain is $\sim 75 \mu\text{m}$. It is noteworthy that the actual grain width may be much larger than it appears in the cross-sectional view because the cross-section is not typically located in the center of the grain. The most significant anisotropic growth occurs at 400 °C for 20 min in this

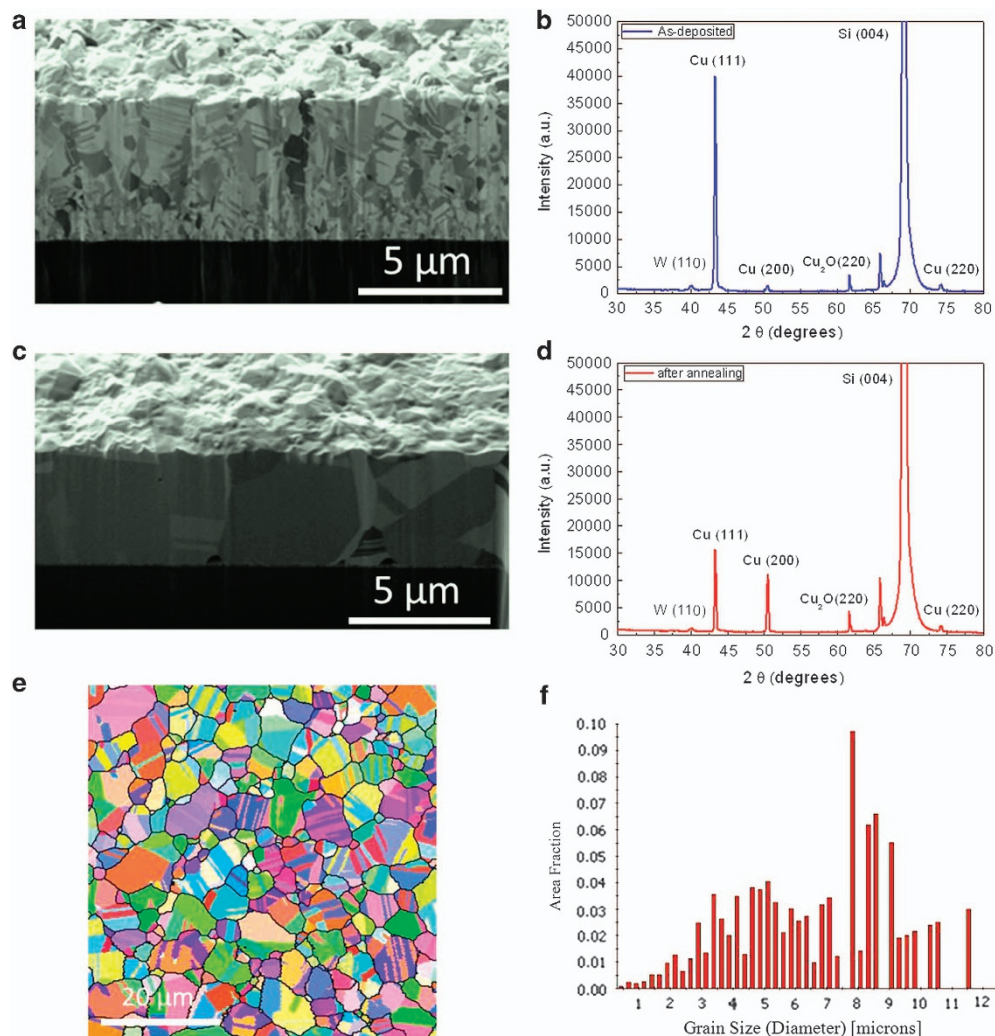


Figure 9 Normal grain growth in Cu films without $\langle 111 \rangle$ -oriented nt-Cu. (a) Cross-sectional focused ion beam (FIB) image for the as-fabricated Cu film at 200 mA cm^{-2} . (b) The corresponding X-ray diffraction spectrum for the Cu film in a. (c) Cross-sectional FIB image for the Cu film after annealing at 450 °C for 30 min and (d) the corresponding X-ray diffraction spectrum for the Cu film in c. (e) Plan-view of the inverse pole figure map from electron backscattered diffraction for the Cu grains annealed at 450 °C for 30 min. (f) Grain size distribution for the grains in e. The average grain size is only $6.6 \mu\text{m}$, which is approximately the thickness of the Cu film.

study. Figure 8e and Supplementary Figure S4 show $\langle 100 \rangle$ grains grown to 170 and 207 μm , respectively, after annealing at 400 $^{\circ}\text{C}$ for 20 min.

To determine whether the $\langle 111 \rangle$ -oriented nanotwins are essential in achieving the extremely anisotropic growth of $\langle 100 \rangle$ crystals, we prepared randomly oriented nanotwins; however, no extremely abnormal grain growth was observed. Figure 9a presents a FIB cross-sectional view of the randomly oriented nanotwins in Cu, and Figure 9b presents the corresponding X-ray diffraction spectrum, where the $\{111\}$, $\{200\}$ and $\{220\}$ reflections of Cu are observed. After annealing at 450 $^{\circ}\text{C}$ for 30 min, normal grain growth occurred. Figures 9c and d present the FIB cross-sectional image and the corresponding X-ray diffraction spectrum, respectively. In addition, an inverse pole figure map of the top view of the grains obtained from EBSD in Figure 9e reveals no extremely anisotropic grains. Figure 9f presents a histogram of the grain size distribution that reveals that the average grain size is 6.6 μm .

Statistical analysis on the grain size distribution was performed for the $\langle 100 \rangle$ Cu grains transformed from $\langle 111 \rangle$ oriented nt-Cu on Si substrates. Figure 10a presents the plan-view of the inverse pole figure map from EBSD for the $\langle 100 \rangle$ Cu grains after the annealing at 400 $^{\circ}\text{C}$ for 1 h. The average misorientation angle was 3.0 $^{\circ}$ for the grains in Figure 10a, as illustrated in Supplementary Figure S5a. There were 738 grains analyzed for this condition. Figure 10b illustrates the fitting of grain size distribution for the crystals annealed at 400 $^{\circ}\text{C}$ for 1 h, which reveals a normal grain size distribution and not a bimodal distribution. Figure 10c presents the plan-view of the inverse pole

figure map from EBSD for the $\langle 100 \rangle$ Cu crystals annealed at 450 $^{\circ}\text{C}$ for 1 h. The average misorientation angle was 2.8 $^{\circ}$ for the grains in Figure 10c, as illustrated in Supplementary Figure S5b. We analyzed 743 Cu grains in this condition. The fitting of the grain size distribution for the crystals annealed at 450 $^{\circ}\text{C}$ for 1 h is shown in Figure 10d. The $\langle 100 \rangle$ Cu grains in this condition also exhibit a normal grain size distribution.

In contrast, in Supplementary Figures S3a and b, we present two connecting images of the cross-section of a sample annealed at 450 $^{\circ}\text{C}$ for 15 min, where the entire sample has been completely transformed into $\langle 100 \rangle$ -oriented grains. In Figure 5d, a plan-view of the top surface is shown with respect to the inverse pole figure map. The color of all of the grains is nearly identical, which indicates that all of the grains are $\langle 100 \rangle$ oriented and that the grain size distribution of the $\langle 100 \rangle$ -oriented grains is very narrow. The average misorientation angle was 4.7 $^{\circ}$ for the grains in Figure 5d, as illustrated in Supplementary Figure S5c. The average grain size is $\sim 283 \mu\text{m}$, and the largest grain was $\sim 470 \mu\text{m}$, as revealed by the histogram of the grain size provided in Figure 5e. The microstructure of a sample before annealing is presented in Figure 1 and Supplementary Figure S1a, where the average grain size is only $\sim 2.4 \mu\text{m}$.

We define the necessary and sufficient conditions to obtain extremely anisotropic crystal growth below. The first condition is a high density of $\langle 111 \rangle$ -oriented nanotwins, which provide the stored energy to serve as the driving force in the growth of the anisotropic and untwinned $\{100\}$ single crystals. Randomly oriented nanotwins will not do serve as this driving force. The second condition is the

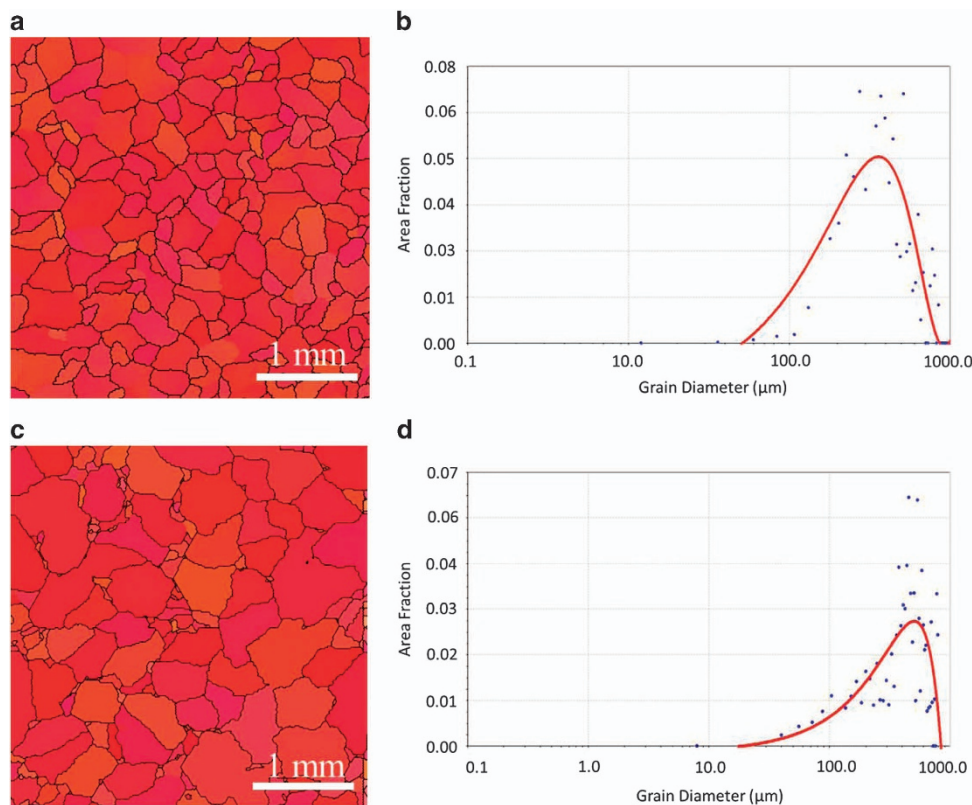


Figure 10 Statistical analysis on grain size distribution for the $\langle 100 \rangle$ Cu grains transformed from $\langle 111 \rangle$ -oriented nt-Cu on Si substrates. (a) Plan-view of the inverse pole figure map from electron backscattered diffraction (EBSD) for the $\langle 100 \rangle$ Cu grains after annealing at 400 $^{\circ}\text{C}$ for 1 h. There were 738 grains analyzed in this condition. (b) Fitting of grain size distribution for the crystals annealed at 400 $^{\circ}\text{C}$ for 1 h. (c) Plan-view of the inverse pole figure map from EBSD for the $\langle 100 \rangle$ Cu crystals annealed at 450 $^{\circ}\text{C}$ for 1 h. We analyzed 743 Cu grains in this condition. (d) Fitting of grain size distribution for the crystals annealed at 450 $^{\circ}\text{C}$ for 1 h. The $\langle 100 \rangle$ Cu grains in both conditions exhibit a mono-sized distribution.

presence of a seeding layer with a low density of $\langle 100 \rangle$ -oriented seeds. When we prepared a seeding layer with a very high density of $\langle 111 \rangle$ -oriented seeds, we obtained a very high intensity of $\langle 111 \rangle$ -oriented nt-Cu; however, upon annealing, we could hardly detect the growth of $\langle 100 \rangle$ grains because there were no $\langle 100 \rangle$ -oriented seeds. The third condition is that the substrate should have a large difference in thermal expansion coefficient from that of the metal film, such that a biaxial strain in the in-plane directions of the film occurs under high-temperature annealing.

It is well known that abnormal grain growth in Cu has been explained by strain relaxation. Strain will clearly have a role in this study of extremely abnormal grain growth. In addition to the biaxial thermal strain mentioned above, it has been reported that in pulse-electroplating nt-Cu, *in situ* strain measurement using the bending beam method reveals that the nt-Cu film is under in-plane biaxial tensile strain.²³ We will not attempt to explain the mechanism of abnormal grain growth here because what we have is the microstructural evolution of an elastically anisotropic inclusion (a two-dimensional grain) within an elastically anisotropic matrix (the oriented nt-Cu).

In summary, we report an extremely abnormal grain growth of $\langle 100 \rangle$ -oriented single crystals of Cu in a matrix of $\langle 111 \rangle$ -oriented and nt-Cu columnar grains; the lateral growth is approximately two orders of magnitude faster than the vertical growth. We have applied this anisotropic growth to produce an array of numerous $\langle 100 \rangle$ -oriented single crystals of Cu microbumps with sizes from 100 to 25 μm on Si wafer surfaces.

CONFLICT OF INTEREST

The authors declare no conflict of interest.

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Supplementary Information accompanies the paper on the NPG Asia Materials website (<http://www.nature.com/am>)