

is the first demonstration of line-specific memory in individual invertebrates. In a study based on only two bacterial strains, maternally transferred immunity seemed to be strain-specific⁹. Invertebrates are occasionally capable of a specific response in the context of allogenic recognition¹⁰, but specificity in these cases is based on compatibility factors that have evolved within the species and so is not directly comparable with defence against a parasite.

The mechanism responsible for the line-specific memory in this copepod remains to be investigated, although binding by lectins is one possibility¹¹. If specific memory in invertebrate defence is more widespread than originally thought, it would have implications for our understanding of innate immunity, host–parasite coevolution, and invertebrate-vector diseases such as malaria.

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Competing financial interests: declared none.

Microfluidic systems

High radial acceleration in microvortices

Microfluidic systems can conveniently be used for rapid analysis of biological samples^{1–5}. Here we describe a single re-circulating flow, or microvortex, that can generate a maximum fluid rotational velocity of up to 12 m s⁻¹ and a corresponding radial acceleration in excess of 10⁶g. Such microvortices may be exploited in centrifugal microdevices to investigate the effects of high radial acceleration on biological and chemical processes.

A microvortex was generated in a diamond-shaped microchamber (Fig. 1a), which was visualized by using fluorescent polystyrene beads (diameter, 2 μm; Fig. 1b, left). The average linear velocity of the pressure-driven flow in the 30-μm-wide channel was 3 m s⁻¹. The diamond-shaped

notch along the side wall of an otherwise smooth channel causes flow detachment at the opening of the notch.

Unlike macroscale lid-driven cavity flow⁶, the combination of extremely tight rotation radii ($r < 50 \mu\text{m}$) and high rotational velocity ($v > 10 \text{ m s}^{-1}$) of fluid flow from constricted microchannels can lead to generation of radial accelerations (v^2/r) as high as 10^7 m s^{-2} . As the average linear velocity in the 30-μm-wide channel was increased from 3 to 45 m s⁻¹ (Reynolds number (Re) = 245 at the microchamber opening), the polystyrene beads (density, $\rho = 1.05 \text{ g cm}^{-3}$) in water were spun out of the microchamber owing to high centrifugal acceleration (Fig. 1b, right).

To illustrate the potential applications of this microvortex system, we flowed red polystyrene beads ($\rho = 1.05 \text{ g cm}^{-3}$) and green silica beads ($\rho, 1.8\text{--}2.0 \text{ g cm}^{-3}$) in a solution of 50 % (w/w) of CsCl ($\rho = 1.56 \text{ g cm}^{-3}$;

Fig. 1c). As the average flow velocity in the 30-μm-wide channel was increased from 1.5 m s⁻¹ (Re = 13 at the cavity's opening) to 20 m s⁻¹ (Re = 122 at the cavity's opening), the red fluorescent beads were concentrated towards the centre of the vortex, whereas the green beads were spun towards the chamber's outer edge. Although the linear flow velocity was 20 m s⁻¹, the corresponding volumetric flow rate was only about 18 μl s⁻¹, owing to the small cross-sectional area of the 30-μm-wide channel.

A simple calculation of the radial acceleration from the average flow rates in the microchannel overestimates the acceleration that we observed experimentally, mainly because of the damping of the flow velocity near the walls of the microchannels for pressure-driven flows. Because the flow that drives the vortex at the opening of the diamond-shaped cavity is near the 30-μm-wide channel wall, the maximum fluid rotational velocity inside the cavity in Fig. 1b, c is only about 20% of the average flow velocity in the channel.

To attain ultrahigh rotational velocities in the microvortex, we designed a trapezoidal microchamber connected to a constricted microchannel that was placed asymmetrically along the width of a 200-μm × 200-μm channel (see supplementary information). Using this design, we measured maximum rotational velocities as high as 12 m s⁻¹ at a distance of 10 μm from the vortex core, which corresponds to a radial acceleration of $1.4 \times 10^7 \text{ m s}^{-2}$ or $1.4 \times 10^6 \text{ g}$ (Fig. 1d). We measured rotation rates (Fig. 1d), both by the time taken for a photolysed (uncaged) plug of fluorescein dye to travel a calibrated distance (5 μm) from a focused nitrogen photolysis laser beam to an argon laser probe volume (Fig. 1d, squares), and by the time taken for the plug of dye to travel completely around the microvortex (Fig. 1d, open circles)⁷ (and see supplementary information).

The ability to use fluid flow with no moving parts to generate large centrifugal forces should prove useful for integrated micro-analytical separations — for example, in sample preparation.

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Supplementary information accompanies this communication on Nature's website.

Competing financial interests: declared none.

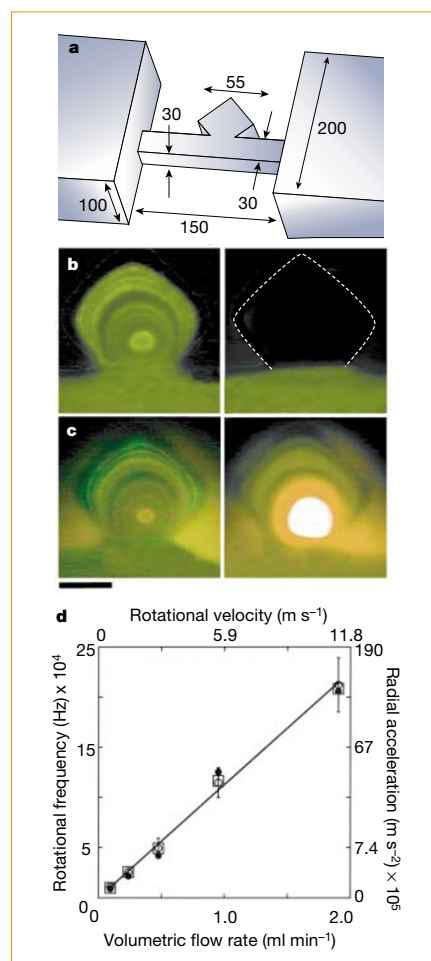


Figure 1 Generation and application of microvortices. **a**, Microchannel used to create the microvortex shown in **b, c**; dimensions are in μm. **b**, Fluorescent beads are spun out of the chamber as the flow velocity increases (from left to right). **c**, Separation of low-density red beads from high-density green beads in the microvortex as the flow velocity increases (from left to right). Scale bar, 25 μm. **d**, The experimentally measured rotational rates (squares and open circles) show excellent agreement with the vorticity results (filled circles) extracted from our simulations (some symbols overlap).