

Fig. 2. Gamma spectrum of iodine-131 in rainfall. O, Source; x, background

to iodine-131 in view of the comparatively short half-lives of the other iodine fission products. A confirmatory measurement on one of the samples was made by γ -ray spectrometer analysis for iodine-131. Good agreement was obtained, as can be seen in Table 2. The spectrum obtained is shown in Fig. 2.

The confirmatory test of the γ -ray spectrometer on the second sample yielded a result of 82 $\mu\mu\text{c./l.}$ of iodine-131 compared with 95.5 $\mu\mu\text{c./l.}$ by β -analysis.

Although these preliminary results indicate clearly the onset of the Russian testing programme, it is impossible at this stage to attempt a more detailed correlation with individual tests.

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Nuclear Spin of Protoactinium-231 from Hyperfine Splitting

THE nuclear spin of protoactinium was first determined by Schuler and Gollnow in 1934¹. They examined the spectrum of protoactinium-231 in the region 4300–5600 Å. and obtained the nuclear spin $I = \frac{3}{2}$ from hyperfine splitting.

More recently the nuclear spin of protoactinium-233 has been determined² from paramagnetic resonance, and the nuclear spin of tetravalent protoactinium-231 from an investigation of paramagnetic resonance and hyperfine structure³. Both authors obtain $I = \frac{3}{2}$.

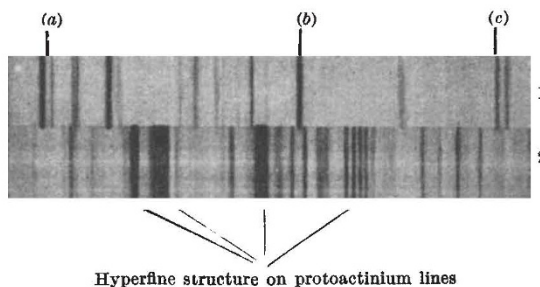


Fig. 1. Protoactinium-231 spectrum with thorium comparison spectrum. 1, Thorium comparison spectrum; 2, protoactinium spectrum; a, thorium, 5869.847 Å.; b, thorium, 5874.981 Å.; c, thorium, 5878.922 Å.

So far, however, the work of Schuler and Gollnow has not been confirmed.

The atomic emission spectrum of protoactinium-231 was recently photographed in the region 4000–9000 Å. using the 21-ft. grating spectrograph at Harwell, and also recorded in the region 0.8–2.5 μ . The light source was an electrodeless discharge tube containing about 3 mgm. of protoactinium iodide and excited by microwaves. The spectrum shows pronounced hyperfine structure, the majority of lines being split into four components. This confirms the value $I = \frac{3}{2}$ obtained by previous workers.

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¹ Schuler and Gollnow, *Naturwiss.*, **22**, 511 (1934).

² Marrus, Nierenberg and Winocur, *Nuclear Phys.*, **23**, 1, 90 (1961).

³ Axe, Stapleton and Jeffries, *Nuclear Phys. Abstr.*, **15**, 3213 (1961).

Rotational Analysis of ⁶⁵Cu³⁵Cl Bands

ALTHOUGH many extensive band systems are known for the CuCl molecule¹, our knowledge of their rotational structure has been meagre. Apart from being heavy, natural cuprous chloride is a mixture of four isotopic species with not very different relative abundances. Very high dispersion and resolution are therefore needed to resolve, even partially, the rotational structure of the bands.

With the availability of compounds containing separated copper and chlorine isotopes, it has become possible to prepare copper chloride with any desired copper and chlorine isotopes. Milligram quantities of ⁶⁵Cu³⁵Cl₂ have been thus prepared and the spectrum of ⁶⁵Cu³⁵Cl excited in an electrodeless discharge tube by a 2,450 mc./sec. Raytheon microwave oscillator². All the known band systems of cuprous chloride were obtained with high intensity. The bands were photographed mostly in the fifth order of the Argonne National Laboratory 30-ft. concave grating spectrograph (dispersion 0.3 Å./mm.). Bands for which rotational analysis has been done are given in Table 1. Spectroscopic data obtained from the present work are given in Table 2. All values in this table, except those of r_e , are in cm.⁻¹ units.

Table 1

System	Bands analysed			
B - X(¹ Π - ¹ Σ)	(1,0);	(0,0);	(0,1)	
C - X(¹ Σ - ¹ Σ)	(1,0);	(0,0);	(0,1)	
D - X(¹ Π - ¹ Σ)	(3,1);	(2,0);	(1,0);	(0,0); (0,1); (0,2)
E - X(¹ Σ - ¹ Σ)	(4,1);	(3,0);	(2,0);	(1,0); (0,0); (0,1)
F - X(¹ Π - ¹ Σ)	(3,0);	(2,0);	(1,0);	(0,0); (0,1); (0,2)