LETTERS TO THE EDITORS

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NOTES ON POINTS IN SOME OF THIS WEEK'S LETTERS APPEAR ON P. 600.

CORRESPONDENTS ARE INVITED TO ATTACH SIMILAR SUMMARIES TO THEIR COMMUNICATIONS.

Nature of the Nebular Red-Shift

FROM an investigation (to be published in *Physica*) of the proper vibrations of expanding spherical space, it follows that—in extremely good approximation—light is propagated with respect to co-moving co-ordinates irrespective of the expansion, except that (a) the time-rate of events is slowed down and (b) all energy portions decrease, both inversely proportional to the radius of curvature.

The slowing down secures the constancy of the velocity of light and entails the nebular red-shift, which from this point of view takes place during the passage. The attempt to decide by observation, whether it is actually due to expansion, rests on two important formulæ, which follow from the new view with great ease. Let l be the linear diameter of a nebula at the moment of emission and χ its angular distance from the observer (linear distance divided by the circumference of space), then the angle $d\theta$ between two geodesics of space, pointing at the moment of emission from the observer to the ends of the diameter, is from pure geometry:

$$d\theta = \frac{l}{R \sin \chi},\tag{1}$$

R being the radius of curvature at the moment of emission. By the theorem quoted above, d0 is also the observed angular diameter of the nebula (Hubble and Tolman, equation 3).

Again, let the energy emitted by the nebula within an appropriately chosen unit of time be E_0 . It will soon assume the shape of a spherical shell of thickness c (say). Let $R_{\rm obs}$ be the radius of space, when this shell reaches the observer. Its surface at this moment is, by pure geometry, $4\pi R^2_{\rm obs}$, $\sin^2 \gamma$. By the theorem quoted above, its thickness then is $c R_{\rm obs}/R$ and its energy is $E_0 R/R_{\rm obs}$. Hence its energy density ρ is

$$\rho = \frac{E_0}{4 \pi c R^4_{\text{obs.}}} \cdot \frac{R^2}{\sin^2 \chi}.$$
 (2)

 ρ is a measure of the bolometric luminosity, observed outside the earth's atmosphere (Hubble and Tolman, equation 4).

My purpose in re-stating here these two important formulæ due to Tolman is to make the following remarks. Both l and E_0 refer to the moment of emission, which is different for two nebulæ observed simultaneously. Should l and E_0 exhibit a general dependence on R, then it would no longer be reasonable to regard them as constants, when equations (1) and (2) are combined (as they actually are) with the hypothesis of uniform spatial distribution of the nebulæ. For the latter, if admitted at all, has to apply to nebulæ which are intrinsically similar at the same moment of time—not at such moments as depend on the accidental position of our galaxy.

As regards l, the question is, whether we are inclined to assume (a) that the distances between the stars within a nebula behave, on the average, like the distances between two points of a rigid body—say, the ends of the Paris metre rod; or (b) like the distance between two distant nebulæ. Clearly the case of the stars is intermediate. To regard l as a constant means to decide for the first alternative. The second one would make l/R constant, giving formula (1) the same form as in the case of a non-recessional explanation of the red-shift (see Hubble and Tolman, equation 3').

and Tolman, equation 3'). As regards E_0 , the possible general decline of the nebular candle-powers has already been mentioned by Hubble and Tolman (see their concluding remarks). To the assumption that the same amount of energy is emitted during every second, there is a peculiarly simple alternative, namely, that the amounts of energy, which have been emitted during a second, remain equal. On account of the decay of travelling energy, this assumption would mean $E_0 \sim 1/R$, which reduces equation 2 to the same form as in the case of a non-recessional explanation of the red-shift (see Hubble and Tolman, equation 4'). I do not mean to suggest $E \sim 1/R$ particularly. I mention it in the way of an example.

These remarks detract nothing from the importance of deciding by observation how d0 and ρ actually behave, if the photographs are interpreted as assuming uniform spatial distribution. I understand that present evidence points to observed luminosities (ρ) decreasing with distance not even quite as rapidly as we should expect (with $E_0 = \text{const.}$) from the non-recessional explanation. If that is so, I should say they rather support the recessional explanation, in spite of its predicting a still more rapid decrease of the ρ 's. The discrepancy, though greater, can here be removed by assuming the E_0 's to decrease with time; an assumption which is very plausible in an expanding universe, which, on the whole, cools down; but not at all plausible in a static one.

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7 Sentier des Lapins, La Panne, Belgium. July 31.

¹ Hubble, E., and Tolman, R. C., Astrophys. J., 82, 302 (1935).

The Forbidden ${}^3P_0{}^{-1}D_2$ Line of O III in the Nebular Spectrum of Nova Herculis 1934

ALTHOUGH the two well-known lines of [O III] $\lambda = 5007 \, \text{A.} \, (^3P_2 - ^1D_2)$ and $\lambda = 4959 \, \text{A.} \, (^3P_1 - ^1D_2)$ are the most prominent features in the spectra of planetary nebulæ and novæ at the nebular stage, the third line of the triplet, corresponding to the $^3P_0 - ^1D_2$