the existence of γ -rays shows that the observed curves must be superpositions of at least two simple spectra, whereas their shapes are not such as could be represented as sums of two K.U. curves with endpoints differing by the energy of the γ -rays.

Physics Department, Cornell University.

F. HOYLE. R. PEIERLS.

H. A. BETHE.

Emmanuel College, Cambridge. University, Birmingham. Dec. 24.

¹ Fermi, E., Z. Phys., 88, 161 (1934).

- ² Kurie, F. N. D., Richardson, J. R., and Paxton, H. C., Phys. Rev., 49, 368 (1936).
- ³ Konopinski, E. J., and Uhlenbeck, G. E., Phys. Rev., 48, 7 (1935).
- ⁴ Richardson, H. O. W., Proc. Roy. Soc., A, 161, 456 (1937).
 ⁵ Cockeroft, J. D., Proc. Roy. Soc., A, 161, 540 (1937).
- ⁶ Walke, H. (in the Press).

C

C

- Walke, H. (II. the Press).
 Fierz, M., Helv. phys. Acta, 10, 123 (1937).
 Hoyle, F., Proc. Roy. Soc., A, 166, 249 (1938).
 Fowler, W. A., Delsasso, L. A., and Lauritsen, C. C., Phys. Rev., 49, 569 (1936). ¹⁰ Richardson, J. R., *Phys. Rev.*, **53**, 610 (1938).
 ¹¹ Burcham, W. E., and Smith, C. L., *Proc. Roy. Soc.*, A, **168**, 176
- (1938). 12 Private communication.
- ¹³ Cockcroft, J. D., and Lewis, W. B., Proc. Roy. Soc., A, 154, 261 (1936).

A New Form of the Baryteron Equation and Some **Related** Questions

It is known that the 2-component spinors introduced by Van der Waerden must be completed to 4-component quantities if linear transformation of their components by the complete Lorentz group including reflections is required. These quantities, often called 'Dirac wave-functions', will be called here 'four-spinors' or 'undors* of the first rank'. Defining the reflection in the right way¹, it is possible to join to every four-spinor ψ another four-spinor² ψ^{ξ} , the components of which are linear combinations of those of ψ^* . This ψ^{\sharp} is called the *charge-conjugated*³ of ψ . For if ψ is a solution of the Dirac equation for a positive particle,

$$mc\psi = \beta\{\varepsilon - (\alpha, \pi)\}\psi,$$
$$\varepsilon = i\hbar(\delta/\delta t) - e\varphi, \quad \overrightarrow{\pi} = -i\hbar\overrightarrow{\nabla} - (e/c)\overrightarrow{A}$$

then $\psi^{\boldsymbol{\epsilon}}$ is a solution of the Dirac equation for a negative particle :

$$mc\psi^{\pounds} = \beta \{ \varepsilon^L - (\alpha.\pi^L) \} \psi^{\pounds},$$

 $\varepsilon^L = i\hbar(\delta/\delta t) + e\varphi, \ \vec{\pi^L} = -i\hbar \vec{\nabla} + (e/c)\vec{A};$

vice versa ($\psi^{\text{e}\text{e}} = \psi$). Four-spinors, which are equal to their own charge-conjugated, we may call neutrinors. These quantities, which are equivalent to a two-spinor together with its complex conjugated, are not only useful in the so-called neutrino theory of light, but also Majorana's theory of neutrons⁴ may be summarized by stating that the quantized fourspinor of a neutral particle is a neutrinor.

An undor-calculus analogous to tensor-calculus and Van der Waerden's spinor-calculus can be developed. An under of the second rank $\Psi_{k,k'}$ (k,k' = 1, 2, 3, 4), that is, a quantity transforming like the product of two four-spinors, may be regarded then as consisting of a scalar, a pseudo-scalar and a pseudo-4-vector (the antisymmetrical part of $\Psi_{k,k'}$) and a regular * Unda = Wave.

functions of protons and neutrons and by α , β and

 α' , β' the Dirac matrices operating on k and k', the baryteron equation^{5,7,8} is given in the following simple form :

$$2m_{0}c\{\Psi_{k,k'} + 2if_{0}p(\psi_{kk}^{\tilde{\Sigma}}\psi_{Pk'} + \psi_{kk'}^{\tilde{\Sigma}}\psi_{Pk})\} =$$

$$\vec{\beta}\{\varepsilon - (\vec{\alpha}.\vec{\pi})\}\Psi_{k,k'} + \beta'\{\varepsilon - (\vec{\alpha'.\pi})\}\Psi_{k,k'}.$$
(1)

The operator f_{op} , which operates on the arguments k and k', becomes a constant c-number (f) only if the two constants f and g, occurring in Kemmer's equations⁵, would happen to be equal. As experiment seems to indicate⁹ that these constants do not differ very much, this formal argument for the advantage of equality of f and g seems to be interesting

The electric charge-density and electric current are given by⁶

$$\rho = (e/8\hbar) \cdot \Psi^*(\beta + \beta')\Psi,$$

$$\vec{j} = (e/8\hbar) \cdot \Psi^*(\vec{\beta\alpha'} + \vec{\beta'\alpha})\Psi.$$

The wave-function of free baryterons can be interpreted to some extent as the Dirac wave-function in the hole-theory of electrons. In both cases the same kinds of difficulty arise.

As for the complications occurring in baryteron theory^{7,8}, as well as in the quantum theory of the Maxwellian field^{10,11}, which are caused by the fact that some of the canonical equations of the field do not contain time derivatives, it is possible to develop a theory of quantization of mixed fields of Fermi-Dirac particles and Einstein-Bose quanta, in which the question of dealing with such complications is answered in a rather general way.

Details will be published elsewhere.

F. J. BELINFANTE.

Instituut voor Theoretische Natuurkunde,

Leiden. Dec. 22.

- * Baovreoog = heavier.
- ¹ Racah, Nuov. Cim., 14, 322 (1937).
- ² Pauli, Ann. Inst. H. Poincaré, 6, 130 (1936). ³ Kramers, Proc. Roy. Acad. Amsterdam, 40, 814 (1937).
- 4 Majorana, Nuov. Cim., 14, 171 (1937).

- ⁵ Kemmer, Proc. Roy. Soc., A, 166, 127 (1938).
 ⁶ Compare De Broglie, "Actual. Scient. Indust.", No. 181 (1934) and No. 411 (1936).
- Yukawa, Proc. Phys. Math. Soc. Japan, 20, 319 and 720 (1938).
 Bhabha, Proc. Roy. Soc., A, 166, 501 (1938).
- ⁹ Fröhlich, Heitler and Kemmer, Proc. Roy. Soc., A, 166, 154 (1938).
- ¹⁰ Heisenberg and Pauli, Z. Phys., 56, 1 (1929).
- 11 Rosenfeld, Ann. Inst. H. Poincaré, 2, 25 (1932).

Exploration of the Sun's Disk and its Astrophysical Results

IN NATURE recently (Dec. 10, p. 1035), I took the opportunity, in connexion with the recent researches on practical relativity by H. E. Ives and G. R. Stilwell, to place on record that on principles scarcely anywhere denied now, a radiating atom, or system of atoms, falling free into the sun or describing a free orbit close around it should not experience any internal influence on the periods of its radiation, the