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## Coherent Charge Transport in Ballistic InSb Nanowire Josephson Junctions

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Hybrid InSb nanowire-superconductor devices are promising for investigating Majorana modes and topological quantum computation in solid-state devices. An experimental realisation of ballistic, phase-coherent superconductor-nanowire hybrid devices is a necessary step towards engineering topological superconducting electronics. Here, we report on a low-temperature transport study of Josephson junction devices fabricated from InSb nanowires grown by molecular-beam epitaxy and provide a clear evidence for phase-coherent, ballistic charge transport through the nanowires in the junctions. We demonstrate that our devices show gate-tunable proximity-induced supercurrent and clear signatures of multiple Andreev reflections in the differential conductance, indicating phase-coherent transport within the junctions. We also observe periodic modulations of the critical current that can be associated with the Fabry-Pérot interference in the nanowires in the ballistic transport regime. Our work shows that the InSb nanowires grown by molecular-beam epitaxy are of excellent material quality and hybrid superconducting devices made from these nanowires are highly desirable for investigation of the novel physics in topological states of matter and for applications in topological quantum electronics.

Study of coherent charge transport through superconductor-normal conductor-superconductor (SNS) Josephson junction devices is of fundamental importance for understanding the properties of superconducting quantum devices<sup>1,2</sup>. When a normal conductor is connected to two closely spaced superconductor leads, Cooper pairs can leak from the superconductor leads into the normal conductor, due to the superconducting proximity effect, and induce a supercurrent through the junction. The advances in nanofabrication technology and material science have made it possible to experimentally realise a variety of Josephson junctions with reduced dimensions. These mesoscopic superconducting hybrid structures realisable with various types of materials have offered great possibilities to investigate novel quantum phenomena and new device concepts including, e.g., tunable supercurrent transistors<sup>3-7</sup>, Andreev bound states<sup>8,9</sup>, Cooper pairs splitters<sup>10,11</sup>, superconducting electronic refrigerators<sup>12</sup>, and topological superconducting states of matter<sup>13-15</sup>. Recently, Josephson junctions based on InSb nanowires have been attracting growing attention owing to their extraordinary transport properties, such as long coherence length, tunable one-dimensional subbands, good contacts with superconductor metals, and strong spin-orbit coupling<sup>16-19</sup>. In particular, signatures of Majorana zero-energy states have been observed in hybrid InSb nanowire-superconductor devices<sup>18-20</sup>. Beside their interest in fundamental physics, Majorana bound states in solid state systems have potential application in topological quantum computation<sup>21-23</sup>. Moreover, the physics of the Josephson junctions made from InSb nanowires is rich due to the interplay between strong spin-orbit coupling and Zeeman splitting in the presence of a magnetic field. For example, various novel physical phenomena, such as the fractional Josephson effect, anomalous Josephson current and bound subgap states<sup>24-27</sup>, have been predicted to be observable in such hybrid quantum devices. Therefore, it is of broad interest and fundamental importance to study the proximity-induced superconducting properties of InSb nanowire-based hybrid structures.

So far, most of the experimental works on InSb nanowire-superconductor junctions have been carried out in the diffusive transport regime<sup>16,18,19</sup>. The combined effects of phase-coherent transport and proximity-induced

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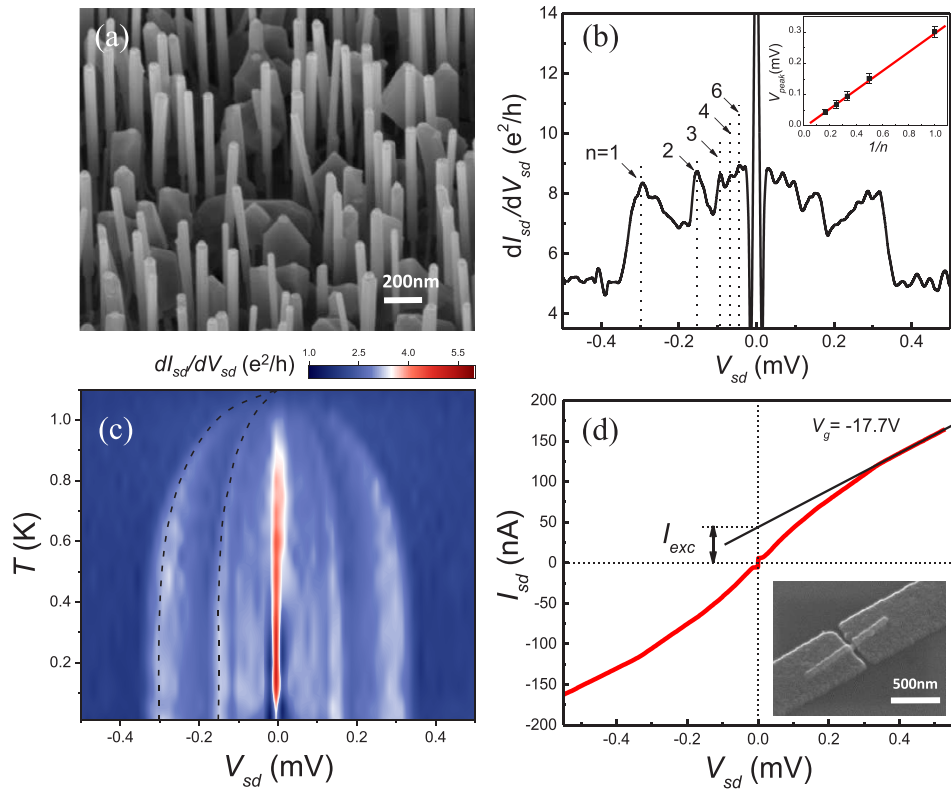
superconductivity in the ballistic transport regime have yet to be systematically explored in InSb nanowire-based Josephson junctions. There are two major challenges in the realisation of a ballistic InSb nanowire-based Josephson junction. One is to achieve highly transparent interfaces between superconductor contacts and an InSb nanowire, and the other is to obtain an InSb nanowire with perfect crystal quality. While the first one is essential to enhance the proximity-induced superconductivity in the nanowire junction, the second one is critical for achieving a long mean free path, desired for ballistic transport, in the nanowire. Recently, we have succeeded in growing high crystal quality InSb nanowires by molecular-beam epitaxy (MBE)<sup>28,29</sup>. The cleanliness of the MBE growth environment enables growth of semiconductor nanowires with high purity and perfect crystal quality. A transport study of MBE-grown InSb nanowires has demonstrated the formation of long single quantum dots with lengths up to 700 nm<sup>30</sup>. Considering the fact that an alternative theoretical scenario, taking into account the effect of disorder, has been put forward to explain the previously observed zero-energy states<sup>31</sup>, these state-of-the-art InSb nanowires are desirable for exclusion of such explanations in Majorana fermion detection experiments and for the study of novel ballistic, quantum coherent transport phenomena.

In this work, we report on realisation and systematic low-temperature transport measurements of superconductor-InSb nanowire-superconductor Josephson junction devices with highly transparent contacts to superconductor leads, in which the InSb nanowires are grown by MBE. This is the first report on a systematic study of the transport properties of MBE-grown InSb nanowire Josephson junction devices. We demonstrate that the devices show proximity-induced supercurrent and clear signatures of multiple Andreev reflections, indicating that the charge transport is phase-coherent within the junctions. We also observe periodic modulations of the critical current and show that the modulations are associated with the resonant states formed by Fabry-Pérot interference in the nanowires. Our results demonstrate that transport in MBE-grown InSb nanowire Josephson junctions is ballistic and phase-coherent, and that such nanowires are appropriate for investigation of novel mesoscopic superconducting phenomena in the presence of strong spin-orbit interaction and for detection and manipulation of Majorana bound states in the solid state.

## Results

**Josephson junctions based on individual InSb nanowires.** Figure 1a shows a scanning electron microscope (SEM) image of an array of as-grown vertically standing InAs/InSb heterostructure nanowires on an InP(111)B substrate. Single nanowire Josephson junction devices with aluminum electrodes are fabricated from the as-grown InSb nanowires on a highly doped silicon substrate (used as a global back gate) covered with a 200-nm-thick SiO<sub>2</sub> capping layer. The inset of Fig. 1d displays a representative single nanowire Josephson junction device. The diameters of the InSb nanowires used in this work range from 50 to 120 nm and the lengths are in the range of 1–2 μm. The source-drain separations  $L$  of the devices are in the range of 60–200 nm. Figure 1b shows the differential conductance  $dI_{sd}/dV_{sd}$  measured for a device with ~56 nm in nanowire diameter and ~60 nm in contact separation (device D1) as a function of source-drain bias voltage  $V_{sd}$  at base temperature  $T = 10$  mK, far below the superconducting critical temperature of aluminum ( $T_c \approx 1.1$  K). The conductance shows a pronounced peak at  $V_{sd} = 0$  V, indicating the occurrence of a supercurrent flowing through the nanowire. The conductance rises when the bias voltage is decreased below  $V_{sd} = 2\Delta/e \approx 300$  μeV, where  $\Delta \approx 150$  μeV and  $2\Delta$  is the superconducting energy gap of aluminum. It is worth noting that the conductance at bias voltages lower than the superconducting energy gap is nearly doubled compared to the normal-state conductance. This is evidence that contact interfaces between the superconductor electrodes and the nanowire are highly transparent and is fully consistent with the Andreev reflection mechanism of transport through a Josephson junction with two almost ideal transparent contact interfaces as described by the Blonder-Tinkham-Klapwijk (BTK) theory<sup>32</sup>. At finite bias voltages, subharmonic gap structures, symmetrically situated around  $V_{sd} = 0$  V, can be clearly observed in the measured conductance. Such subharmonic gap structures are commonly observed in an SNS junction<sup>4,5,16</sup> and could be understood by invoking multiple Andreev reflection (MAR) processes. In a MAR process, an electron is reflected multiple times, gaining an energy of  $eV_{sd}$  in each reflection, until its energy is above the superconducting energy gap. This process generates a series of conductance peaks within the superconducting energy gap, with different peaks corresponding to different orders of successive Andreev reflections<sup>33,34</sup>. As shown in the inset of Fig. 1b, the voltage positions  $V_{peak}$  of the differential conductance peaks fall on top of the expected linear dependence  $V_{peak} = 2\Delta/ne$ , where  $n$  is an integer number. Each data point in the inset is obtained from an average of several  $dI_{sd}/dV_{sd}$  traces at different gate voltages and the error bar attached to the data point denotes the standard deviation. The evolutions of these subharmonic conductance peaks with temperature are shown in Fig. 1c. The positions  $V_{peak}$  of the differential conductance peaks show systematic shifts towards lower bias voltages with increasing temperature. At temperatures above the critical temperature  $T_c$  of aluminum, these subharmonic gap structures disappear. Also, as expected, the positions of the subharmonic gap structures follow the BCS-type temperature dependence, in which the superconducting energy gap varies with temperature approximately as  $\Delta(T) \approx \Delta(0) \sqrt{\cos\left[\frac{\pi}{2}\left(\frac{T}{T_c}\right)^2\right]}$  (see the dashed lines in Fig. 1c for the evolutions of the peak positions for  $V_1 = 2\Delta/e$  and  $V_2 = \Delta/e$ ). The observation of MARs implies the coherent transfer of multiple charges in our junction device.

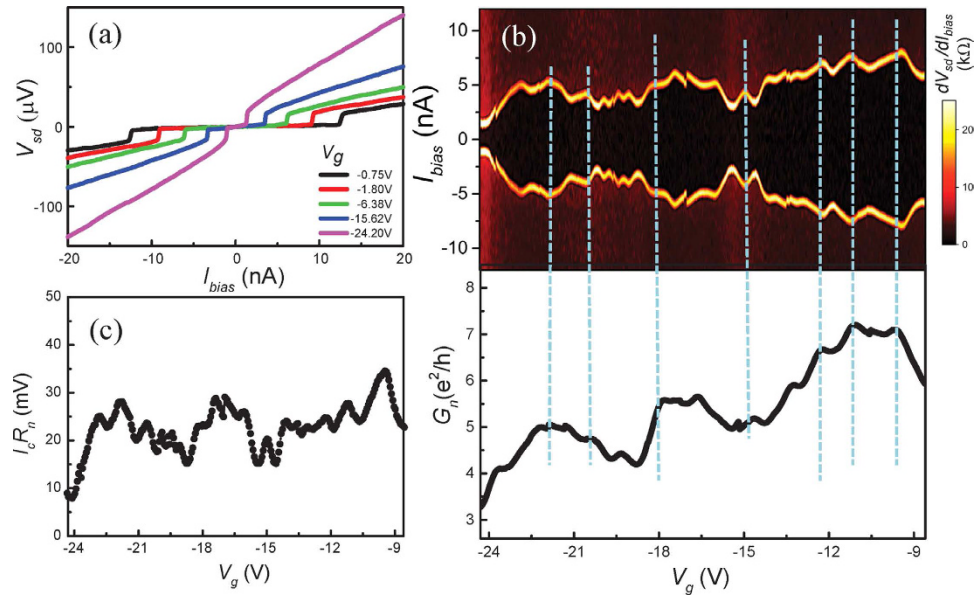
The transparency of the interfaces between the InSb nanowire and the aluminum contacts can be estimated from the excess current  $I_{exc}$ —a non-zero current offset at  $V_{sd} = 0$  V as obtained by the linear extrapolation of the current-voltage characteristics at high bias voltages ( $V_{sd} > 2\Delta/e$ )—due to the transfer of Cooper pairs through the junction via Andreev reflection processes. In the BTK theory for Andreev reflection, the strength of the potential barrier at a superconductor-normal conductor interface can be characterized by a dimensionless scattering parameter  $Z$ , which relates to the transmission coefficient at the interface as  $T_r = 1/(1 + Z^2)$ <sup>35,36</sup>. The scattering parameter  $Z$  can be estimated from the ratio  $eI_{exc}R_n/\Delta$ , where the normal state resistance  $R_n$  and  $I_{exc}$  can be determined from the linear fit to the measured current-voltage curve at high bias voltages and its intercept with the



**Figure 1. Multiple Andreev reflections in an InSb nanowire Josephson junction device.** (a) SEM image of as-grown InSb nanowires grown via an InP/InAs stem technique on an InP(111)B substrate. These nanowires are 50–90 nm in diameter. Gold seed particles are visible on top of individual nanowires. (b) Differential conductance  $dI_{sd}/dV_{sd}$  versus bias voltage  $V_{sd}$  measured for device D1 made from an InSb nanowire with a diameter of  $\sim 56$  nm at gate voltage  $V_g = -15.25$  V and base temperature  $T = 10$  mK. The arrows and vertical dotted lines indicate conductance peaks at the voltage positions of  $eV_n = 2\Delta/n$  ( $n = 1, 2, \dots$ ), corresponding to multiple Andreev reflection processes. The inset shows the peak position versus the inverse of the peak index  $1/n$ . The data points are averaged for several  $dI_{sd}/dV_{sd}$  traces of different gate voltages and the error bars denote the standard deviation. The solid line is a linear fit to the peak positions and is seen to go through the origin as expected. (c)  $dI_{sd}/dV_{sd}$  as a function of  $V_{sd}$  and temperature  $T$  at zero magnetic field. The dashed lines follow the temperature evolutions of the conductance peaks due to multiple Andreev reflection processes, as expected from the temperature dependence of  $\Delta(T)$  predicted by the BCS theory. The high zero-bias conductance peak originates from the proximity induced superconductivity in the nanowire. (d) Current-voltage characteristics measured at  $V_g = -17.7$  V and  $T = 10$  mK. The linear fit of the current-voltage curve at high bias voltages ( $V_{sd} > 2\Delta/e$ ) extrapolates to a finite value of the current at zero bias voltage, i.e., excess current  $I_{exc}$ . The inset shows an SEM image of the corresponding InSb nanowire-based Josephson junction. Here, the separation between the two contacts is  $\sim 60$  nm.

current axis at zero bias voltage, respectively. The current-voltage characteristics of the junction measured at gate voltage  $V_g = -17.7$  V and base temperature  $T = 10$  mK are shown in Fig. 1d. Using  $\Delta = 150 \mu\text{eV}$ , we obtain  $Z \sim 0.42$ , which corresponds to  $T_r = 0.85$  at the interfaces, indicating that the nanowire-aluminium contacts are highly transparent with the transmission close to the ideal value of unity. To the best of our knowledge, this high value of the transparency is the best result ever achieved for InSb nanowire devices contacted by superconductor electrodes. We attribute this high transparency at the nanowire-superconductor interfaces to both the high material quality of the MBE-grown InSb nanowire and the improvement in the contact fabrication technique (see Methods).

In Fig. 2a we show current-biased, four-probe, measurements of the voltage-current characteristics of the device at different gate voltages  $V_g$  and base temperature  $T = 10$  mK. A characteristic dc Josephson response is observed—each  $V_{sd} - I_{bias}$  curve exhibits a zero-voltage region when the current  $I_{bias}$  is below a certain value  $I_c$ , providing a clear evidence for the proximity-induced superconductivity in the hybrid superconductor-nanowire device. We can observe an induced supercurrent through a hybrid superconductor aluminium-InSb nanowire device with a junction length up to 200 nm (see Supplementary Fig. S1). Beyond  $I_c$ , switching occurs from the superconducting state to the dissipative resistive state, leading to the abrupt appearance of a finite source-drain voltage. To show the dependence of  $I_c$  on  $V_g$  in more detail, we plot in the top panel of Fig. 2b the differential resistance  $dV_{sd}/dI_{bias}$  as a function of the applied gate voltage  $V_g$  and bias current  $I_{bias}$ . In the central dark region,  $dV_{sd}/dI_{bias} = 0$  and the nanowire is in the superconducting state with the critical current  $I_c$  given by the thin bright

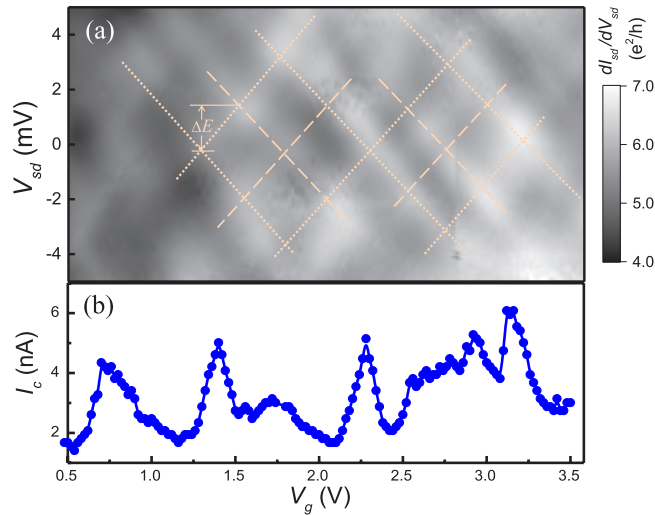


**Figure 2. Gate-modulated supercurrent in an InSb nanowire Josephson junction.** (a) Measured voltage  $V_{sd}$  as a function of bias current  $I_{bias}$  for device D1 at  $T = 10$  mK and different gate voltages  $V_g$ . Here a gate-tunable supercurrent is seen. (b) Upper panel: Differential resistance  $dV_{sd}/dI_{bias}$  as a function of  $I_{bias}$  and  $V_g$ . The central dark area is the region with  $dV_{sd}/dI_{bias} = 0$ , i.e., the device in the superconducting state. The bright boundaries correspond to the transitions from the superconducting state to the dissipative state. The critical current  $I_c$  measured by the bright boundaries, oscillates as a function of  $V_g$ . Lower panel: Normal state conductance  $G_n$  of device D1 extracted from linear fits to the  $I_{sd} - V_{sd}$  curves at the high bias voltage region with  $V_{sd} > 2\Delta/e$ . Here, the correlations between the modulations of the critical current and the oscillations of the normal state conductance are clearly seen as indicated by vertical dashed lines. (c) Product of the critical current  $I_c$  and the normal state resistance  $R_n$  of the device extracted as a function of  $V_g$ .

lines of the high resistance in the figure. Here it is easily seen that the magnitude of the critical current  $I_c$  is gate voltage-dependent. In the areas outside the central dark region, the nanowire is in the dissipative state. The bottom panel of Fig. 2b displays the evolution of the corresponding normal state conductance  $G_n$  over the same  $V_g$  range, where  $G_n$  is determined from linear fits of the measured  $I_{sd} - V_{sd}$  curves at high source-drain voltages ( $V_{sd} > 2\Delta/e$ ). It can be readily seen that the critical current is closely correlated with the normal state conductance  $G_n$ . Such correlations between  $I_c$  and  $G_n$  have been established for SNS junctions both in the diffusive and in the ballistic transport regime<sup>37,38</sup>. Based on field-effect electrical measurements, we are able to determine an electron density of  $n_e \simeq 9 \times 10^{17} \text{ cm}^{-3}$  and a mobility of  $\mu \simeq 2000\text{--}4000 \text{ cm}^2/\text{V} \cdot \text{s}$  in the nanowire. These values give an electron mean free path  $l_e \sim 40\text{--}80$  nm, which is comparable to the contact separation in the device. The superconducting coherence length is given by  $\xi = (\hbar D/\Delta)^{1/2} \simeq 500$  nm, where  $D$  is the diffusive coefficient. Therefore, our junction operates in the short junction ( $L \leq \xi$ ) and ballistic or quasi-ballistic regime ( $L \leq l_e$ ). According to theory, the product  $I_c R_n$  of the critical current and normal state resistance in the clean short-junction limit should be  $\pi\Delta/e \approx 470 \mu\text{V}$ <sup>38</sup>. As shown in Fig. 2c, the measured values,  $I_c R_n \approx 10\text{--}30 \mu\text{V}$ , are however much smaller than expected. Such strong suppression of the  $I_c R_n$  product has been commonly observed in a Josephson junction device and has been attributed to premature switching in such a resistively and capacitively shunted Josephson junction device<sup>3–5,16,39</sup>.

**Fabry-Pérot interference in ballistic InSb nanowire Josephson junctions.** It is seen in Fig. 2b that the modulations of  $I_c$  are strongly correlated with the variations of the normal state conductance of the junction in the open conduction region where the normal state conductance lies in a range of 3 to  $8 e^2/h$ . Such correlated  $I_c$  modulations and  $G_n$  variations have been previously observed in several semiconductor-based Josephson junctions<sup>3,4,40,41</sup> and have been attributed to the gate voltage induced changes in the carrier density. Below, we will show further that the observed modulations of  $I_c$  and variations of  $G_n$  arise from the phase-coherent transport through the Fabry-Pérot-type resonant states formed in the clean nanowire junction region and show that, as observed in carbon nanotube-based ballistic SNS junctions<sup>5,42,43</sup>, the critical current and the normal state conductance exhibit regular periodic oscillations.

Figure 3a displays a gray scale plot of the differential conductance  $dI_{sd}/dV_{sd}$  for device D2, where the nanowire's diameter is  $\sim 60$  nm and the contact separation is  $\sim 95$  nm, as a function of  $V_g$  and  $V_{sd}$ . A small magnetic field of  $B = 50$  mT is applied perpendicularly to the device substrate in order to suppress superconductivity in the electrodes. The differential conductance exhibits a characteristic “checkerboard” pattern (see dashed lines in Fig. 3a for a guided view), showing quasi-periodic oscillations with both  $V_{sd}$  and  $V_g$ . The oscillations can be understood qualitatively by ballistic electron transport through the resonant states formed within the electron cavity confined by the two metal contacts in analogy to an optical Fabry-Pérot interferometer. In a ballistic device,



**Figure 3. Fabry-Pérot interference in a ballistic InSb nanowire Josephson junction.** (a) Normal-state differential conductance  $dI_{sd}/dV_{sd}$  measured for device D2 with  $L \sim 95$  nm in contact separation and  $D \sim 60$  nm in nanowire diameter as a function of source-drain voltage  $V_{sd}$  and gate voltage  $V_g$  at  $T = 10$  mK and a magnetic field of 50 mT. Here, a Fabry-Pérot-like interference pattern is observable, as indicated by dashed and dotted lines in the figure. The smallest energy spacing of the Fabry-Pérot resonances is marked as  $\Delta E$  in the plot. (b) Critical current  $I_c$  of the device measured in the same gate voltage range as in (a).

the phase-coherent electron transport in the nanowire between the contacts can lead to transmission resonances when the Fermi wave vector  $k_F$  of electrons in the nanowire matches the condition for resonance. By tuning either  $V_{sd}$  or  $V_g$ , a shift in the Fermi energy of electrons in the nanowire leads to a change in the Fermi wave vector, giving rise to a shift in the condition for resonance along a crisscrossing mesh line. Though similar patterns of periodic conductance oscillations have previously been observed in short junction devices<sup>44–49</sup>, there has not been a systematic report for observation of such patterns in a device made from an InSb nanowire. With a parabolic electron dispersion in a nanowire, the smallest energy spacing  $\Delta E$  of Fabry-Pérot resonances can be expressed as  $\Delta E = \hbar^2 \pi^2 / 2m^* L_c^2$ , where  $\hbar$  is the Planck constant,  $m^*$  is the effective mass and  $L_c$  is the length of the resonant cavity<sup>47</sup>. For device D2, the observed value of this energy spacing is  $\Delta E \simeq 2\text{--}3$  meV, corresponding to an effective cavity length  $L_c \approx 75\text{--}115$  nm, in good agreement with the geometrical length of the junction, i.e., the contact separation, in the device. The observation of a regularly oscillating conductance interference pattern implies that the charge transport through the nanowire in the junction is in the ballistic regime.

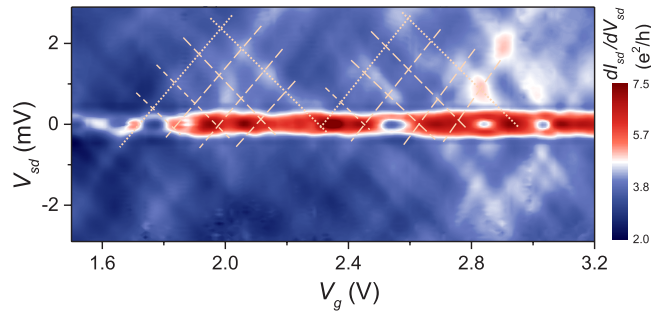
In Fig. 3b we plot the critical current  $I_c$ , extracted from voltage-current characteristics of device D2 measured at zero magnetic field, over the same  $V_g$  range as in Fig. 3a. The critical current  $I_c$  displays clearly quasi-periodic variations with the gate voltage  $V_g$ . In a nanowire electron Fabry-Pérot cavity, quantum interference leads to a set of resonant levels. In the case of the nanowire being in the proximity-induced superconducting state, the Josephson current can flow through these discretely resonant levels, giving rise to periodic modulations of  $I_c$ <sup>38,50,51</sup>. As shown in Fig. 3b, the critical current exhibits peaks at the same gate voltages where the device shows maximal normal state conductance values, providing a compelling evidence that such Fabry-Pérot constructive interference manifests itself in the superconducting states. Considering a Josephson junction in an electromagnetic environment, the critical current through quantized energy levels in the resonant regime is given by<sup>5</sup>

$$I_c \sim \left[ 1 - \left( 1 - \frac{G_n}{4G_0} \right)^{1/2} \right]^{3/2}. \quad (1)$$

This correlation between  $I_c$  and  $G_n$  explains our observation shown in Fig. 3a,b. In addition, we note that its deviation from the linear dependence of  $I_c$  on  $G_n$  predicted for an ideal SNS junction could explain the observed  $V_g$  dependence of  $I_c R_n$ , as shown in Fig. 2c.

## Discussion

As shown in Fig. 3a, we can clearly identify two sets of periodic oscillations in the interference pattern, highlighted by dashed lines and dotted lines, respectively. As we will discuss below, these two sets of periodic oscillations originate from the multimode transport in the InSb nanowire. In our devices, the shape of the nanowires can be viewed as a cylinder where the diameter of the nanowire  $D$  is comparable to the channel length  $L$ . As a result, both the transverse quantization and the longitudinal Fabry-Pérot interference contribute to the formation of the resonant states in the nanowires. The relative magnitudes of the energy quantization resulting from the transverse confinement and the longitudinal Fabry-Pérot interference can be approximately determined based on the aspect ratio  $\xi \sim D/L$ . In device D2 with  $L \sim 95$  nm and  $D \sim 60$  nm ( $\xi > 1$ ), the fast oscillations with a period of  $\sim 1.6$  V in  $V_g$  originate from the longitudinal Fabry-Pérot interference and the slow oscillations correspond to the transverse modes. Such resonant conductance and critical current oscillations are also observed in the other nanowire devices with different aspect ratios. In Fig. 4, we display the differential conductance measured as a function of  $V_g$  and  $V_{sd}$  for device D3 with  $L \sim 65$  nm and  $D \sim 90$  nm ( $\xi < 1$ ) at zero magnetic field. At low bias voltages ( $V_{sd} \leq 2\Delta/e$ ), the proximity-induced supercurrent exhibits quasi-periodic oscillations with  $V_g$ . At



**Figure 4. Charge stability diagram of an InSb nanowire Josephson junction at  $B = 0$ .** Differential conductance  $dI_{sd}/dV_{sd}$  versus  $V_{sd}$  and  $V_g$  measured for device D3 with  $L \sim 65$  nm in contact separation and  $D \sim 90$  nm in nanowire diameter at zero magnetic field and  $T = 10$  mK. At bias voltages below  $2\Delta/e$ , the differential conductance oscillates quasi-periodically, indicating that the supercurrent is modulated, as a result of the formation of resonant states in the nanowire junction by Fabry-Pérot interference. The dashed and dotted lines in the figure are guides to the eyes, highlighting the checkerboard pattern of Fabry-Pérot conductance oscillations at high bias voltages.

high bias voltages ( $V_{sd} > 2\Delta/e$ ), we also observed two sets of Fabry-Pérot oscillations. As expected, the larger diameter of the nanowire results in smaller resonant energy level spacings and thus a shorter period of conductance oscillations. Similar complex interference patterns have been previously observed in the devices made from two-dimensional electron gas and graphene, and have been attributed to the two-dimensional geometry of the cavities<sup>48,49,52–55</sup>. Finally, it is worthy to note that our devices operate in the high conductance regime, where more than one subbands are occupied in the nanowires, which could give rise to a beating behaviour in conductance oscillations<sup>47</sup>.

In summary, we have investigated low temperature transport properties of Josephson junction devices made from high crystal quality, MBE-grown InSb nanowires with highly transparent interfaces between superconductor electrodes and the nanowires. We have observed that the devices show gate-tunable proximity-induced supercurrent and signatures of multiple Andreev reflections in the differential conductance, as expected for phase-coherent transport through the nanowire junctions. We have also observed that the devices exhibit characteristic Fabry-Pérot oscillations in the normal state conductance, a manifestation of ballistic transport through the nanowire junctions, and modulations of the supercurrent associated with the resonant states formed by Fabry-Pérot interference in the nanowires. This work presents a first report on a systematic study of ballistic InSb nanowire based Josephson junction devices and provides an important step towards development of topological quantum computation devices based on hybrid InSb nanowire-superconductor structures.

## Methods

**InSb nanowire growth and device fabrication.** Single crystalline InSb nanowires used in this study were grown by gas source MBE on semi-insulating InP(111)B substrates, using a known stem technique<sup>29,56–58</sup>. The grown InSb segments were confirmed to be 100% perfectly untwined zinc blende single crystal nanowires, as usually observed for gold-assisted growth by other techniques<sup>56,57</sup>. All nanowires showed some lateral overgrowth but no measurable tapering. Details on the growth and characterization of the nanowires can be found elsewhere<sup>28</sup>.

After mechanical transfer from a growth chip to a Si/SiO<sub>2</sub> substrate prepatterned with contact pads and alignment markers on top of the SiO<sub>2</sub> layer, individual InSb nanowires were located using SEM. The superconducting electrodes consisting of 5-nm-thick titanium and 90-nm-thick aluminum were defined on the located individual InSb nanowires by electron-beam lithography (EBL), electron-beam evaporation of metal and lift-off processes. To lower the contact resistance between InSb nanowires and the metallic electrodes, the native oxide layers on the surfaces of the InSb nanowires were removed using an ammonium polysulfide ((NH<sub>4</sub>)<sub>2</sub>S<sub>x</sub>) solution prior to metal deposition<sup>59</sup>. Nanowire devices with resistance at room temperature in a range of 5–30 kΩ were selected for low-temperature transport measurements in this work. Comparing with earlier works<sup>16,19</sup>, we used longer time for etching and passivation treatment in order to achieve highly transparent contact interfaces. Before the devices were cooled down for low-temperature transport measurements, we pumped the chamber for ~24–48 hours in order to eliminate adsorbed molecules on the nanowire surfaces and the substrate<sup>60</sup>.

**Low-temperature measurement.** All measurements were carried out in an Oxford <sup>3</sup>He/<sup>4</sup>He-dilution refrigerator equipped with a superconducting magnet, using a standard four-terminal measurement configuration. The current and voltage preamplifiers were calibrated by standard resistors. In order to minimise the electronic noise, we used a series of  $\pi$ -filters, copper-powder filter and RC filters at different temperature stages for covering different frequency ranges<sup>4,5</sup>. Whenever needed, a magnetic field was applied perpendicularly to the substrate plane and thus to the nanowires in this work.

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## Author Contributions

N.K. and H.Q.X. conceived and designed the experiments. P.C. developed and conducted the growth of the nanowires. S.L. and D.X.F. fabricated the devices. S.L., D.X.F. and N.K. performed the electrical measurements. L.B.W. and Y.Q.H. helped with low-temperature measurements. S.L. and N.K. analyzed the measurement data. S.L., N.K. and H.Q.X. co-wrote the manuscript. H.Q.X. supervised the project. All authors discussed the results and commented on the manuscript.

## Additional Information

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