

# SCIENTIFIC REPORTS

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## Direct frequency comb optical frequency standard based on two-photon transitions of thermal atoms

Received: 25 February 2015  
Accepted: 16 September 2015  
Published: 13 October 2015

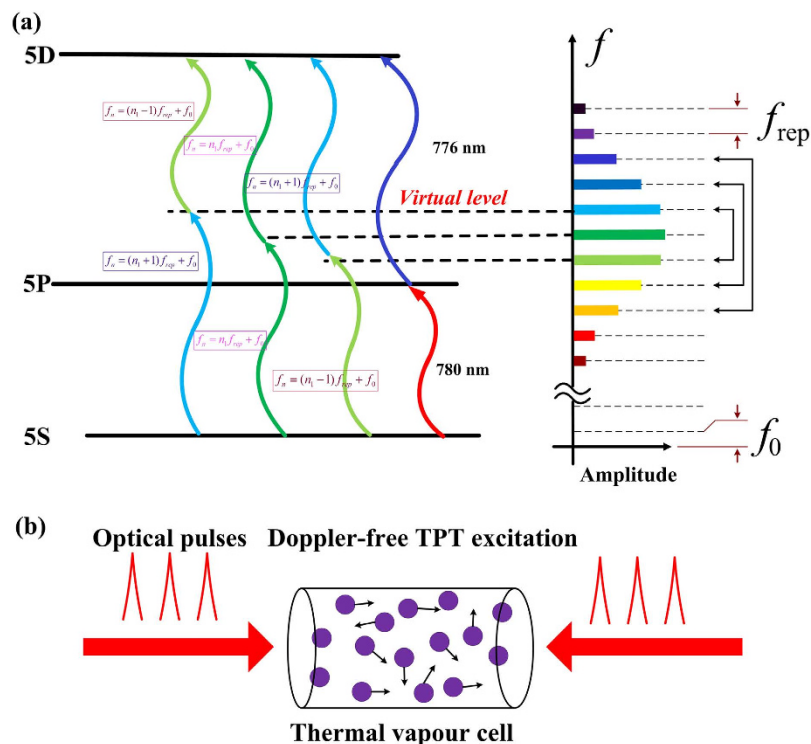
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Optical clocks have been the focus of science and technology research areas due to their capability to provide highest frequency accuracy and stability to date. Their superior frequency performance promises significant advances in the fields of fundamental research as well as practical applications including satellite-based navigation and ranging. In traditional optical clocks, ultrastable optical cavities, laser cooling and particle (atoms or a single ion) trapping techniques are employed to guarantee high stability and accuracy. However, on the other hand, they make optical clocks an entire optical tableful of equipment, and cannot work continuously for a long time; as a result, they restrict optical clocks used as very convenient and compact time-keeping clocks. In this article, we proposed, and experimentally demonstrated, a novel scheme of optical frequency standard based on comb-directly-excited atomic two-photon transitions. By taking advantage of the natural properties of the comb and two-photon transitions, this frequency standard achieves a simplified structure, high robustness as well as decent frequency stability, which promise widespread applications in various scenarios.

With the development of optical clocks<sup>1–3</sup>, extraordinary results in fundamental research<sup>4,5</sup> as well as practical applications including satellite-based navigation and ranging<sup>6</sup> have been demonstrated in the last decade. Benefited by the development of frequency-stabilization techniques (ultrastable optical cavities)<sup>7–9</sup> for continuous-wave lasers and optical frequency combs<sup>10–12</sup>, traditional optical clocks based on optical transitions in ultra-cold neutral atoms and trapped ions<sup>2,13,14</sup> have achieved a great advance in optical frequency standards. As a result, the state-of-the-art optical clock has reached  $10^{-18}$  level both in accuracy and stability<sup>3</sup>. However, in the meantime, ultrastable optical cavities, laser cooling and particle trapping techniques are susceptible to environments, which hinder the widespread use of optical clocks in many practical applications.

In a conventional optical clock scheme, a pre-stabilised continuous wave laser with narrow linewidth is employed as the probe laser to excite the narrow clock transitions of either a single ion or neutral atoms. To prevent Doppler effects, collision shift and other motion induced disturbances to the clock transition resonances, it is necessary to use laser-cooling and particle-trapping technologies, such as magneto-optical trap and ion trap, to laser-cool the atoms and confine their motion into the Lamb-Dicke regime<sup>2,14</sup>. On the other hand, the probe laser is required to be pre-stabilised to a high-finesse optical reference cavity for linewidth narrowing so as to achieve precise transition resonance probing and laser locking<sup>8,9</sup>. Furthermore, an optical frequency comb operates as the clockwork to faithfully transfer the high stability and accuracy of the optical standard to countable microwave frequency<sup>1,15</sup>.

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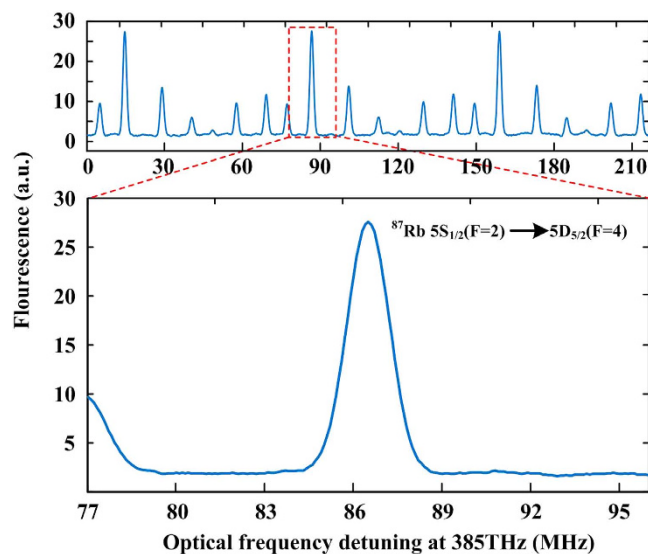
**Figure 1. Three-level scheme used to describe the two-photon transitions excited by frequency comb.** (a) Schematic of the joint participation of comb lines for TPT excitations. Once the comb is resonant with the TPT resonance, comb line pairs can jointly excite TPTs. The 5P state enables the occurrence of stepwise TPTs, involving the  $5S \rightarrow 5P$  (780 nm) and  $5P \rightarrow 5D$  (776 nm) transitions. (b) Doppler-free TPT excitations by two counter-propagating laser beams. For the elimination of Doppler-broadening TPTs background, various coherent quantum techniques are employed in this scheme (see methods).

Much attention of the optical clock research is paid on pushing the accuracy and stability to the quantum limit<sup>16</sup>. Hence, novel material of optical cavity with high elastic modulus and large conductivity is explored to narrow the laser linewidth<sup>8</sup>. Moreover, technologies for particle cooling and trapping, e.g. optical lattice with magic wavelength cooling, have been continuously improved to obtain narrower and more stable transition line. All these efforts enable uncertainty and instability of state-of-the-art optical clock<sup>3</sup> to reach the order of  $10^{-18}$ . However, these devices related to high stable optical cavity, laser cooling and particle trapping usually make the optical clocks occupy large size. Meanwhile, they are highly susceptible to environmental influences, including acoustic noise, mechanical vibrations, temperature drifts, etc. These entire characteristics make optical clocks cannot work continuously for a long time, and restrict the use of optical clocks in various application scenarios beyond the research labs. Until now, the time-keeping clocks used for Coordinated Universal Time (UTC) are still microwave clocks (like Microsemi 5071A), whose stability is much worse than that of optical clocks.

The orientation of this article is to explore a simplified and robust scheme of optical frequency standard that could expedite practical and commercial applications, and ultimately to make a time-keeping optical clock. To achieve this goal, this scheme utilizes the optical frequency comb, instead of the continuous wave laser in conventional scheme, as the probe laser to directly excite the atomic clock transitions, and finally frequency locks one of the comb lines to the centre of transition resonance. With the optical comb simultaneously operating as the probe laser as well as the perfect frequency stability transfer gear between optical and microwave domains, this scheme omits the continuous wave probe laser and relevant inter-laser frequency link servos, which are essential parts in conventional optical clocks. Moreover, in contrast to continuous wave lasers, optical combs have naturally narrow linewidth, usually about kilohertz<sup>17</sup> with simple optimizing, and even on the order of subhertz<sup>18–20</sup> after stabilised to similarly narrow references. This natural advantage enables the direct frequency comb optical frequency standards to avoid the use of delicate optical cavities.

In view of the large amount of comb lines in a frequency comb, each of comb lines shares rather low power (usually on the order of microwatts), which is insufficient for effective atomic transition excitation. To obtain high signal-to-noise ratio (SNR) of transition resonance probing, it is preferable that more comb lines could contribute to the transition excitations. With regard to this requirement, two-photon transition<sup>21</sup> (TPT) would serve as a superior candidate for direct frequency comb optical frequency standards, where, as illustrated in Fig. 1a, a series of comb line pairs with the same sum frequency





**Figure 3. High quality TPTs spectrum of  $^{87}\text{Rb}$  with eliminated Doppler broadening background.** (a) The observed spectrum of  $^{87}\text{Rb}$  Doppler-free TPTs directly excited by the frequency comb with the  $f_{\text{rep}}$  scanned around 144.198 MHz. (b) The  $5S_{1/2}(F=2)$  to  $5D_{5/2}(F=4)$  TPT transition. This TPT has the highest SNR among all of the transitions, and is selected for the stabilization of the frequency comb.

pulses, so as to raise the second-harmonic-generation efficiency of the comb. The comb is then frequency doubled in a periodically poled lithium niobate (PPLN) crystal, resulting in 40 mW second-harmonic (SH) comb centred at 778 nm. TPTs excited by each unidirectional beam (Fig. 1b) give rise to a Doppler-broadening background, which blurs the Doppler-free signals<sup>29,33</sup>. Hence, we utilise the “split-pulse” technique<sup>29</sup> (involving a grating, a lens and a pair of mirrors in Fig. 2) to eliminate the nonresonant type of TPT background. Besides, we greatly eliminate the stepwise (resonant) type of two-photon transition background by tuning the  $f_{\text{rep}}$  to destroy the resonant condition<sup>34</sup>. More details about the background elimination are described in Methods. After frequency shifted by an acousto-optic modulator (AOM, used for comb line locking), the SH comb is focused in the thermal  $^{87}\text{Rb}$  cell. The pulses are reflected back by a mirror placed at a distance of  $c/2f_{\text{rep}}$  (where  $c$  is the speed of light in the space) from the center of the cell, to ensure the two pulses collide temporally in the middle of the Rb cell. The Doppler-free TPTs take place in the centre of the cell when the returned pulses overlap with the counter-propagating pulses.

**Doppler-free two-photon transitions spectrum.** The TPT signals are monitored by detecting the 420 nm fluorescence from cascade decay via the 6P state with a photomultiplier tube (PMT). By scanning the  $f_{\text{rep}}$ , we can obtain the DFCS of the TPTs (Fig. 3a), which illustrates high-resolution  $5S \rightarrow 5D$  transition lines with a  $f_{\text{rep}}$ -scanning repetition period of 27 Hz and weak Doppler-broadening background below 5% intensity of the strongest transition line. Each of the TPT lines performs a linewidth of  $\sim 2$  MHz in optical domain (Fig. 3b), which is mainly due to the 1.5-MHz residual Doppler broadening and 300-kHz transit time broadening with a laser beam waist diameter of  $\sim 150 \mu\text{m}$ . The rest linewidth broadening is related to power broadening and beam misalignment. The Doppler effect can be suppressed to  $\sim 1 \times 10^{-3}$  in thermal Rb atomic system.

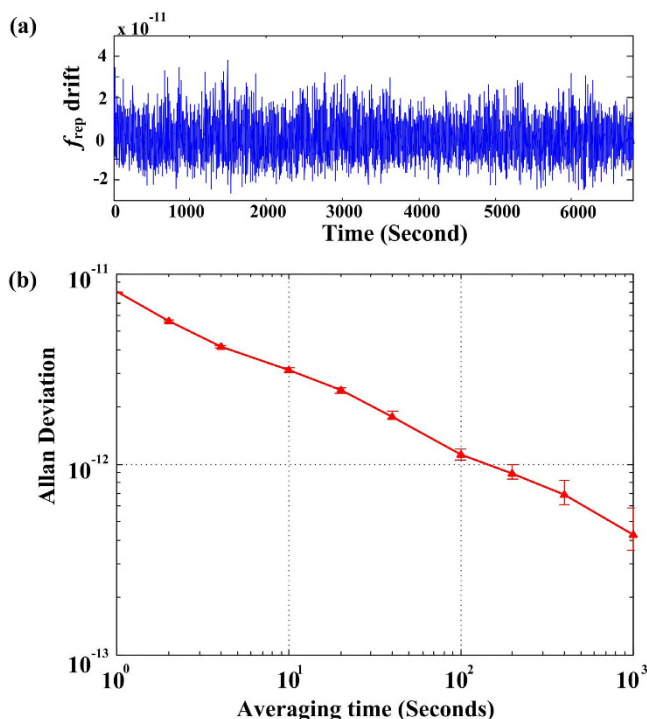
The  $5S_{1/2}(F=2) \rightarrow 5D_{5/2}(F=4)$  transition is chosen as the optical frequency standard. Before the determination of this unperturbed two-photon transition frequency, the evaluation of systematic shifts needs to be investigated. These systematic shifts are presented in Table 1. The main systematic effect in the present work is light shift. This shift scales linearly with the average power of the laser, and can be corrected by performing measurements at different optical powers and extrapolating to zero. In our experimental conditions, this shift is about  $-47.3$  kHz at 22-mW average power.

According to previous study<sup>35</sup>, a collisions shift about  $-181 \pm 16$  Hz is expected at the temperature of  $60 \pm 1$  °C. This is difficult to quantify shifts from impurities in the vapour cell since the pressure of impurities is hardly changing with the temperature. We take a conservative upper limit for the pressure shift equal to 1 kHz. The smaller shifts including the relativistic shift, the black-body radiation shift and the collisions shift are also quantified. The Rb cell is placed in a specially designed  $\mu$ -metal magnetic shield for reducing the influence of external fields, we estimate an upper limit of 703-Hz magnetic shift, providing 0.4 Gauss maximum residual magnetic field. In the locking process, the electronics induces approximately  $-1.6$  kHz frequency shift.

Taking into account the evaluated systematic effects and 5-kHz statistical uncertainty related to the absolute measurement, the  $5S_{1/2}(F=2)$  to  $5D_{5/2}(F=4)$  two-photon transition is measured to be 385 284

Effects	Shifts at 778 nm
Light shift	$-43.2 \pm 3$ kHz
Relativistic shift	$-360 \pm 1$ Hz
Collisions shift	$-181 \pm 16$ Hz
Collisions shift due to background gases	$< 1$ kHz
Black-body radiation shift	$-330 \pm 4$ Hz
Magnetic shift	$< 703$ Hz
Electronic shifts	$-1.6 \pm 0.2$ kHz

**Table 1.** Summary of the systematic frequency shifts.



**Figure 4.** The stability of the TPT stabilised comb. (a) Time records of the fractional frequency drift of the stabilised repetition rate. (b) Fractional frequency instability of the stabilised repetition rate.

566 373(10) kHz. This result is well consistent with previous measurements<sup>27,35</sup>. The accuracy of this frequency standard can reach  $2.5 \times 10^{-11}$ .

**Frequency instability of the direct frequency comb optical frequency standard based on two-photon transitions of thermal  $^{87}\text{Rb}$  atoms.** After getting the high-resolution TPT lines, the next step is to lock comb lines to the  $5S_{1/2}$  ( $F=2$ )  $\rightarrow$   $5D_{5/2}$  ( $F=4$ ) transition, which has the maximum SNR among all the TPTs of  $^{87}\text{Rb}$ . For this purpose, an audio frequency signal is applied on the microwave driver (80 MHz) of the AOM. By detecting the phase signal from the 420 nm fluorescence, a digital lock-in-amplifier is employed to generate the error signal, and feeds it back into the actuator of the repetition rate (see Methods). With the  $f_0$  previously frequency locked to the  $f_{\text{rep}}$ , the comb is fully stabilised to the TPT resonance. The curve in Fig. 4(a) shows the fractional frequency drifts of the fundamental repetition rate at 144 MHz, which is stabilised within 6 mHz. Figure 4(b) shows the measured fractional frequency instabilities in terms of overlapping Allan deviation. The instability reaches  $8.0 \times 10^{-12}$  and  $4.3 \times 10^{-13}$  for averaging times of 1 s and 1,000 s, respectively. The frequency instability of our system is mainly limited by the instability of the fibre laser, which induces frequency noise from the environment and the light shift caused by variations in the optical power.

## Discussion

The demonstrated optical frequency standard takes advantage of the narrow-linewidth optical comb to directly excite the narrow TPTs of thermal atoms, and then frequency locks the comb to these



transitions. The optical comb enables the use of various coherent quantum techniques for significant Doppler-broadening elimination and SNR improvement<sup>29,33,36</sup>. Although the currently demonstrated system is still a tableful system, most parts are fibre-based components, and can be easily integrated into a small box. Moreover, the robustness of the whole system can be greatly improved by using an all polarization-maintaining, all fibre frequency comb<sup>37,38</sup> as the probe laser. This is also our future work. The current stability demonstrated here is only a first step in this direction, it is a promising indicator that more advances can be made to achieve  $10^{-13}$  instability with respect to the further reduction of Doppler-broadening and improvement of SNR.

This proof-of-concept system can also be expanded to other narrower TPTs, like the ultraviolet 1S-2S two-photon resonance with a natural linewidth of  $\sim 1$  Hz in atomic hydrogen<sup>16</sup>, and can obtain a better stability and accuracy. When combining with the cold atomic system, it can be easy to realize  $10^{-15}$  instability and uncertainty<sup>39</sup>, which is very close to the optical clock based on the ultrastable optical cavities, laser cooling and particle trapping techniques. Ultimately, the direct frequency comb optical frequency standards based on TPTs of thermal atoms are likely to benefit many practical applications, such as timing and synchronization systems, ultra-broadband and secure communications systems.

## Methods

**Elimination of the Doppler broadening background.** We use a passive Er: fibre femtosecond pulse laser with a repetition frequency of  $\sim 144$  MHz and an optical bandwidth of about 30 nm at a centre wavelength of 1,556 nm, as the probe laser and frequency standard output. In the observation of the Doppler-free Rb TPT transition directly excited by two counterpropagating laser beam, the undesired Doppler broadening transitions including direct TPTs (D-TPTs) and sequential TPTs (S-TPTs) are also happening at the same time and result in a Doppler-broadening background, and degrade the stability. Therefore, various coherent quantum techniques are utilized to estimate the Doppler-broadening background. As Fig. 2 shows, a 1200 grooves/mm grating is used to disperse the beam. Two mirrors placed parallel to the Fourier plane are used to reflect  $>778$  nm part and  $<778$  nm part of the frequency comb spectrum. The  $<778$ -nm and  $>778$ -nm part are introduced by different time delays, therefore the two dispersed pulses cannot arrive at the vapour cell at the same time, so that they are inhibited to simultaneously excite the unidirectional D-TPTs.

For the S-TPTs Doppler-broadening background excited by one unidirectional beam, two of the comb lines are required to be resonant with the  $5S \rightarrow 5P$  (780 nm) and  $5P \rightarrow 5D$  (776 nm) transitions, respectively. When  $f_{\text{rep}}$  is a subharmonic of the frequency difference  $\Delta f$  between the two-step  $5S \rightarrow 5P \rightarrow 5D$  transitions, i.e.,  $f_{\text{rep}} = \Delta f / N$  ( $N$  is an integer), there are always some velocity groups of atoms to be excited by different mode pairs of the comb and contribute to the S-TPTs Doppler-broadening background. For the elimination of S-TPTs background, the  $f_{\text{rep}}$  is far detuned from the resonance values  $\Delta f / N$ . More detailed information about these coherent quantum techniques can be found in ref. 34.

**Stabilisation of the carrier-envelope offset frequency.** To detect the  $f_0$ , we employ the conventional  $f$ -to- $2f$  self-referencing technique, as shown in Fig. 2. The output of the femtosecond laser is amplified to 200 mW. The amplified pulses are injected into a highly nonlinear fibre (HNLF) to generate a supercontinuum spectrum. The spectrum around 2060 nm in the supercontinuum spectrum is doubled by second-harmonic generation in a PPLN. This generated 1030-nm light beats with the fundamental 1030-nm light in an avalanche photodiode (APD). The beating signal  $f_0$  at 40 MHz is then amplified and filtered, mixed with 1-GHz microwave, and divided by 32 in frequency. This signal is digitally phase-locked with a signal from direct digital synthesizer referenced to the  $f_{\text{rep}}$ , by controlling the pump current of the fibre comb.

**Stabilisation of the repetition rate.** Before the optical pulses enter the Rb cell, we firstly frequency shift them by 80 MHz in the AOM. For the phase locking, this 80-MHz signal is frequency modulated by an audio frequency ( $\sim 780$  Hz). After frequency-shifted, modulated, the counter-propagating pulses collide in the centre of the cell and excite the Doppler-free TPTs. We use a PMT to detect the 420-nm fluorescence from the cascaded  $5D \rightarrow 6P \rightarrow 5S$  spontaneous radiation. The electronic signal from the PMT contains the phase information of the audio frequency, and is delivered into a digital lock-in-amplifier to generate the error signal. We feed it back into the intracavity piezoelectric transducer to stabilise the repetition rate.

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## Acknowledgements

This work was supported in part by the Nature Science Foundation of China under Grant No.61535001 and No.61371074, and sponsored by the Research Fund of ZTE Corporation.

## Author Contributions

J.Y.Z. and Z.G.Z. developed the concept. J.T.W., S.Y.Z. and Z.G.Z. designed the fibre frequency comb and optical paths. J.Y.Z., Y.L.Z., S.Y.Z. and J.X.L. designed the RF signal-processing circuits and electronic servo systems. S.Y.Z., J.T.W. and J.Y.Z. collected the data. S.Y.Z. and W.P.Y. did the theoretic calculation. All authors joined in the discussion and provided comments. S.Y.Z. and J.T.W. contributed equally to this work.

## Additional Information

**Competing financial interests:** The authors declare no competing financial interests.

**How to cite this article:** Zhang, S. Y. *et al.* Direct frequency comb optical frequency standard based on two-photon transitions of thermal atoms. *Sci. Rep.* **5**, 15114; doi: 10.1038/srep15114 (2015).



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