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# Facet-dependent electrochemical properties of Co<sub>3</sub>O<sub>4</sub> nanocrystals toward heavy metal ions

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We revealed an interesting facet-dependent electrochemical behavior toward heavy metal ions (HMIs) based on their adsorption behaviors. The (111) facet of  $\text{Co}_3\text{O}_4$  nanoplates has better electrochemical sensing performance than that of the (001) facet of  $\text{Co}_3\text{O}_4$  nanocubes. Adsorption measurements and density-functional theory (DFT) calculations reveals that adsorption of HMIs is responsible for the difference of electrochemical properties. Our combined experimental and theoretical studies provide a solid hint to explain the mechanism of electrochemical detection of HMIs using nanoscale metal oxides. Furthermore, this study not only suggests a promising new strategy for designing high performance electrochemical sensing interface through the selective synthesis of nanoscale materials exposed with different well-defined facets, but also provides a deep understanding for a more sensitive and selective electroanalysis at nanomaterials modified electrodes.

iming at a more sensitive electroanalysis, a variety of nanomaterials or nanocomposites modified electrodes were very often explored for electrochemical detection of very trace levels of toxic heavy metal ions (HMIs). Very recently, nanoscale metal oxides as novel modifiers have been reported in electrochemical detection of HMIs<sup>1-5</sup>. Compared with traditional modifiers such as noble metals and biomolecules, the electrodes fabricated with nanoscale metal oxides are easy to synthesize with low cost<sup>1,3,5</sup>. However, as is often claimed, increased currents and increased analytical sensitivity are simply reflections of an increased microscopic surface area and not electrocatalytic activity or any other nano-effects. That is, the mechanism of using nanoscale metal oxides in electrochemical sensing of HMIs has not been proposed by considering their insulating property. Therefore, the design and implementation of new experimental approach combined with theoretical studies are extremely needed.

For a crystalline nanoscale material, different facets may have different geometric and electronic structures, and exhibit different physical and chemical properties  $^{6-8}$ . Much more attention has been given to investigate the facet effects on catalysts, photocatalysts, electrocatalysts, Li-ion battery, supercapacitors and so on  $^{9-24}$ . Take  $Co_3O_4$  nanoscale materials as an example; Li *et al.* investigated the facet effect of  $Co_3O_4$  on catalytic property for methane combustion  $^{10}$ . Xiao *et al.* reported that the exposed facets of  $Co_3O_4$  nanocrystals are very important for Li<sup>+</sup> transport in Li-ion battery  $^{12}$ . Although oriented  $Co_3O_4$  nanostructures have attracted wide attention, the direct experimental detection of HMIs with controlled crystalline morphology and orientation remains a significant challenge- that is, the facet effect in electrochemical sensor especially detection of HMIs is not proposed. Further, till now nanoscale  $Co_3O_4$  has not been used to examine the electrochemical sensing of HMIs. It should be also pointed out that plenary theoretical investigations are expected to help to get deeper insight into the crystal facet effect, but most reports on the crystal facet effect did not adequately combined first-principles theoretical studies at atomic level with the experimental results  $^{13,14,17,19-21,24}$ .

In this work, we report the facile synthesis of  $\mathrm{Co_3O_4}$  with two different shapes. Using this non-conductive nanomaterial-modified electrode, we try to demonstrate the facet-dependent electrochemical behaviour of  $\mathrm{Co_3O_4}$  nanocrystals toward HMIs by combining the adsorption measurements and the density-functional theory (DFT) calculations. It has been recognized that the (111) facet of  $\mathrm{Co_3O_4}$  nanoplates has better electrochemical sensing performance than that of the (001) facet of  $\mathrm{Co_3O_4}$  nanocubes. Adsorption measurements and DFT calculations reveal that adsorption of Pb(II) is responsible for the difference of electrochemical



properties. To the best of our knowledge, this is the first study to investigate the effect of facet on electrochemical sensing behaviour toward HMIs. Our combined experimental and theoretical studies provide a solid hint to explain the mechanism of electrochemical detection of HMIs using nanoscale metal oxides. This strategy may be extended to other electrochemical sensors based on nanoscale metal oxides.

### Results

Two simple routes have been designed to fabricate Co<sub>3</sub>O<sub>4</sub> nanocubes and nanoplates (Supporting Information, Figure S1). Co<sub>3</sub>O<sub>4</sub> nanocubes were synthesized by a one-step hydrothermal method with Co(CH<sub>3</sub>COO)<sub>2</sub> as the only reactant. The Co<sub>3</sub>O<sub>4</sub> nanoplates were obtained by a solvothermal method in ethylene glycol (EG) followed by a calcination treatment in air. Figure 1 shows the representative scanning electron microscopy (SEM), transmission electron microscopy (TEM) and high resolution TEM (HRTEM) images of Co<sub>3</sub>O<sub>4</sub> nanoplates and nanocubes. The TEM and HRTEM images of Co<sub>3</sub>O<sub>4</sub> nanocubes are presented in Figure 1a and b, respectively. As seen, the particle size of Co<sub>3</sub>O<sub>4</sub> nanocubes is about 20 ~ 40 nm. HRTEM (Figure 1b) and corresponding selected area electron diffraction (SAED) pattern (Inset in Figure 1b) indicate that the lattice fringe is 0.28 nm and the nanocube is exposed with six (001) facets. A typical SEM image taken for Co<sub>3</sub>O<sub>4</sub> nanoplates is shown in Figure 1c. The low magnification SEM and TEM images in Figure S2 demonstrated the homogeneity of Co<sub>3</sub>O<sub>4</sub> nanoplates. It is seen that most of the nanoplates display well-defined hexagonal shape. The width and the thickness of the plates is determined to be 700  $\sim$ 900 nm and 50  $\sim$  60 nm, respectively. The HRTEM image in Figure 1d clearly shows that the spacing between lattice fringes with an angle of  $60^{\circ}$  is 0.28 nm, which is consistent with the ( $\overline{2}02$ ) and  $(0\overline{2}2)$  planes of cubic Co<sub>3</sub>O<sub>4</sub>. And it is also confirmed by the SAED pattern (Inset in Figure 1d) revealing a single crystal structure. Thus, the resulting dominant exposed plates of Co<sub>3</sub>O<sub>4</sub> nanoplates are (111) facets. Powder X-ray diffraction patterns (XRD) (Supporting Information, Figure S3) of both nanocubes and nanoplates have identical peaks, which can be perfectly indexed to that of cubic spinel Co<sub>3</sub>O<sub>4</sub> (Joint Committee on Powder Diffraction Standards (JCPDS) card no. 42-1467). No impurities have been detected, indicating the formation of pure cobalt oxides. The crystal sizes were calculated according to Sherrer equation based on the (311) diffraction peak of XRD data. The crystal size of Co<sub>3</sub>O<sub>4</sub> nanocubes was 32.5 nm, which was consistent with the TEM observation. However, the crystal size of Co<sub>3</sub>O<sub>4</sub> nanoplates was calculated to be 21.3 nm, which may be due to its porous structure. The specific surface areas of Co<sub>3</sub>O<sub>4</sub> nanocubes and nanoplates have been measured by the Brunauer-Emmett–Teller (BET) method (Supporting Information, Figure S4). The measured specific surface areas for Co<sub>3</sub>O<sub>4</sub> nanocubes and nanoplates are 20 and 13 m<sup>2</sup> g<sup>-1</sup>, respectively.

In general, the obtaining of metal oxides requires a basic medium, and usually NaOH is employed for this purpose. However, in the fabrication of  $\text{Co}_3\text{O}_4$  nanocubes, the hydrolysis of acetate anions from  $\text{Co}(\text{CH}_3\text{COO})_2$  becomes a good source of hydroxyl anions (Equation (1)). The  $\text{O}_2$  dissolved in the deionized water is supposed to act as the main oxidant. Partial Co(II) is transformed into Co(III) and the coexistence of Co(II) and Co(III) under the ambient environment prefers the formation of  $\text{Co}_3\text{O}_4$  (Equation (2)). The low amount of hydroxyl anions in the solution may promote the formation of  $\text{Co}_3\text{O}_4$  nanocubes with (001) planes. Similar result has also been reported by Xiao et al<sup>12</sup>.

$$CH_3COO^- + H_2O \rightarrow CH_3COOH + OH^-$$
 (1)

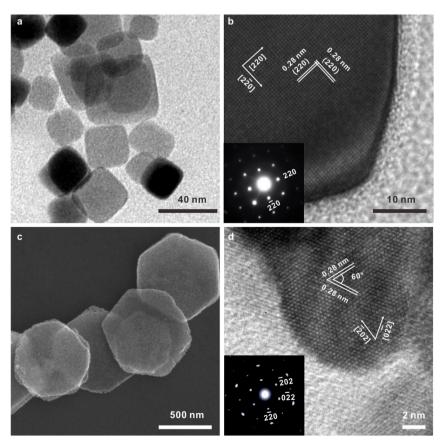


Figure 1 |  $Co_3O_4$  nanocrystals. (a-b), TEM and HRTEM image of typical  $Co_3O_4$  nanocubes. Inset in panel (b): SAED pattern of nanocubes. (c-d), SEM and HRTEM image of typical  $Co_3O_4$  nanoplates. Inset in panel (d): SAED pattern of nanoplates.



$$6Co^{2+} + 12OH^{-} + O_2 \rightarrow 2Co_3O_4 + 6H_2O$$
 (2)

From the SEM and TEM images of the  $\text{Co}_3\text{O}_4$  nanoplates precursors before calcination (Data not shown), we can see that the plate-like structures of  $\text{Co}_3\text{O}_4$  are inherited from these precursors. The typical XRD pattern of the precursors is shown (Supporting Information, Figure S3, black line). A strong diffraction peak around  $10^\circ$  in the XRD pattern is the characteristic of metal glycolates<sup>25–28</sup>. It has been assumed that EG would lose its two protons and the dianion complex with metal center<sup>25</sup>. Xia et al. has discussed the oligomerization process of metal glycolates. Longer chains of cobalt glycolate oligomers could self-assemble into ordered bundles (i.e., nanoplates) through van der Waals interactions and then precipitate out from the reaction medium<sup>25</sup>. After calcination,  $\text{Co}_3\text{O}_4$  exposed mainly with (111) facets are obtained. However, the reason for the exposed (111) facets after calcination is not clear.

Co<sub>3</sub>O<sub>4</sub> nanocubes and nanoplates modified glassy carbon electrodes (GCEs) are first electrochemically characterized by cyclic voltammetric (CV) and electrochemical impedance spectrum (EIS) (Supporting Information, Figure S5). As compared with the bare GCE, the anodic and cathodic peaks decrease at the Co<sub>3</sub>O<sub>4</sub> modified electrode, which demonstrates that Co<sub>3</sub>O<sub>4</sub> nanomaterials have been modified onto the surface of the GCE (Supporting Information, Figure S5a). It also indicates that the rate of electron transfer at the electrode surface is hindered with the attachment Co<sub>3</sub>O<sub>4</sub> onto the GCE surface. In a typical Nyquist plot, the semicircle proton corresponding to the electron-transfer resistance ( $R_{\rm et}$ ) at higher frequency range while a linear part at lower frequency range represents the diffusion limited process<sup>3</sup>. The value corresponding to the bare GCE is about 27  $\Omega$  (Supporting Information, Figure S5b). And the value of  $Co_3O_4$  nanocubes modified GCE (48  $\Omega$ ) is comparable to that of  $Co_3O_4$  nanoplates (50  $\Omega$ ) (Supporting Information, Figure S5b). In this case, the surface area of the modified electrodes were calculated to be 0.0534 (nanoplates), 0.0519 (nanocubes), and 0.07065 (bare GCE) cm<sup>2</sup>. This may be due to the insulating property of Co<sub>3</sub>O<sub>4</sub> nanostructures, as such, the direct contribution of microscopic surface area could be avoided, which is very helpful to understand that the difference of electrochemical behavior is indeed from the effect of crystal facet (as will be discussed in the following).

Among all of the HMIs, lead has been identified as one of the most toxic heavy metals because of its detrimental effects on the human nervous system, blood circulation system, kidneys, and reproductive system<sup>29</sup>. Subsequently, the performances of these two types of Co<sub>3</sub>O<sub>4</sub> with different exposed facets in electrochemical detection of heavy metal ions are comprehensively investigated. Pb(II) is used as a probe heavy metal ion. We first examined the square wave stripping voltammograms (SWASV) responses to Pb(II) in 0.1 M NaAc-HAc (pH = 5.0) at  $Co_3O_4$  nanoplates and nanocubes modified electrodes (Figure S6). There is almost no obvious response at bare GCE, and at the meanwhile, a weak stripping response for nanocubes/GCE is observed. However, a strong and well-defined peak at −0.584 V is clearly seen for nanoplates modified GCE. And the peak current obtained is about 8 times and 13 times that of nanocubes/GCE and bare GCE, respectively. Moreover, the stripping peak shifts toward more positive potential and become progressively less symmetrical. This is the consequence of having more ions on the electrode (this is fully consistent with that more ions are adsorbed by nanoplates and release to the bare GCE, see Discussion section), and therefore requiring a longer sweep to remove the metal from the surface entirely. Figure 2 shows that the current densities increase linearly versus the Pb(II) concentrations. For nanocubes/GCE, the linearization equation is  $j/(\mu A \text{ cm}^{-2}) = -122.9 + 174.4 \text{ c/}\mu M$ . While for nanoplates/GCE, the linearization equation is j/( $\mu A \text{ cm}^{-2}$ ) = -131.1+ 372.7 c/ $\mu$ M. The results indicate that Co<sub>3</sub>O<sub>4</sub> nanoplates with (111) facet exhibit better electrochemical detection performance than

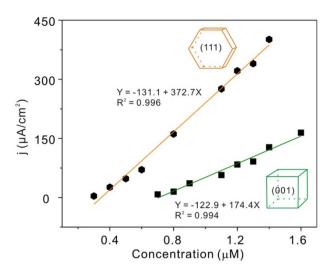


Figure 2 | Electrochemical properties of  $\text{Co}_3\text{O}_4$  nanocrystals. Calibration plots of  $\text{Co}_3\text{O}_4$  nanoplates and  $\text{Co}_3\text{O}_4$  nanocubes modified electrode toward Pb(II) at different concentrations in 0.1 M NaAc-HAc solution (pH = 5.0). Considering the weak response at nanocubes modified electrode, the initial concentration is from 0.7  $\mu\text{M}$ .

Co<sub>3</sub>O<sub>4</sub> nanocubes with (001) facet. The sensitivity of Co<sub>3</sub>O<sub>4</sub> nanoplates/GCE (372.7  $\mu$ A cm<sup>-2</sup>/ $\mu$ M) is over 2 times that of Co<sub>3</sub>O<sub>4</sub> nanocubes/GCE (174.4  $\mu$ A cm<sup>-2</sup>/ $\mu$ M). The SWASV responses toward Pb(II) at various concentrations on Co<sub>3</sub>O<sub>4</sub> nanoplates and nanocubes modified electrodes were seen in Figure S7. The limit of detection (LOD) as low as 0.12 nM and 0.16 nM (3 $\sigma$  method) for Co<sub>3</sub>O<sub>4</sub> nanoplates and nanocubes, respectively was achieved. This meets the requirements of the World Health Organization (WHO) maximum permissible limit for lead concentration in drinking water of 10  $\mu$ g L<sup>-1</sup>. The LODs are better than existing methods based on noble metal and ion-imprinted polymers, such as gold nanofilm (0.1  $\mu$ M)<sup>30</sup> and nano-sized Pb<sup>2+</sup> imprinted polymer (0.6 nM)<sup>31</sup>, etc. The LODs are also superior to some oxides, such as tube-in-tube SnO<sub>2</sub> (1.6 nM)<sup>32</sup> and SnO<sub>2</sub>/graphene nanocomposites (0.18 nM)<sup>33</sup>.

### **Discussion**

We suggest that the difference in electrochemical behavior may have a relationship with the adsorption capacities of different nanoscale Co<sub>3</sub>O<sub>4</sub> toward Pb(II). To confirm this contribution, the adsorption measurement is conducted in 0.1 M NaAc-HAc (pH = 5.0) in order to be consistent with electrochemical sensing conditions. Co<sub>3</sub>O<sub>4</sub> nanocubes and nanoplates are incubated in Pb(II) solutions with different initial concentrations for 24 h to reach the adsorption equilibrium. The adsorption isotherms are shown in Figure 3a and b. Both the adsorption isotherms fit the Langmuir isotherm very well (Supporting Information, Figure S8a and b). The maximum adsorption capacity  $(q_{\text{max}})$  of Co<sub>3</sub>O<sub>4</sub> nanoplates (33.8 mg g<sup>-1</sup>) is about 19 times that of  $Co_3O_4$  nanocubes (1.8 mg g<sup>-1</sup>). In previous result we have shown that Co<sub>3</sub>O<sub>4</sub> is non-conductive material (Supporting Information, Figure S5). We believe that the higher adsorption performance of Co<sub>3</sub>O<sub>4</sub> nanoplates than that of Co<sub>3</sub>O<sub>4</sub> nanocubes lead to its better performance in electrochemical detection. It is important to emphasize that as the adsorption experiments were done in solution with dispersed Co<sub>3</sub>O<sub>4</sub>, while the Co<sub>3</sub>O<sub>4</sub> were just coated onto the surface of the GCE, the electrochemical sensitivities for two types of Co<sub>3</sub>O<sub>4</sub> would be not proportional to the adsorption capacities of them.

The above results and discussions experimentally show that  $\text{Co}_3\text{O}_4$  nanocubes and nanoplates exhibit obviously different electrochemical sensing and adsorption performances.  $\text{Co}_3\text{O}_4$  nanoplates with (111) planes exhibit larger sensitivity and adsorption capacity for heavy metal ions. Besides, we suggest that Pb(II) on the (111) facet are easier to diffuse onto GCE for reduction and



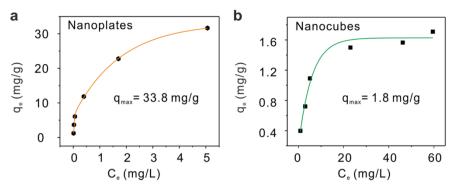


Figure 3 Adsorption measurements. Adsorption isotherms of Pb(II) onto (a), Co<sub>3</sub>O<sub>4</sub> nanoplates and (b), Co<sub>3</sub>O<sub>4</sub> nanocubes.

stripping reaction than those on the (001) facet. Regarding the above surface area analysis of two kinds of  ${\rm Co_3O_4}$  (nanocubes > nanoplates), this further confirms that the differences in their electrochemical sensing performances are due to the facet effect.

To better understand the effect of facet on electrochemical behavior of Pb(II) on Co<sub>3</sub>O<sub>4</sub> nanoplates and nanocubes, we perform DFT calculations to explore the adsorption and diffusion behaviors of Pb on Co<sub>3</sub>O<sub>4</sub>(001) and (111) surfaces. The adopted computational models are similar to the recent report on Li storage capability of Ti<sub>3</sub>C<sub>2</sub> and  $Ti_3C_2X_2$  (X = F, OH) monolayer<sup>34</sup>. The adsorption configurations of Pb atom on Co<sub>3</sub>O<sub>4</sub>(001) and (111) surfaces are optimized without any symmetry constraint. In the most energetically stable configuration of Pb/Co<sub>3</sub>O<sub>4</sub>(001) system, the Pb atom is four-coordinated by two oxygen and two Co surface atoms of Co<sub>3</sub>O<sub>4</sub>(001) surface (Supporting Information, Figure S9), and the Pb-O and Pb-Co bond lengths are 2.40 and 2.27 Å, respectively. For Pb/Co<sub>3</sub>O<sub>4</sub>(111) adsorption system, the Pb atom locates at the three-fold hollow site (Supporting Information, Figure S10), and the Pb-O bond length is 2.27 Å. Two corresponding adsorption configurations of Pb atom on  $Co_3O_4(001)$  and (111) surfaces are illustrated in top panel of Figure 4 labeled with i and ii, respectively. The calculated adsorption energy of Pb atom on Co<sub>3</sub>O<sub>4</sub>(111) surface is -4.11 eV, which is larger than that on  $Co_3O_4(001)$  surface (-3.70 eV). The relative large adsorption energy for Pb/Co<sub>3</sub>O<sub>4</sub>(111) system mainly originates from the relative short Pb-O bond length. The Bader charge analysis<sup>35</sup> (Supporting Information, Figure S11) shows that the adsorbed Pb atom carries positive +0.36 and +0.10 |e| (e, the electron charge) on Co<sub>3</sub>O<sub>4</sub>(111) and (001) surfaces, respectively, which gives the quantitative proof of the relative strong Pb-substrate interaction for Pb/ Co<sub>3</sub>O<sub>4</sub>(111) system. Note that there are four equivalent stable adsorption sites in a  $Co_3O_4(001) - (\sqrt{2} \times \sqrt{2})$  cell and  $Co_3O_4(111) - (1 \times 1)$  cell (Supporting Information, Figure S9 and S10), but the area of the former cell is larger than that of the latter one. In addition, the adsorption energy for Pb atom locating at the second stable adsorption site (labeled with II' symbol within twelve cyan circles in Supporting Information, Figure S10c) on  $Co_3O_4(111)$  surface is predicted to be -3.81 eV, which is also larger than that of the most stable adsorption configuration of Pb/ Co<sub>3</sub>O<sub>4</sub>(001) system. On the other hand, we have examined the full-coverage of Pb on the stable adsorption sites. The adsorption energy of each Pb on Co<sub>3</sub>O<sub>4</sub>(001) and (111) surfaces is predicted to be -3.38 and -3.69 eV, respectively. It is clear that they are consistent with the calculated results of low-coverage of Pb on Co<sub>3</sub>O<sub>4</sub> facets. Thus, the Pb atom adsorption ability on Co<sub>3</sub>O<sub>4</sub>(111) surface is significantly larger than that on Co<sub>3</sub>O<sub>4</sub>(001) surface. Moreover, the high adsorption energies and the large charge transfer suggest that

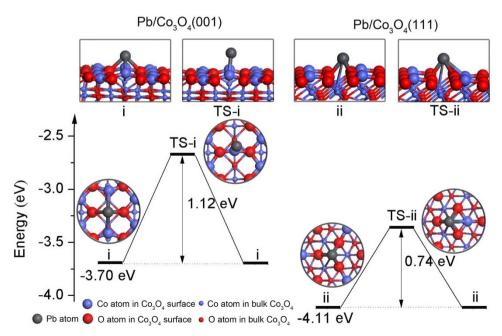


Figure 4 | DFT calculations. Top panel: side view of the optimized stable adsorption and transition-state (TS) structures for Pb on  $Co_3O_4(001)$  and (111) surfaces. Bottom panel: the TS barriers of Pb on  $Co_3O_4(001)$  and (111) surfaces, and top view of these optimized stable adsorption and TS structures are also shown. After overcoming the TS barriers, Pb reaches the nearest neighboring stable adsorption site.

Q

the adsorbed Pb can form a strong Coulomb interaction with the  $Co_3O_4$  surface.

Now we turn to calculate the energy barriers of Pb diffusing on  $\text{Co}_3\text{O}_4(001)$  and (111) surfaces using the climbing image nudged elastic band (CI-NEB) method<sup>36</sup>. The Pb has two diffusion paths from the most stable adsorption site to its nearest neighboring stable site on  $\text{Co}_3\text{O}_4(001)$  surface, while there are six diffusion directions for Pb/Co<sub>3</sub>O<sub>4</sub>(111) adsorption system (Supporting Information, Figure S12). The geometric structures of the transition-states for Pb on  $\text{Co}_3\text{O}_4(001)$  and (111) surfaces are shown in Figure 4 labeled with TS-i and TS-ii, respectively. The calculated diffusion energy barrier of Pb on  $\text{Co}_3\text{O}_4(001)$  and (111) surfaces are predicted to be 1.12 and 0.74 eV, respectively. The relative low transition-state barrier results in the Pb fast diffusing on  $\text{Co}_3\text{O}_4(111)$  surface. These DFT results confirm that  $\text{Co}_3\text{O}_4(111)$  surface can capture more Pb than  $\text{Co}_3\text{O}_4(001)$  surface, and the adsorbed Pb diffuse more easily on the  $\text{Co}_3\text{O}_4(111)$  surface, consistent well with experimental observations.

Based on the above experimental and DFT calculated results, a schematic illustration of how adsorptive nanoscale materials with different exposed crystal facets could be designed for electrochemical sensing interface is depicted in Figure 5. As for bare GCE, there are not nanoscale materials which can effectively capture HMIs and the weaker stripping peak is obtained. Large amount of HMIs could be adsorbed onto the surface of nanoscale materials and then diffuse to the GCE surface. The more target HMIs are adsorbed onto the surface of nanoscale material, the more HMIs diffuse to the GCE and the stronger the stripping peak current obtains as a result. As Co<sub>3</sub>O<sub>4</sub> nanoplates exposed with (111) facet can adsorb more metal ions than Co<sub>3</sub>O<sub>4</sub> nanocubes with (001) facet and the metal ions on the surface of (111) facet can diffuse more easily onto the GCE, Co<sub>3</sub>O<sub>4</sub> nanoplates modified GCE obtain larger stripping peak current. In contrast to previous studies<sup>37–41</sup> focusing on elevating the specific surface area of nanoscale materials in electrochemical sensing; our findings demonstrate that the construction of nanoscale materials onto electrode with well-defined exposed facets is crucial when considering the design of high-performance electrochemical sensor.

In summary,  $Co_3O_4$  nanocubes and nanoplates have been successfully synthesized by facile methods by using different solvents. The predominantly exposed facets are (001) in the  $Co_3O_4$  nanocubes and (111) in the  $Co_3O_4$  nanoplates, respectively. Studies of their sensing properties revealed that the as-prepared  $Co_3O_4$  nanocrystals exhibit interesting facet-dependent electrochemical behaviors toward HMIs

based on their adsorption behaviors. The (111)-bound  ${\rm Co_3O_4}$  nanoplates are superior to (001)-bound  ${\rm Co_3O_4}$  nanocubes. Adsorption measurements revealed that  ${\rm Co_3O_4}$  (111) facets can adsorb more metal ions than the (001) facets. DFT calculations suggest that  ${\rm Co_3O_4}$ (111) facets exhibit a relative larger adsorption energy, more adsorption sites, and a relative lower transition-state barrier than  ${\rm Co_3O_4}$ (001) facets, which may be the predominant reason accounting for the facet-dependent electrochemical sensing behavior toward HMIs. Both adsorption experiments and DFT results agree well with the electrochemical sensing results. This study not only suggests a promising new strategy for designing high performance electrochemical sensing interface through the selective synthesis of nanoscale materials exposed with different well-defined facets, but also provides a deep understanding for a more sensitive and selective electroanalysis at nanomaterials modified electrodes.

### **Methods**

**Materials.** All reagents were commercially available from Sinopharm Chemical Reagent Co., Ltd (China) with analytical grade. Stock solution used in electrochemical measurement of Pb(II) was prepared by dissolving Pb(NO<sub>3</sub>)<sub>2</sub> in deionized water. 0.1 M acetate buffer (NaAc-HAc) solution of pH = 5.0 was prepared with NaAc and HAc. Ultrapure fresh water obtained from a Millipore water purification system (MilliQ, specific resistivity > 18 MΩ cm, S.A., Molsheim, France) are used in all runs.

Synthesis of  $\text{Co}_3\text{O}_4$  nanocubes. 0.25 g of  $\text{Co}(\text{CH}_3\text{COO})_2\cdot 4\text{H}_2\text{O}$  was loaded into a 23 mL poly(tetrafluoroethylene) (PTFE)-lined stainless steel autoclave, which was then filled with 18 mL water. The autoclave was sealed and maintained at 200°C for 12 h, and then cooled down to room temperature. The final products were centrifuged, rinsed with distilled water and ethanol for several times to remove any impurities.

Synthesis of  $\text{Co}_3\text{O}_4$  nanoplates. 0.25 g of  $\text{Co}(\text{CH}_3\text{COO})_2 \cdot 4\text{H}_2\text{O}$  was loaded into a 23 mL poly(tetrafluoroethylene) (PTFE)-lined stainless steel autoclave, which was then filled with 18 mL ethylene glycol. The autoclave was sealed and maintained at 200°C for 12 h, and then cooled down to room temperature. The final products were entrifuged, rinsed with distilled water and ethanol several times to remove any impurities. The as-prepared precursors were finally calcined at 350°C in air for 3 h.

Characterization. The SEM images were taken by a FEI Quanta 200 FEG field emission scanning electron microscope. The TEM and HRTEM images analyses were carried out on a JEM-2010 microscope. XRD was performed on a D/MaxIIIA X-ray diffractometer (Rigaku Co., Japan), using Cu K $\alpha$  ( $\lambda$ K $\alpha$ 1 = 1.5418 Å) as the radiation source. The nitrogen adsorption and desorption isotherms at 77 K were measured with a Micromeritics ASAP 2020 M analyzer. The Brunauer, Emmett, and Teller (BET) equation was used to obtain the specific surface areas. The Pb(II) concentrations were determined in the liquid phase using inductively coupled plasma atomic emission spectrometry (ICP-AES, Jarrell-Ash model ICAP 9000). Electrochemical experiments were recorded using a CHI 660D computer-controlled

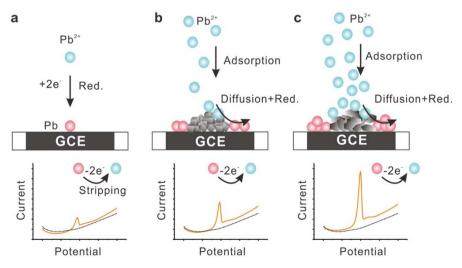


Figure 5 | Schematics of how adsorptive nanoscale materials exposed with different crystal facets could be designed to enhance the performance of electrochemical sensing. (a), Bare GCE. (b),  $Co_3O_4$  nanocubes modified GCE. (c),  $Co_3O_4$  nanoplates modified GCE. Facet-dependent electrochemical properties of  $Co_3O_4$  nanocrystals toward heavy metal ions (e.g.  $Pb^{2+}$ ) are investigated. The (111)-bound  $Co_3O_4$  nanoplates are superior to (001)-bound  $Co_3O_4$  nanocubes. Adsorption measurements reveal that  $Co_3O_4$  (111) facet can adsorb more metal ions than the (001) facet. DFT calculations demonstrate that Pb on  $Co_3O_4$ (111) facet exhibits larger adsorption energies, more adsorption sites, and faster diffusion than on (001) facet.



potentiostat (ChenHua Instruments Co., Shanghai, China). A conventional three-electrode system consisted of a glassy carbon working electrode (GCE, 3 mm diameter), an Ag/AgCl as the reference electrode and a platinum wire as the counter electrode.

Fabrication of modified electrode. 2 mg of the as-prepared  $Co_3O_4$  nanocubes and nanoplates were dispersed in 4 mL water with ultrasonic agitation to give a homogeneous solution. Prior to each modification, the bare GCE was sequentially polished with 1.0  $\mu$ m and 0.05  $\mu$ m alumina power slurries to a mirror shiny surface, and then sonicated with 1:1 HNO<sub>3</sub> solution, absolute ethanol and deionized water. The construction of  $Co_3O_4$  on the surface of GCE was performed as follows: 5.0  $\mu$ L of  $Co_3O_4$  solution dripped onto the surface of a freshly polished GCE and then evaporating it at room temperature in the air.

**Electrochemical measurements.** The electrochemical measurement was carried out in SWASV mode for Pb(II) detection in 0.1 M NaAc-HAc solution. A deposition potential of -1.2 V was applied for 180 s to the working electrode under stirring. The SWASV responses were recorded between -1.0 to -0.2 V with step potential of 5 mV, amplitude of 20 mV, and frequency of 25 Hz. A desorption potential of 0 V for 210 s was performed to remove the residual metals under stirring condition.

**Adsorption experiments.** Experiments were carried out at 298 K in 10 mL polyethylene centrifuge tubes containing 1 g·L<sup>-1</sup> adsorbent and various concentrations of Pb(II) for 24 h. The pH values of all these Pb(II) solutions 5.0  $\pm$  0.2. After adsorption equilibrium, the adsorbent was separated by centrifugation. The Pb(II) concentrations remaining in the solution were analyzed. The amount of metal adsorbed ( $q_e$ ) was calculated according to the following equation:

$$q_e = \frac{(C_0 - C_e)V}{m} \tag{3}$$

where  $C_0$  and  $C_e$  represent the initial and equilibrium Pb(II) concentrations (mg L<sup>-1</sup>), respectively, V is the volume of the solutions (mL), and m is the amount (mg) of adsorbent. A Langmuir isotherms model was used to analyze the experimental data. The mathematical expressions of the Langmuir isotherm is

$$\frac{C_{\rm e}}{q_{\rm e}} = \frac{1}{q_{\rm m}K_{\rm L}} + \frac{C_{\rm e}}{q_{\rm m}} \tag{4}$$

where  $q_{\rm m}$  and  $K_{\rm L}$  are Langmuir constants, representing the maximum adsorption capacity of adsorbents (mg/g) and the energy of adsorption, respectively. The values of  $q_{\rm m}$  and  $K_{\rm L}$  can be calculated from the slope and intercept of plots of  $\frac{C_{\rm e}}{q_{\rm e}}$  versus  $C_{\rm e}$ .

Computational methods. The DFT calculations were performed by using the Vienna ab-initio simulation package (VASP)<sup>42,43</sup> with the Perdew-Burke-Ernzerh parameterization of the generalized gradient approximation (GGA) adopted for the exchange correlation potential<sup>44</sup>. An energy cut of 400 eV were consistently used in our calculations. The atomic positions are fully relaxed with the conjugate gradient procedure until the residual forces vanished within 0.02 eV/Å. A  $2\times2\times1$  Monkhorst-Pack k-point was used to sample the surface Brillouin zone. Because  $\text{Co}_3\text{O}_4$  is a strongly correlated antiferromagnet and Pb ion is one of HMIs, we also performed test DFT +  $U^{45,46}$  (U = 3.3 eV and J = 0 eV according to previous reports<sup>47</sup>) calculations as well as PBE calculations with relativistic corrections<sup>48</sup> to Pb. As shown in Table S1, we can confirm that our PBE results are qualitatively correct. The transition states were obtained by using the climbing image nudged elastic band (CI-NEB) method<sup>36</sup>.

**Computational models.** The optimized lattice constant ( $a_0$ ) of the bulk  $Co_3O_4$  is 8.12 Å, which is close to the experimental value (8.08 Å)<sup>47</sup>. The bulk  $Co_3O_4$  shows antiferromagnetic properties. The  $1\times1\times10$  and  $1\times1\times11$  supercells were used to model the  $Co_3O_4$  (001) and (111) surfaces, respectively, in which a 12 Å of vacuum is adopted. During the structural optimizations, we allowed the atoms in the top four layers in the supercells to relax fully, and fixed all other atoms in the bulk configuration. The side and top views of the optimized  $Co_3O_4$  (001) and (111) surfaces are shown in **Figures S9** and **S10**, respectively. These optimized geometric structures agree well with the previous reports on the surface atoms arrangement of the  $Co_3O_4$  nanocrystals with different crystal planes  $^{10,47,49}$ .

**Adsorption energy.** To compare the Pb adsorption ability on two different crystal planes, we define the Pb adsorption energy as  $E_{\rm ads} = E_{\rm Pb/surface} - (E_{\rm surface} + E_{\rm Pb})$ . Here,  $E_{\rm surface} = E_{\rm Pb/surface} = E_{\rm Pb/$ 

**Bader charge analysis and diffusion paths.** To address the Pb-substrate interaction, we conduct the Bader charge analysis for Pb atom adsorbing on  $\text{Co}_3\text{O}_4$  (111) and (001) surfaces. The corresponding results are presented in **Figure S11**. For clarity, the diffusion directions of Pb on  $\text{Co}_3\text{O}_4$  (111) and (001) surfaces are illustrated in **Figure S12**.

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### **Author contributions**

X.J.H. conceived the electrochemical studies of nanocrystals with different facets. Q.X.L. designed the calculations. X.Y.Y. carried out all the experiments and Q.Q.M. performed calculation and contributed equally to this work. X.Y.Y. and Q.Q.M. co-wrote the paper. T.L., Y.J., B.S., and J.H.L. contributed data analysis and interpretation equally.

# **Additional information**

 ${\bf Supplementary\ information\ accompanies\ this\ paper\ at\ http://www.nature.com/scientific$ reports

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