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Role of substrate temperature on the performance of BaTiO₃/Si photodetector prepared by pulsed laser deposition

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In this study, the pulsed laser deposition (PLD) method was employed to fabricate nanostructured BaTiO₃ films on glass and silicon substrates at varying temperatures. The structural analysis confirmed the formation of crystalline nanostructured BaTiO₃ with mixed tetragonal and hexagonal phases, and the film deposited at 150 °C has the best crystallinity and largest particle size. The optical energy gap of the BaTiO₃ nanostructure decreases from 3.94 to 3.84 eV, with increasing substrate temperature from 60 to 150 °C. Photoluminescence spectra of BaTiO₃ films deposited at 25, 60, 100, and 150 °C exhibit emission peaks centered at 450, 512, 474, and 531 nm, respectively. Raman spectra of BaTiO₃ films show E (LO), A (TO), E (LO) + TO, and B1 vibration modes. Hall measurements reveal that the mobility of the BaTiO₃ film increases with temperature up to 100 °C and then decreases at 150 °C. The current–voltage characteristics of the BaTiO₃/p-Si heterojunction, deposited over a temperature range of 25 to 150 °C, were investigated in the dark and under illumination. The heterojunctions exhibit rectifying properties, with the best rectification factor observed for the heterojunction prepared at 100 °C. The values of the ideality factor for the heterojunctions fabricated at 25, 60, 100, and 150 °C were 4.3, 3.8, 2.8, and 5, respectively. The study reveals an improvement in both the figures of merit and the photodetector performance with increased substrate temperature. The responsivity increases from 2.2 to 9.25 A/W as the deposition temperature rises from 25 to 100 °C. The detectivity (D*) and external quantum efficiency (EQE) of the photodetector prepared at the optimum substrate temperature of 100 °C, were found to be 4.62×10^{12} Jones and 114%, respectively, at 500 nm.

Keywords BaTiO₃, Pulsed laser deposition, Nanostructure film, Substrate temperature, Photodetector

Barium titanate (BaTiO₃) is a significant functional semiconductor within the perovskite class, renowned for its outstanding pyroelectric, ferromagnetic, high mechanical and chemical stability, as well as electro-optic properties. These attributes render it well-suited for diverse industrial and technological uses¹. BaTiO₃ displays tetragonal symmetry, marked by the displacement of titanium and oxygen in opposite directions along the elongated axis (c-axis)². Its applications span a broad spectrum, including but not limited to piezoelectric infrared sensors, capacitors, ultrasonic transducers, dye-sensitized solar cells, ferroelectric random-access memories (FRAMs), electro-optic switches, waveguides, thermistors, and actuators^{3–5}.

The choice of synthesis technique depends on the desired properties for the final application, and the selected preparation route has a significant impact on the structure and properties of barium titanate materials^{6,7}. Various methods are employed for the synthesis of the BT system, including sol–gel^{8,9}, solid-state¹⁰, coprecipitation¹¹, chemical vapor deposition (CVD)¹², hydrothermal^{13,14}, electrophoretic deposition¹⁵, physical radio frequency sputtering¹⁶, laser molecular beam epitaxy (MBE), flash evaporation, and pulsed laser deposition (PLD)^{17–19}.

Pulsed-laser deposition (PLD) is a unique method that distinguishes itself and appears as one of the most encouraging approaches for producing thin films using multicomponent materials. Its merits encompass simplicity, cost-effectiveness, the capability for large-area film deposition, the maintenance of precise film stoichiometry, the attainment of excellent crystallinity in deposited films, and rapid processing. Noteworthy among the extensively investigated materials are high-temperature superconductors, compound semiconductors, dielectrics, ferroelectrics, electro-optic and extremely large magneto-resistance oxides, polymers, and diverse heterostructures. PLD has notably been applied in the deposition of thin films, such as BaTiO₃. Recently, Serralta-Macias et al.²⁰ reported the preparation of lead-free antiferroelectric perovskite ultrathin films (0.92(Bi_{0.54}Na_{0.46})).

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TiO₃-0.08BaTiO₃) on p-type silicon using the pulsed laser deposition technique. Silicon-based heterojunction photodetectors exhibit very attractive optoelectronic and photovoltaic properties compared to p-n silicon photodiodes. They have high responsivity, fast response, low noise, long-wavelength detection, a simple route, cost-effectiveness, and no high-temperature processing needed²¹. Several high-performance silicon-based heterojunction photodetectors have been produced using pulsed-laser deposition (PLD) methods. To the best of our knowledge, there is no reported data on the fabrication of a BaTiO₃/Si heterojunction photodetector. In this context, we fabricated a BaTiO₃/Si heterojunction photodetector using the pulsed laser deposition method, omitting the use of a buffer layer. In order to enhance the photodetector's performance, we conducted an investigation into the impact of substrate temperature on the structural, optical, and electrical properties of the BaTiO₃ film. Additionally, we evaluated the figures of merit associated with the BaTiO₃/Si photodetector.

Experimental work

A homemade pulsed laser deposition (PLD) system was utilized, comprising a Q-switched Nd:YAG laser operating at a wavelength of 532 nm (second harmonic) with a pulse duration of 7 ns, along with a glass bell jar. This system was employed to deposit BaTiO₃ thin films onto cleaned glass and silicon substrates. The BaTiO₃ pellet, measuring 1 cm² in diameter and 2 mm in depth, was created by compressing high-purity (99.99%) BaTiO₃ powder obtained from Sigma-Aldrich using a hydraulic compressor at 5 tons. The distance between the target and substrate was consistently set at 3 cm for all experiments. The BaTiO₃ films were deposited by irradiating the BaTiO₃ pellet with an energy density of 8.9 J/cm² and 200 pulses under a vacuum pressure of 10⁻⁴ mbar. The films were deposited on substrates heated to various temperatures: 25, 60, 100, and 150 °C. The glass substrate was initially cleaned with distilled water and alcohol. A polished single crystal 1 cm² p-type silicon substrate with an orientation of (111) and electrical resistivity of 3–5 Ω cm was used after being cleaned with distilled water and treatment with diluted HF. The film structure was analyzed using a Shimadzu XRD-6000 X-ray diffractometer. Atomic force microscopy (AFM), specifically the Digital Instruments Nanoscope II Scanning Probe Microscope (AFM/SPM DualScope™ DS/Germany), was employed to investigate the surface morphology and particle size distribution of the deposited films. A field emission scanning electron microscope (FESEM) equipped with energy-dispersive X-ray (EDX) from Imaging-EDS-Mapping/Germany was used to examine the chemical composition of the BaTiO₃ thin film. For optical properties measurements, a double-beam UV-Vis spectrophotometer (UV-DRS Shimadzu UV-2550) was used. Raman shift analysis was conducted with a Raman Spectroscopy system (Raman Malvern HORIBA XploRA PLUS/UK). Photoluminescence measurements were carried out at room temperature using an argon ion laser with a wavelength of 320 nm (PL Varian CARY ECLIPSE) as the excitation source. Hall effect measurements were performed to determine the electrical conductivity and mobility of the deposited films. Ohmic contacts were established by depositing In film on the BaTiO₃ film and Al electrode on the backside of the silicon substrate through a square thin metal mask using a thermal evaporation system. Figure 1a illustrates the schematic diagram of the BaTiO₃/Si heterojunction photodetector. The film thickness deposited on silicon substrates ranged from 600 to 1400 nm, depending on the substrate temperature. The current–voltage characteristics of the BaTiO₃/Si heterojunction were measured under dark and white illumination conditions. The responsivity of the BaTiO₃/Si photodetector was measured in the spectral range of 400–900 nm using a Jobin Yvon monochromator, beam splitter, and halogen lamp under reverse bias voltage. The power of the light at a certain wavelength was calibrated with a silicon power meter. Figure 1b shows the schematic diagram of the setup used to measure the dark, illuminated I–V characteristics, and spectral responsivity.

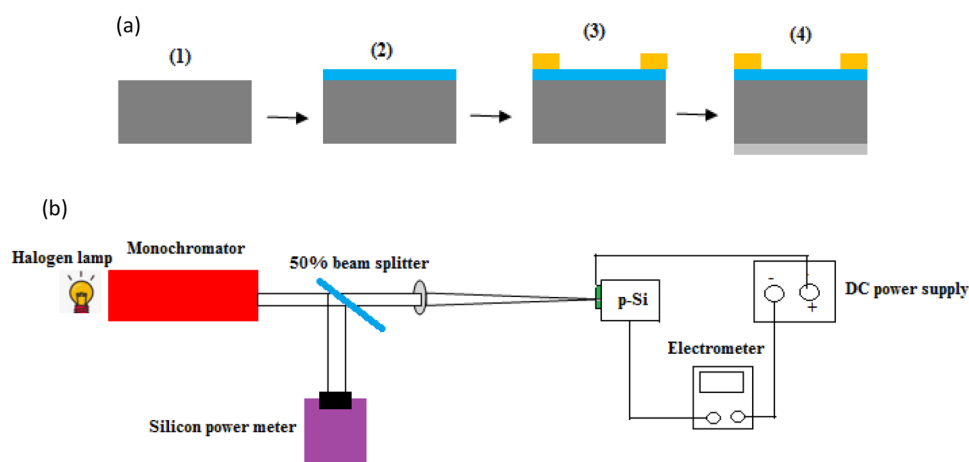


Figure 1. (a) Fabrication steps of BaTiO₃/Si photodetector (1) Cleaned silicon substrate, (2) deposition of BaTiO₃ film, (3) evaporation of In- electrode on BaTiO₃ film, and (4) Evaporation of Al-electrode on backside of Si and (b) Schematic diagram of dark, illuminated I–V characteristics and responsivity measurements setup.

Results and discussion

The XRD patterns of the BaTiO₃ pellet and films deposited at 25, 60, 100, and 150 °C are illustrated in Fig. 2. The film deposited at 25 °C exhibits seven peaks that confirm the formation of a hexagonal-tetragonal co-existing BaTiO₃ film. These peaks are observed at $2\theta = 22^\circ, 31.4^\circ, 38.9^\circ, 40.52^\circ, 45.1^\circ, 51.1^\circ,$ and 56° , corresponding to the (100), (110), (006), (114), (200), (210), and (211) reflection planes, respectively. According to JCPDs # 00-005-0626²² and 00-034-0129²³, the BaTiO₃ film deposited at room temperature (RT) was crystalline and exhibited seven diffraction peaks corresponding to the (100), (110), (006), (114), (200), (210), and (211) reflection planes, respectively. These peaks along (100), (110), (200), (210), and (211) belong to crystalline BaTiO₃ with a tetragonal phase, while the peaks along (006) and (114) planes indeed belong to hexagonal BaTiO₃. A small shift in diffraction angles of the XRD peaks of the nanostructured films compared to the bulk is due to the presence of stacking fault^{24,25}.

The lattice constant of tetragonal BaTiO₃ was calculated and found to be $a = b = 0.4032$ nm, and the c/a ratio was 0.997 nm, which is close to the lattice constant for bulk tetragonal BaTiO₃ ($a = b = 0.3994$ nm) with a c/a ratio of 1.011 nm, according to JCPDs # 00-005-0626. The lattice constants of hexagonal BaTiO₃ were $a = b = 0.57116$ nm, and $c = 1.3931$ nm. The ratio of c/a is 2.4391. These lattice parameters and the c/a ratio are in good agreement with those for bulk BaTiO₃: $a = b = 0.57248$ nm, $c = 1.39673$ nm, and c/a ratio of 2.4397, according to JCPDs # 00-034-0129. Increasing the deposition temperature results in the presence of a new peak found at 32.8° corresponding to the (112) plane, as well as a slight increase in the intensity of the other XRD peaks. This indicates that the crystallinity of the nanostructured film improved after increasing the substrate temperature. This improvement could be attributed to the fact that the kinetic energy of BaTiO₃ nanoparticles increases as the substrate temperature rises, giving them the ability to move very quickly across the film surface and create low surface energy structures²⁶. No XRD peaks related to other elements or contaminants have been found in the XRD patterns, indicating the purity of the deposited nanostructured films. The average crystallite size of the BaTiO₃ film was estimated using the Debye–Scherrer formula:

$$D = \frac{0.9\lambda}{\beta \cos\theta} \quad (1)$$

where λ is the Cu K α X-ray wavelength ($\lambda = 0.154$ nm), θ is the diffraction angle, and β is the full width at half maximum FWHM of the XRD peak. Table 1 lists the values of the average crystallite size of the films deposited at various substrate temperatures.

The crystallite size increased when the deposition temperature increased as a result of the enhancement of film crystallinity. The strain ϵ and dislocation density δ formed in the film were calculated using the following equations:

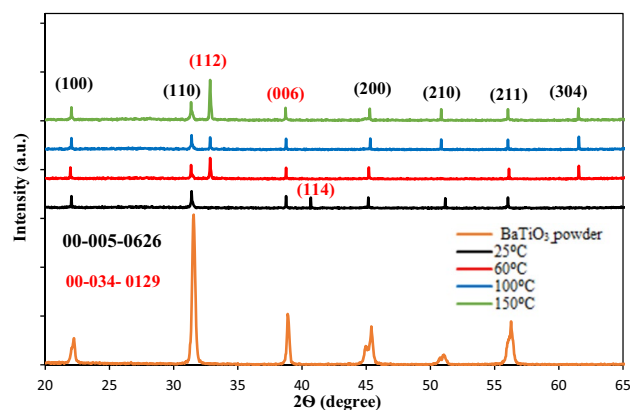


Figure 2. X-ray diffraction patterns of BaTiO₃ powder and BaTiO₃ nanostructure deposited at various substrate temperatures.

Substrate temperature (°C)	Lattice constants for tetragonal structure (nm)			Lattice constants for hexagonal structure (nm)			Crystallite size (D) (nm)	Strain (ϵ)	Dislocation density (nm^{-2})
	a = b	c	c/a	a = b	c	c/a			
25	0.403	0.402	0.997	0.571	1.3931	2.409	28.6	3.05×10^{-3}	1.22×10^{-3}
60	0.404	0.401	0.992	0.578	1.3934	2.406	42.8	2.44×10^{-3}	5.46×10^{-4}
100	0.403	0.401	0.992	0.578	1.3937	2.407	45.05	3.05×10^{-3}	4.83×10^{-4}
150	0.403	0.403	0.999	0.578	1.395	2.4107	55.6	3.67×10^{-3}	3.23×10^{-4}

Table 1. Effect of substrate temperature on crystallite size, strain, and dislocation of the film.

$$\varepsilon = \frac{\beta \cos \theta}{4} \quad (2)$$

$$\delta = \frac{1}{D^2} \quad (3)$$

It can be seen from Table 1 that the strain and dislocation density in the film reduce as the deposition temperature rises due to the decreasing density of the structural defects accompanied by the Fig. 3 shows the FESEM images of the nanostructured BaTiO₃ films deposited at various substrate temperatures²⁷. The film deposited at 25 °C, as shown in Fig. 3a, exhibits the formation of randomly distributed high density of nanoparticles with different sizes, with most grains having a spherical shape. Some larger agglomerated nanoparticles are observed to be formed on the surface, representing the last deposited grains. The average grain size was estimated using Image J software and found to be 16 nm. The film deposited at 60 °C (Fig. 3b) also shows the formation of a high density of nanoparticles, but with a small number of agglomerated particles, and the average particle size was 23 nm. Films deposited at substrate temperatures of 100 °C, as shown in Fig. 3c, exhibit larger grains with an average size of 44 nm, as well as the formation of particulates and some white color droplets due to the laser splashing effect. The film deposited at 150 °C was compact, dense, thicker, and have droplets with an average grain size of 91 nm, as shown in Fig. 3d. Increasing the substrate temperature leads to an increase in the amount of material ablated and deposited on the substrate. Furthermore, the adhesion of the nanostructured film was checked by scratching technique and found to be improved compared to that deposited at room temperature. No cracks or voids have been observed on the film surface. Figure 4 shows the cross section FESEM image of the BaTiO₃ film deposited on a silicon substrate at substrate temperatures of 25 and 100 °C. The inset of Fig. 4 is the particle size distribution of the samples deposited at various temperatures. The samples have a narrow and nearly Gaussian distribution.

The images confirm the formation of spherical nanoparticles and a clear boundary between the film and silicon substrate. The nanostructured film thickness was measured and found to increase with substrate temperature,

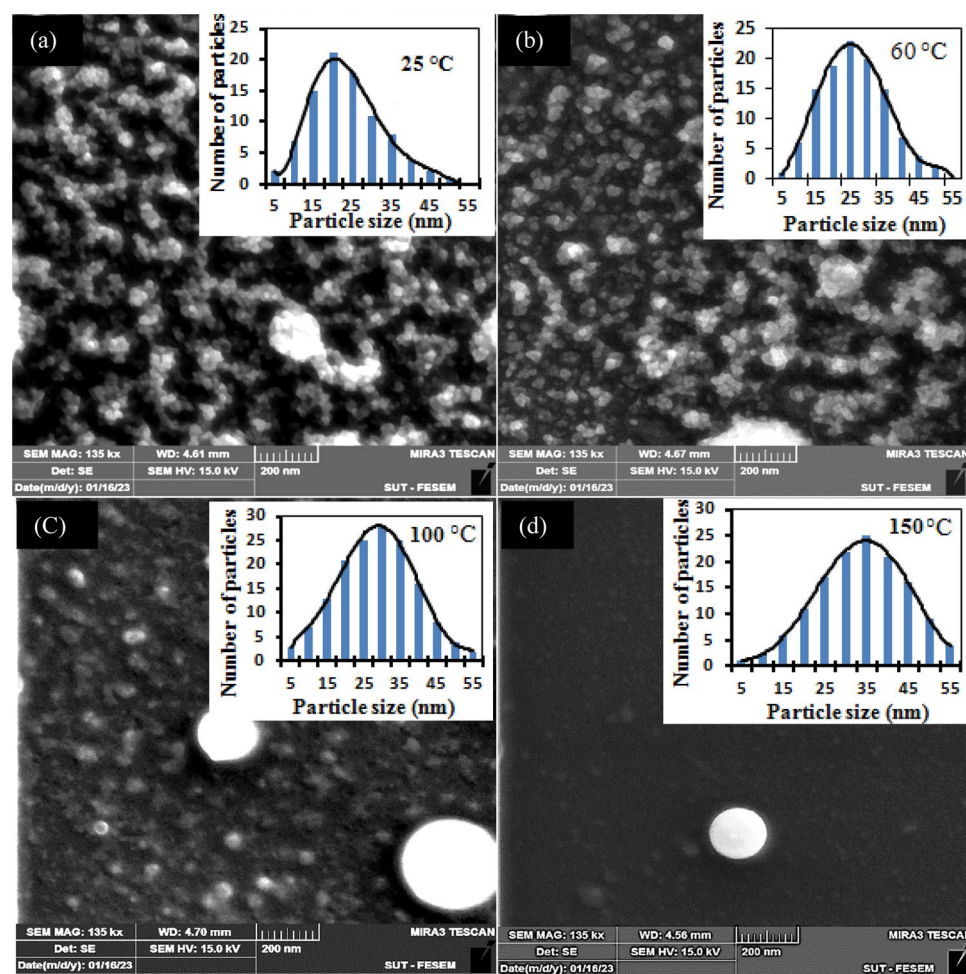


Figure 3. FESEM images of BaTiO₃ nanoparticles deposited at substrate temperature of (a) 25, (b) 60, (c) 100, and (d) 150 °C. Inset are the particle size distribution of the samples.

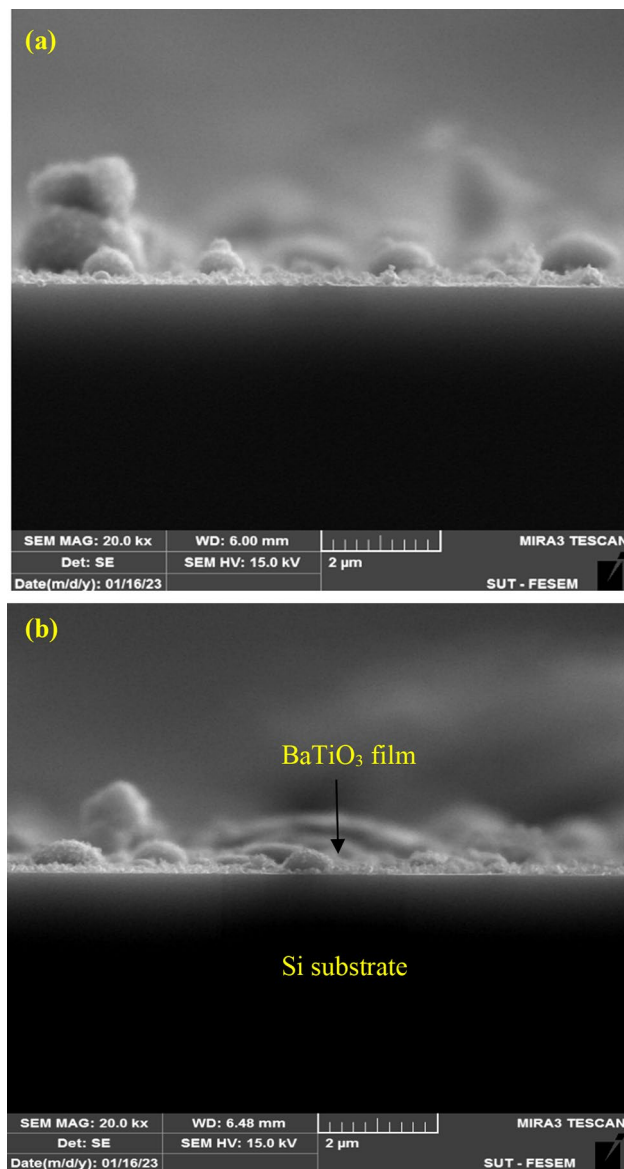


Figure 4. Cross section FESEM images of BaTiO₃ nanostructures deposited at (a) 25 and (b) 100 °C.

as shown in Fig. 5. This increase could be due to the increasing ablated and deposited grains with an increase in substrate temperature. As shown in Fig. 5, the thickness of the nanostructured film demonstrated a linear increase with substrate temperature. The calculated increment in film thickness per degree Celsius of substrate temperature was found to be 6.4 nm/°C.

Figure 6 shows the EDX spectra of the BaTiO₃ film deposited at various substrate temperatures. All spectra confirm the presence of peaks related to the Ba, Ti, and O elements, which are the main elements of the BaTiO₃ film. The Au peaks come from the deposition of the film on the BaTiO₃ film for SEM investigation. The weight percentages of the elements are shown in Table 2. The best stoichiometry was found for the film deposited at 60 °C. No other elements were detected in the EDX spectra. Decreasing the amount of barium with the rise in substrate temperature, due to the thermal decomposition of barium as well as the vaporization of the barium component more than the other components, could lead to a barium-deficient film.

Figure 7 illustrates the EDX mapping images of the films deposited at various substrate temperatures, revealing that the Ba, Ti, and O elements are distributed across the film surface in all samples. The film deposited at 60 °C shows more grain agglomeration compared to the other samples.

Figure 8 shows 3D AFM images of the BaTiO₃ films. The film topography depended strongly on the substrate temperature. The film deposited at room temperature (Fig. 8a) shows the presence of rounded grains, as depicted in Fig. 8b, distributed over the entire surface.

The AFM image of the film deposited at $T_s = 60$ °C reveals the formation of triangular, uniformly distributed grains with different heights, and the root mean square (RMS) of surface roughness is higher than that of the film deposited at room temperature, as shown in Table 3. Increasing the substrate temperature to 100 °C leads to

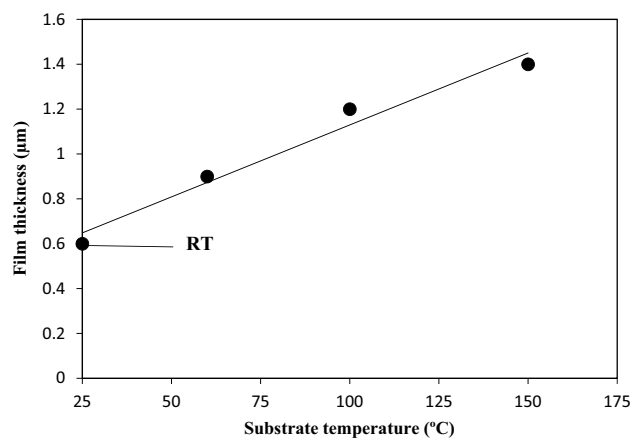


Figure 5. Variation of film thickness with substrate temperature.

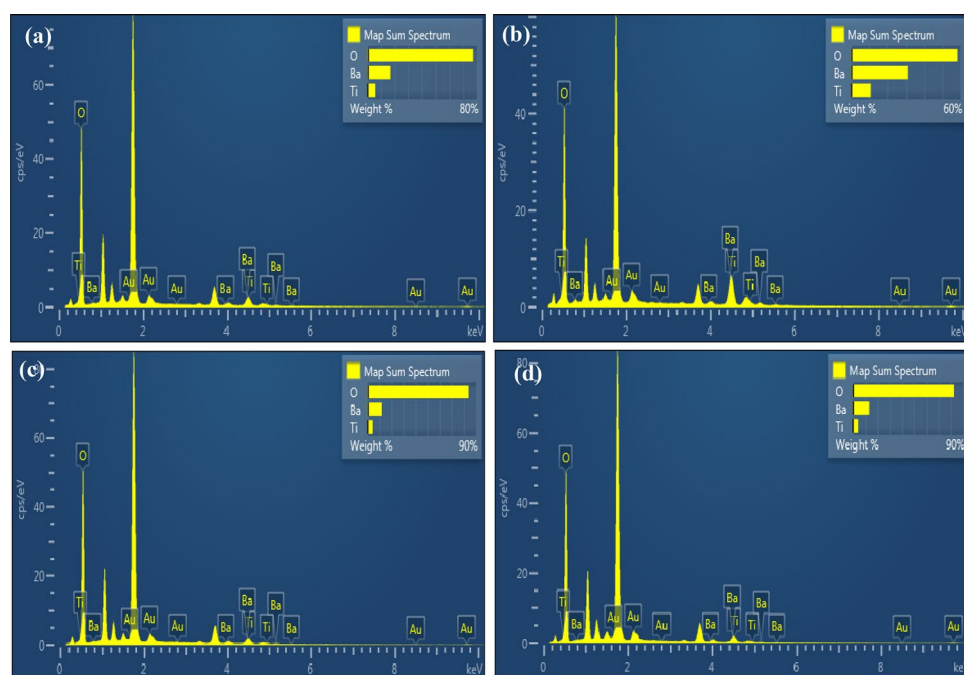


Figure 6. EDX spectra of BaTiO₃ nanostructure films deposited at (a) 25, (b) 60, (c) 100, and (d) 150 °C.

Element	Weight%			
	25°C	60°C	100°C	150°C
O ₂	77.2	58.2	83	82.5
Ti	5.8	10.6	4.3	4.5
Ba	16	30.8	11.8	13

Table 2. Weight percentage of BaTiO₃ film as a function of substrate temperature.

the formation of a small amount of large triangular grains and irregular grains on the film surface. Finally, the AFM image of the film deposited at $T_s = 150$ °C shows large irregular and rounded grains with some triangular grains. The RMS of surface roughness increases with substrate temperature, as shown in Table 3, due to the high kinetic energy and mobility of the deposited adsorbed grains on the substrate^{22,28}. The grain size of the film increases with substrate temperature, which is in good agreement with XRD data.

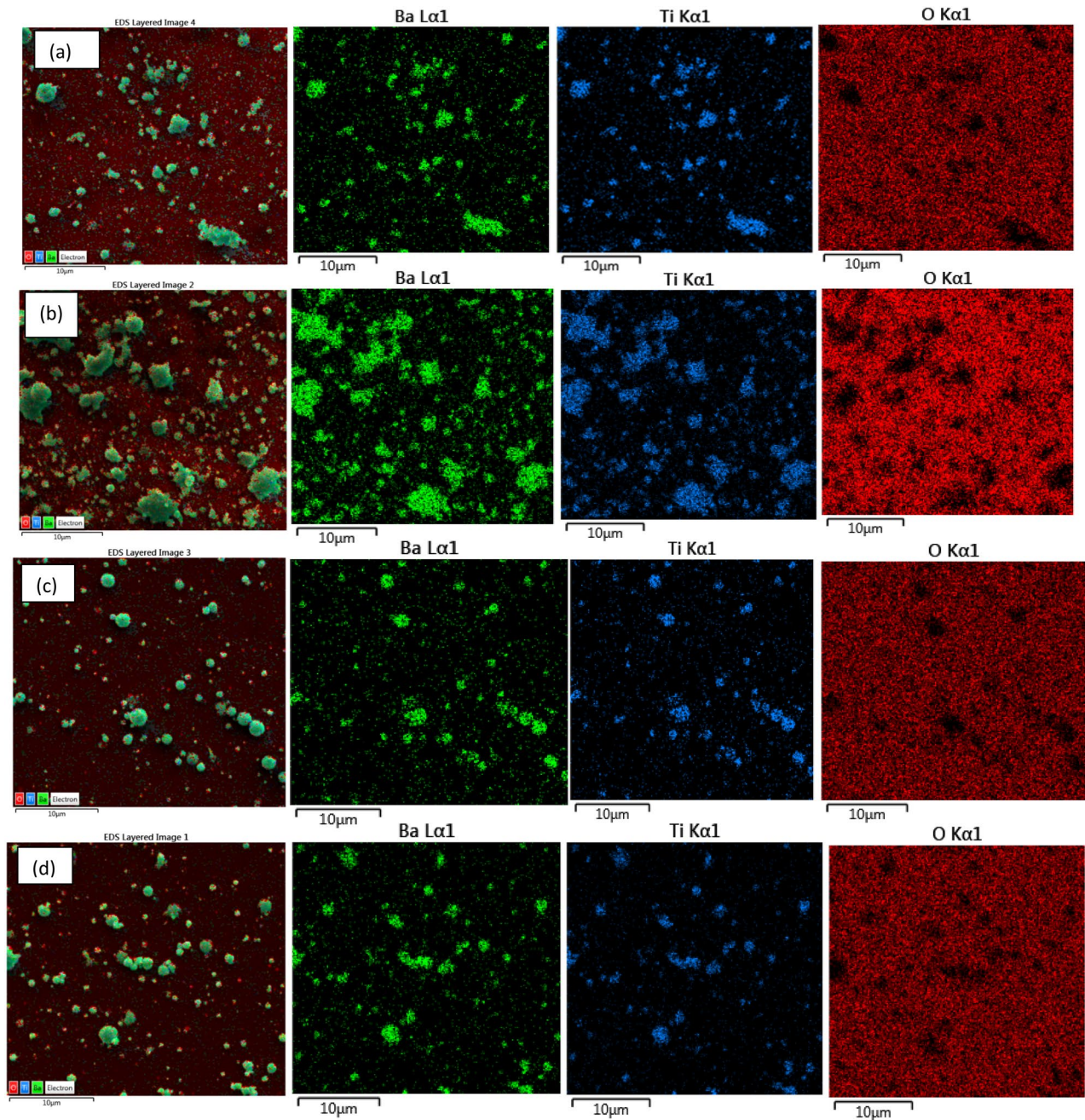


Figure 7. Elemental EDX mapping images of BaTiO₃ films deposited at substrate temperature of at (a) 25 °C (b) 60 °C (c) 100 °C, and (d) 150 °C.

Figure 9a shows the effect of substrate temperature on the optical transmission of the BaTiO₃ film. The optical transmittance of the film decreases as the substrate temperature increases due to the increased thickness and RMS surface roughness of the film with increasing substrate temperature. The average optical transmission of the films deposited at Ts = 25, 60, 100, and 150 °C was 85%, 78%, 76%, and 59.5%, respectively. The transmission of the film deposited at 150 °C has decreased remarkably due to the increase in surface roughness and the formation of droplets and particulates^{29,30}, which increases the scattering losses, and these agree well with AFM and SEM results.

The optical transmission of the films increases significantly after 352 nm and then tends to saturate, with a small blue shift detected for the sample deposited at Ts = 100 °C. The optical energy gap (E_g) of the BaTiO₃ film was calculated using the following Tauc's formula:

$$(\alpha h\nu)^2 = A(h\nu - E_g) \quad (4)$$

where A is a constant depending on the material, α is the absorption coefficient, $h\nu$ is the photon energy. The value of E_g can be obtained by plotting $(\alpha h\nu)^2$ against $h\nu$ and extrapolating the linear portion of the curve to the

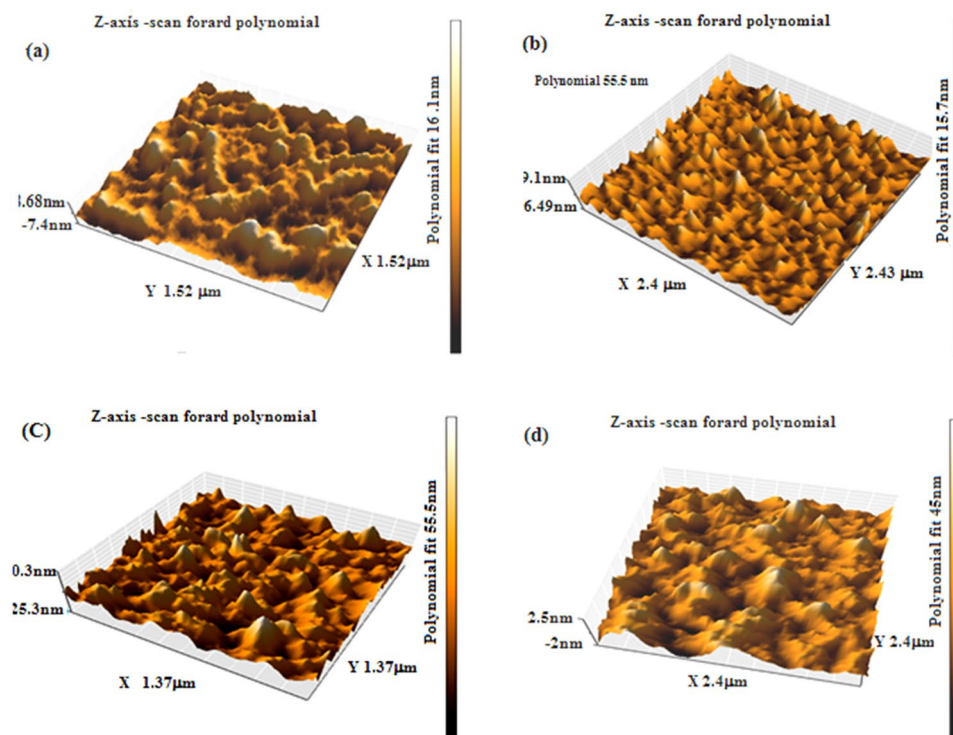


Figure 8. 3D AFM images of BaTiO₃ films deposited at (a) 25, (b) 60, (c) 100 °C, and (d) 150 °C.

Ts (°C)	RMS of roughness (nm)	Grain size (nm)
25	11.17	36
60	19.91	50
100	22.37	68
150	30.50	74

Table 3. RMS of surface roughness of the film and grain size obtained from AFM analysis.

photon energy axis, as shown in Fig. 9b. The optical band gap of the BaTiO₃ films deposited at 25, 60, 100, and 150 °C was 3.94 eV, 3.91 eV, 3.87 eV, and 3.84 eV, respectively.

The obtained values of the energy gap are larger than that of bulk BaTiO₃ (3.2 eV) due to the quantum size effect, and they are also in good agreement with reported data^{28,30,31}. This means that the energy gap of the film decreases as the substrate temperature increases, as shown in Fig. 10. We can attribute the decrease in the optical energy gap with substrate temperature to an increase in the grain size of the film with increasing Ts, as previously shown in the SEM results.

Raman spectra of the BaTiO₃ film deposited at different substrate temperatures are depicted in Fig. 11. As shown, the Raman band at 220 cm⁻¹ is indeed the E (LO) (longitudinal optic) vibration mode. Raman peaks located at 260, 288, and 480 cm⁻¹ are related to the A (TO) vibration mode, and the peaks observed at 414, 438 cm⁻¹ belong to the E (TO)/E(OL) mode. The peak at 343 cm⁻¹ is assigned to E (LO) + TO, B1) mode. The broad peaks give an indication of the polycrystalline nature of the deposited films and the degree of crystallinity. As reported, these obtained vibration modes arise from the displacement of Ti⁴⁺ ions brought on due to the effect of heating the substrate at high temperatures^{32–35}. The full width at half maximum (FWHM) of the Raman bands was found to increase with increasing substrate temperature, and this could be correlated to the film's crystallinity quality.

Figure 12 illustrates the photoluminescence (PL) spectra of BaTiO₃ films deposited on glass substrates at varying substrate temperatures, utilizing an excitation source with a 320 nm wavelength. At a deposition temperature of 25 °C, the film displays four broad emission peaks at 365 nm (3.39 eV), 450 nm (2.75 eV), 525 nm (2.36), and 660 nm (1.87 eV). The first emission peak is due to band-to-band transition and the third and fourth peaks are originated from magnetic-dipole transition³⁶. Conversely, films deposited at substrate temperatures of 60 °C, 100 °C, and 150 °C exhibit prominent emission peaks at 512 nm (2.42 eV), 474 nm (2.61 eV), and 531 nm (2.33 eV), respectively. The other emission peaks can be attributed to the structural defects, deep level emission, and traps in the energy level system of the film as well as due to the small off-stoichiometry of the deposited films. The origin of blue emission is linked to delocalized electronic levels proximate to the valence

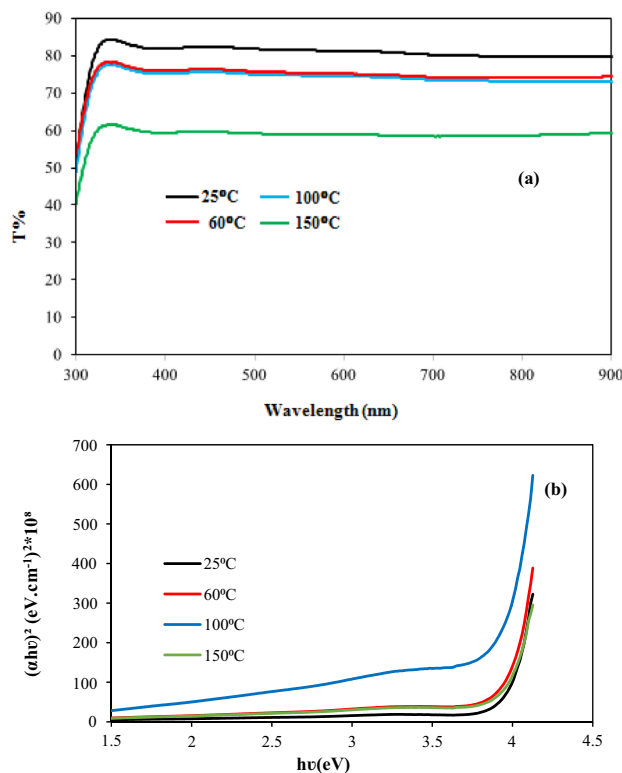


Figure 9. (a) Optical transmission plot and (b) $(\alpha h\nu)^2$ against photon energy of BaTiO₃ films deposited at various substrate temperatures.

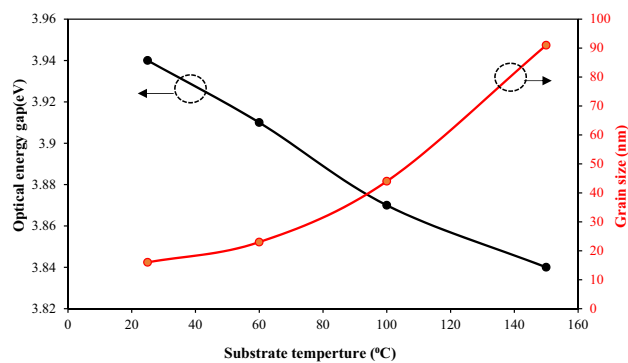


Figure 10. Dependence of optical energy gap of film on substrate temperature.

and conduction bands, while the green emission results from the radiative recombination of an electron in the deep oxygen vacancy energy band with a hole in the valence band^{37,38}. The intensity of PL emission peaks increases with substrate temperature due to the improvement in crystallinity of the films and a reduction in the optical band gap are responsible for the increase in visible emission intensity that is seen between 25 and 150 °C.

The electrical characteristics of the BaTiO₃ films, including carrier concentration, mobility, electrical resistivity, and conduction type, were examined by Hall effect. The results reveal that the Hall coefficient of the films was negative, indicating that the films are n-type. The variation of the mobility and carrier concentration of the film with substrate temperature is shown in Fig. 13. The mobility of the film increased from 144 to 171 cm²/V·s as the substrate temperature rises from 25 to 100 °C as a result of increasing grain size and decreasing grain boundaries and decreased after this temperature due to the formation of droplets and particulates. The carrier concentration of the electrons decreases as the substrate temperature increases due to increasing mobility. Figure 14 illustrates that the electrical resistivity of the film decreases when the substrate temperature rises from 25 to 150 °C; it decreases from 1.24×10^5 to $0.0169 \times 10^5 \Omega \text{ cm}$ due to the increasing grain size and electron mobility of the film^{39–42}.

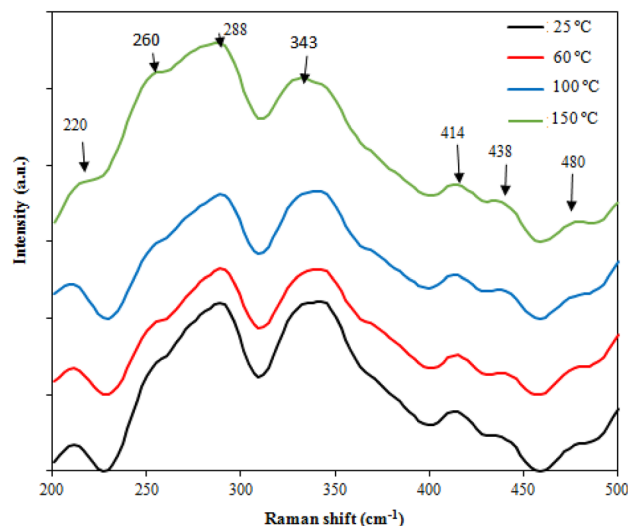


Figure 11. Raman spectra of BaTiO₃ films prepared at different deposition temperatures.

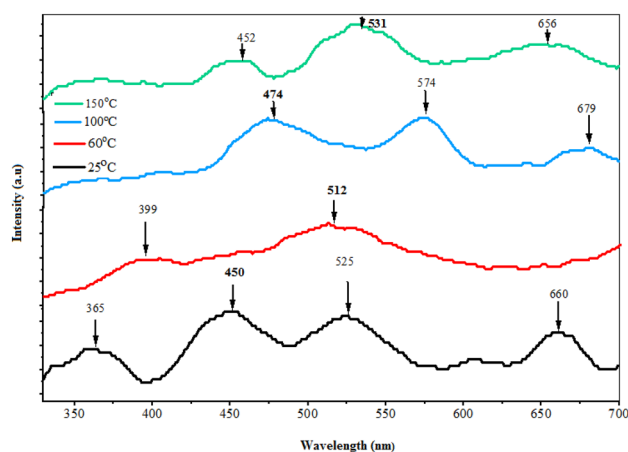


Figure 12. Room temperature PL spectra of BaTiO₃ films deposited at different substrate temperatures.

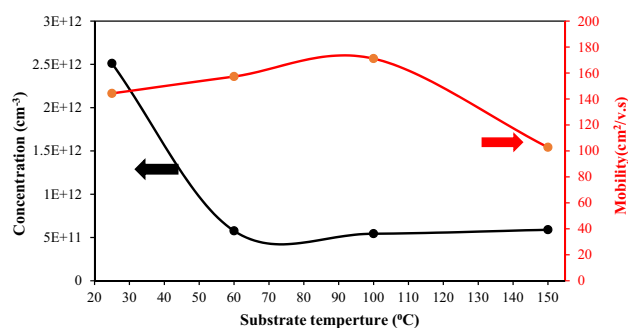


Figure 13. Carrier concentration and mobility of BaTiO₃ film against substrate temperature.

Figure 15 shows the dark I–V characteristics of the n-BaTiO₃/p-Si heterojunction measured over the voltage range of – 5 to + 5 V. All fabricated heterojunctions demonstrate rectification properties, and the rectification factor was enhanced after increasing the substrate temperature to 100 °C due to the decreasing surface leakage current and increasing forward current as a result of reducing the electrical resistivity of the film. The forward current increases exponentially with bias voltage, and the substrate temperature affects the turn-on voltage

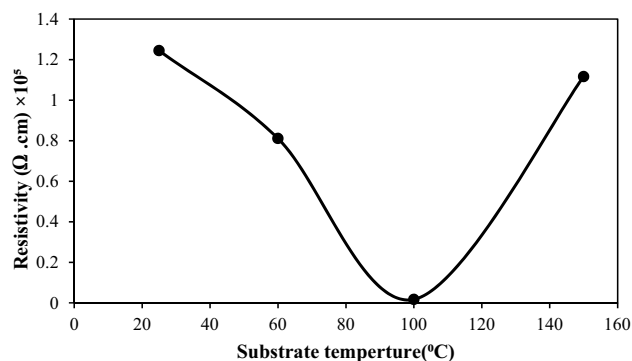


Figure 14. Electrical resistivity as a function of substrate temperature.

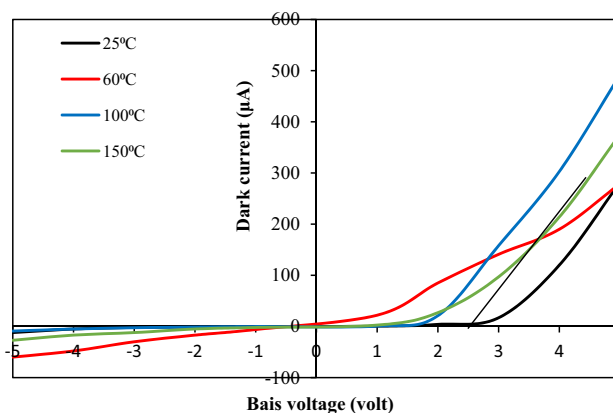


Figure 15. Dark I–V characteristics of BaTiO₃/Si heterojunctions deposited at various substrate temperatures.

(V_T), ranging from 0.8 to 2.75 V depending on the electrical resistivity of the film. The forward currents of the heterojunctions deposited at 25, 100, and 150 °C show the domination of diffusion current, particularly at bias voltages larger than the turn-on voltage, while the heterojunction deposited at $T_s = 60$ °C illustrates a forward current suffering from recombination processes due to structural defects and high series resistance.

The ideality factor (n) of the n-BaTiO₃/p-Si heterojunction was calculated by using the following diode equation:

$$n = \frac{q}{k_B T} \frac{\Delta V}{\ln \frac{\Delta I}{I_s}} \quad (5)$$

where q is the electron charge, k_B is the Boltzmann constant, and I_s is the saturation current of the heterojunction. The value of I_s was estimated by extrapolating the linear part of the semilogarithmic forward current–voltage characteristics to the current axis. The values of the ideality factor, turn-on voltage, and saturation current density as a function of substrate temperature are shown in Table 4. As shown in Table 4, the ideality factor is larger than unity, and the best value has been obtained for the heterojunction deposited at 100 °C, indicating that the BaTiO₃-Si interface has a low density of surface states and traps^{43–47}.

Figure 16 illustrates the I–V characteristics of the BaTiO₃/Si heterojunction when the photodetector is illuminated with white light of varying intensities (100, 130, and 150 mW/cm²). The generation of electron–hole (e–h) pairs causes the photodetector’s current to increase when illuminated with light due to the absorption of light

T_s (°C)	V_T (V)	J_s (nA/cm ²)	n
25	1.2	570	4.3
60	3	450	3.8
100	1.8	150	2.8
150	2.5	630	5

Table 4. Effect of substrate temperature on V_T , J_s , and n of the BaTiO₃/Si heterojunction.

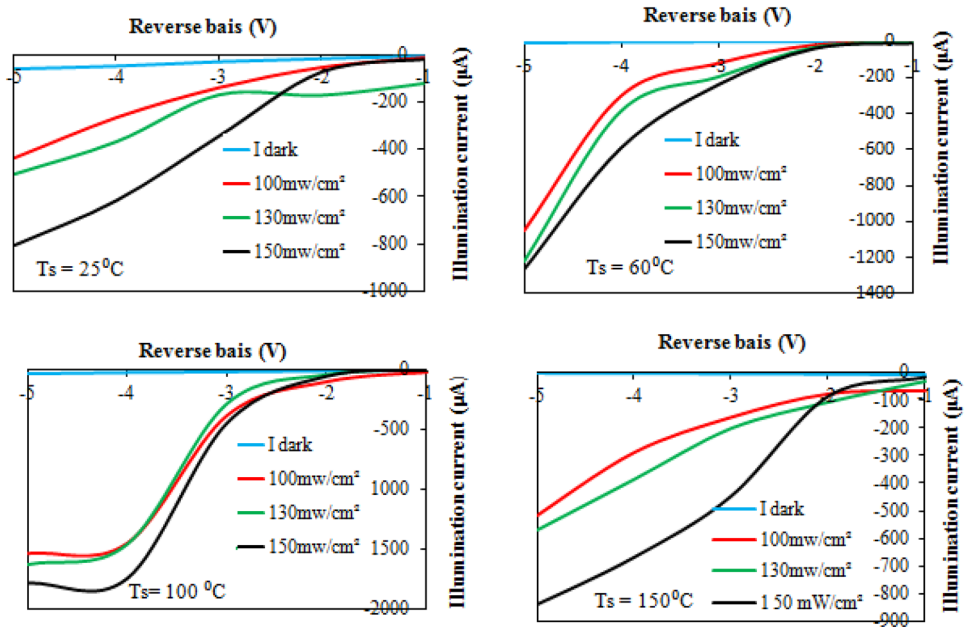


Figure 16. Illuminated I-V characteristics of n-BaTiO₃/p-Si photodetectors.

in the depletion region of the heterojunction. The increasing number of photogenerated e-h pairs as a result of increasing the light intensity leads to an increase in photocurrent as light intensity increases. No saturation in the photocurrent was detected after increasing the light intensity, indicating that the photodetectors have linearity properties. As depicted, raising the substrate temperature to 60 and 100 °C results in an increase in photocurrent due to enhanced film mobility and widening of the depletion region. However, an increase in substrate temperature to 150 °C leads to a decrease in photocurrent, attributed to nanoparticle agglomeration, the presence of droplets, and particulates. The sensitivity of the photodetectors as a function of substrate temperature is shown in Fig. 17, and the maximum sensitivity was around 180 for the photodetector deposited at Ts = 100 °C. We can attribute this result to the improvement of the junction characteristics with minimal structural and interface defects, as well as an increase in the diffusion length (L_D) of the minority carriers. The diffusion length increases as the mobility (μ) increases according to the following relationship:

$$L_D = \left(\frac{KT}{q} \mu \tau\right)^{0.5} \tag{6}$$

where τ is the minority carrier lifetime of the photodetector and T is the operating temperature.

The spectral responsivity of BaTiO₃/p-Si heterojunction photodetectors fabricated at different substrate temperatures is shown in Fig. 18. The photodetector’s responsivity R_λ was calculated from the following equation:

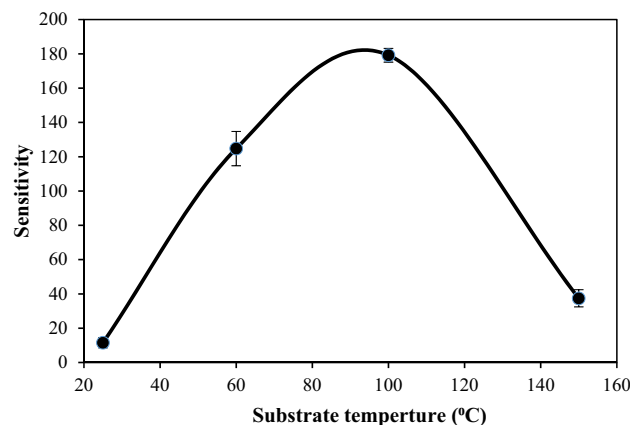


Figure 17. Impact of substrate temperature on sensitivity of BaTiO₃/p-Si photodetector.

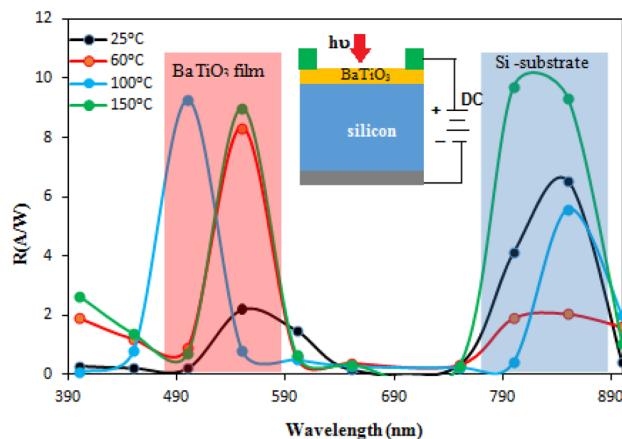


Figure 18. Spectral responsivity of BaTiO₃/p-Si heterojunction photodetectors prepared at various substrate temperatures at bias voltage of -5V. Inset is schematic diagram of the photodetector under reverse bias with electrical configuration.

$$\mathfrak{R} = \frac{I_{ph}}{P} \quad (7)$$

where P is the light power and I_{ph} is the photocurrent. We have observed two response peaks for all fabricated photodetectors. The photodetector deposited at room temperature exhibits two peaks located at 550 nm and 850 nm with responsivity values of 2.2 and 6.5 A/W, respectively. The first peak belongs to light absorbed in the depletion region located in the BaTiO₃ (absorption edge of BaTiO₃), and the origin of the second peak is the absorption edge of the silicon substrate^{48–50}. Considering the absorption edge of nanostructured BaTiO₃, the peak response should be around 350 nm rather than 550 nm. The absorption coefficient at 350 nm for the film is significantly higher compared to the wavelength of 550 nm, and the corresponding absorption depth is very small. This suggests that the 350 nm wavelength will be absorbed on the film surface, and the photogenerated carriers can rapidly recombine due to the presence of surface states and dangling bonds on the surface of the nanoparticles. Increasing the substrate temperature results in a significant enhancement of responsivity, accompanied by an observed shift in peak response. The responsivity of the photodetector increased from 2.2 A/W at 550 nm to 8.3, 9, and 9.2 A/W after increasing the substrate temperature to 60, 100, and 150 °C, respectively. This improvement in responsivity after increasing the substrate temperature is attributed to the enhancements in the structural, optical, and electrical properties of the film^{51,52}.

The shift in peak of response after increasing the temperature to 100 °C can be attributed to the increase in the depletion layer thickness moving toward the BaTiO₃ film and quantum size effect⁵³, thereby enhancing the detection of short wavelengths. The presence of bias voltage plays a crucial role in increasing the responsivity by preventing e–h recombination and widening the depletion region width^{54–56}.

The effect of the substrate temperature on the specific detectivity of the photodetector is demonstrated in Fig. 19a. The specific detectivity (D^*) of the photodetector is defined as the minimum detectable power and it can be estimated by the following formula:

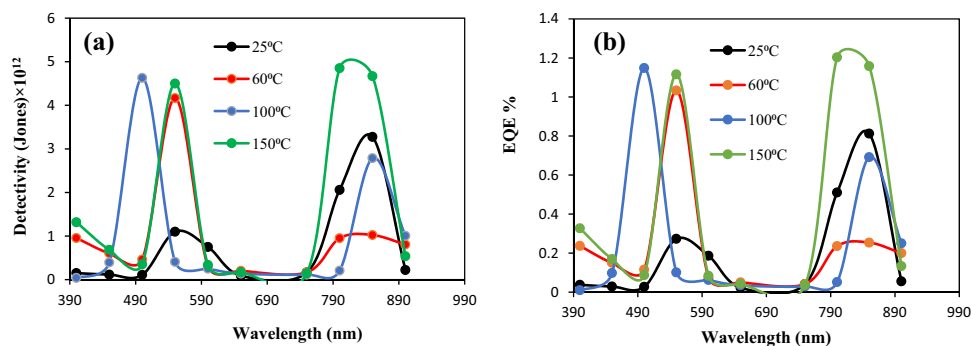


Figure 19. (a) D^* and (b) EQE of n-BaTiO₃/p-Si heterojunction photodetectors prepared at various substrate temperatures.

$$D^* = R_\lambda (A/2qI_d)^{\frac{1}{2}} \quad (8)$$

here, A is the sensitive area of the photodetector, and I_d is the dark current. The highest value of D^* was 4.62×10^{12} Jones at 500 nm for the photodetector prepared at $T_s = 100^\circ\text{C}$ and corresponded to a noise equivalent power (NEP) of 0.21 PW. Figure 19b illustrates the external quantum efficiency (EQE) of the photodetector deposited at different substrate temperatures. As shown in Fig. 19b, the photodetector fabricated at 100°C has EQE exceeding unity, and this result is due to light trapping, minimal structural defects, minimal grain boundaries, widening the depletion layer thickness, low density of surface states, and recombination^{57–59}. Table 5 provides a comparison of the figures of merit for the BaTiO₃/Si photodetector deposited at 100°C with those of heterojunctions based on silicon photodetectors reported in the literature.

To understand the mechanism of light detection in BaTiO₃/Si photodetector, the energy band diagram was constructed, as shown in Fig. 20. The conduction band offset between the BaTiO₃ nanostructure and silicon substrate was estimated as $\Delta E_C = \chi_{\text{BT}} - \chi_{\text{Si}}$ ($4.8 - 4.05 = 0.75$ eV, and the valence band offset as $\Delta E_V = (E_{\text{gBT}} - E_{\text{gSi}}) - (\chi_{\text{BT}} - \chi_{\text{cSi}}) = (3.87 - 1.12) - 0.75 = 2$ eV. Upon the light incident on the BaTiO₃ film, it will be absorbed in the depletion layer region of the n-BaTiO₃, which leads to the generation of an e–h pair and an consequently increased the photocurrent of the photodetector. The presence of the internal electric field as a result of intimate contact between BaTiO₃ film and silicon substrate prevents e–h pair recombination^{62,63}. The light exciting the electron from valence band to conduction band produces a free electron in conduction and a free hole in valence band, and these carriers conduct current in photodetector.

Conclusions

In this study, high-performance n-BaTiO₃/p-Si heterojunction photodetector using the laser deposition route was successfully fabricated without using buffer layer. The structural, optical, and electrical properties of BaTiO₃ was studied as a function of substrate temperature. X-ray diffraction studies confirmed that nanostructured BaTiO₃ was polycrystalline with a mixture of tetragonal and hexagonal phases. The particle size increases with an increase in the substrate temperature. The optical energy gap of the film decreased from 3.94 to 3.84 eV as T_s increases from 25 to 100°C due to an increase in the grain size. AFM investigation confirms that the surface roughness of the film increases as the substrate temperature increases, enhancing the light trapping of the film. Based on the obtained results, the optimal substrate temperature was found to be 100°C . Examining the impact of substrate temperature on the BaTiO₃/Si photodetector's figures of merit, we observed a superior device at 100°C with a responsivity of 9.2 A/W, a specific detectivity of 4.62×10^{12} Jones, and an external quantum efficiency of 114%. Furthermore, we constructed the illuminated energy band diagram for the n-BaTiO₃/p-Si heterojunction photodetector to understand the photodetection mechanism.

Collectively, our results emphasize the potential of these photodetector for applications requiring the detection of very weak light signals, affirming its effectiveness in specialized and demanding environments.

Photodetector type	Preparation method	Peak wavelength (nm)	Responsivity R_λ (A/W)	External quantum efficiency EQE %	Detectivity D^* (Jones)	References
Bi ₂ O ₃ /Si	Laser ablation in liquid	610	0.23	0.5×10^2	1.5×10^{11}	50
PbI ₂ /Si	pulsed laser deposition	610	0.4	0.5×10^2	20.4×10^{12}	51
ZnO/Si	Laser ablation	450	9.9	2.951×10^2	4.3×10^{12}	60
HgI ₂ /Si	pulsed laser deposition	400	1.09	3×10^2	3.6×10^{12}	61
This work	pulsed laser deposition	500	9.2	1.14×10^2	4.62×10^{12}	

Table 5. Figures of merit of some heterojunction-based silicon photodetectors compared with fabricated BaTiO₃/Si at 100°C (best photodetector).

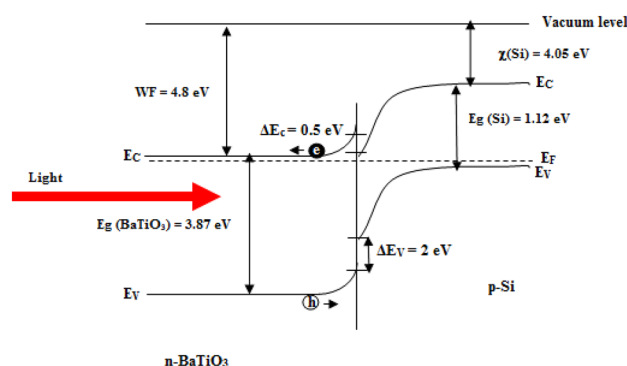


Figure 20. Illuminated energy band diagram of BaTiO₃/Si heterojunction photodetector prepared at 100°C .

Data availability

The datasets generated during and/or analyzed during the current study are available from the corresponding author (R.A. Ismail)) on reasonable request.

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Author contributions

R.A.I. and S.S.S. conceived of the presented idea. R.A.I. and S.S.S. supervised the finding of this work. All authors discussed the results. N.Z.A. conducted the experiments. R.A.I. and N.Z.A. provided critical feedback and helped shape the research, analysis and manuscript.

Competing interests

The authors declare no competing interests.

Additional information

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