

# SCIENTIFIC REPORTS



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## Influence of tensile-strain-induced oxygen deficiency on metal-insulator transitions in NdNiO<sub>3-δ</sub> epitaxial thin films

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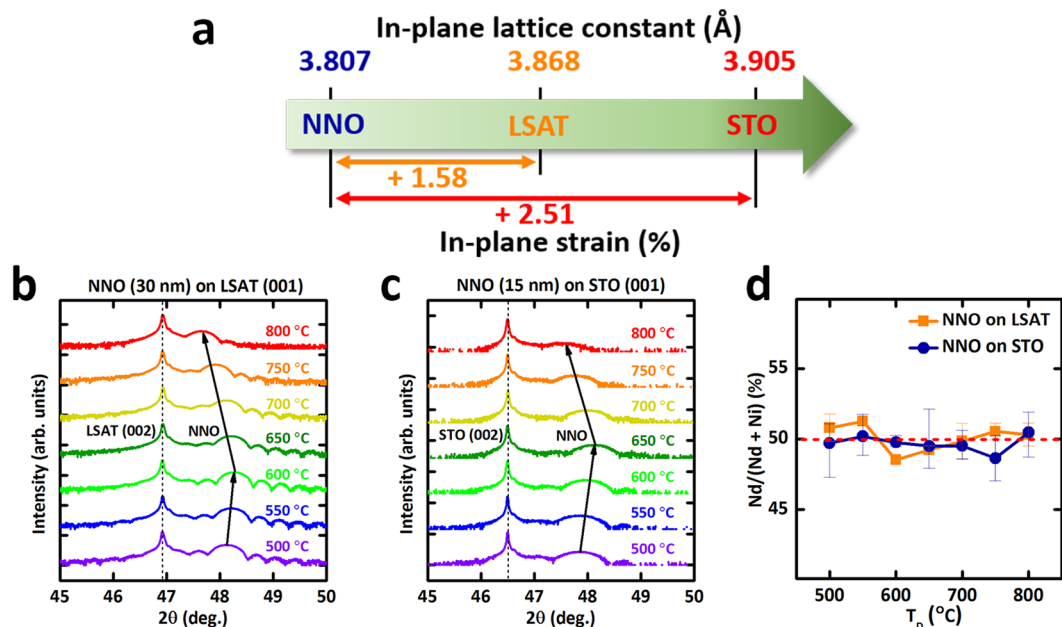
We report direct evidence that oxygen vacancies affect the structural and electrical parameters in tensile-strained NdNiO<sub>3-δ</sub> epitaxial thin films by elaborately adjusting the amount of oxygen deficiency ( $\delta$ ) with changing growth temperature  $T_D$ . The modulation in tensile strain and  $T_D$  tended to increase oxygen deficiency ( $\delta$ ) in NdNiO<sub>3-δ</sub> thin films; this process relieves tensile strain of the thin film by oxygen vacancy incorporation. The oxygen deficiency is directly correlated with unit-cell volume and the metal-insulator transition temperature ( $T_{MI}$ ), i.e., resulting in the increase of both unit-cell volume and metal-insulator transition temperature as oxygen vacancies are incorporated. Our study suggests that the intrinsic defect sensitively influences both structural and electronic properties, and provides useful knobs for tailoring correlation-induced properties in complex oxides.

Correlated transition metal oxides with partially-filled  $d$  electrons undergo intriguing metal-insulator electronic phase transition (MIT)<sup>1,2</sup> as a result of electron interactions<sup>3</sup>. The capability to control those exotic phenomena with external stimuli is a major subject in correlated oxide heterostructure research due to their extreme sensitivity of those materials systems near metal-insulator phase boundary. For example, in these materials, strain and doping can cause small changes in the crystal structures or charge density, and these changes can shift the phase boundaries, and thereby lead to large change in the electrical and optical properties<sup>4-6</sup>. Among the diverse stimuli, oxygen vacancies are inevitable defects in perovskite oxides (ABO<sub>3-δ</sub>, where  $\delta$  is the oxygen deficiency) during the synthesis of oxide thin films. The quantity of oxygen vacancies can be stabilized in the range of  $\delta$ , and thus have a strong influence on the crystal and electronic structures. Because oxygen vacancies have a strong influence on the  $d$ -band filling, as well as structural changes, the ability to control the degree of oxygen deficiency may provide the opportunity to tailor the electrical<sup>7</sup>, optical<sup>8</sup> or electrochemical device functionalities<sup>9-11</sup> of ABO<sub>3-δ</sub>.

Bulk rare-earth nickelates (RNiO<sub>3</sub>, where  $R$  is a trivalent rare-earth ion) have strongly-correlated unpaired electrons in Ni<sup>3+</sup>  $e_g$  band, and therefore undergo a first-order phase transition to an insulating state at 200 K due to strongly correlated unpaired electrons in Ni<sup>3+</sup>  $e_g$  band under cooling<sup>12</sup>. Because nickel ions tend to be stabilized as Ni<sup>2+</sup> valence state, instead of Ni<sup>3+</sup>, bulk RNiO<sub>3</sub> is prone to oxygen deficiency unless it is synthesized under extremely high partial pressure  $p(O_2)$  of oxygen<sup>7,13</sup>. For this reason, the MIT and optical absorption coefficient of RNiO<sub>3</sub> are both sensitively dependent on  $\delta$ <sup>6,7,13-15</sup>.

Likewise, RNiO<sub>3</sub> epitaxial films synthesized using vacuum growth techniques (e.g., pulsed laser deposition<sup>16</sup>) are also prone to the oxygen deficiency due to the low  $p(O_2)$  during the process. RNiO<sub>3</sub> thin films show inconsistent MIT characteristics, and can have different  $T_{MI}$  even at the same lattice mismatch<sup>6,17-20</sup>. For example, NdNiO<sub>3</sub> (NNO) films grown on LaAlO<sub>3</sub> substrates have  $T_{MI}$  ranging from 175 K to 0 K<sup>6,18,20,21</sup>, which indicates the difficulty of synthesizing RNiO<sub>3</sub> that has the desired cation ratio ( $R/Ni$ ) and the degree of oxygen deficiency ( $\delta$ ). In the previous studies, oxygen deficiency in RNiO<sub>3</sub> epitaxial films has been controlled by varying  $p(O_2)$  during their growth<sup>13,15,20</sup>. However,  $p(O_2)$  can significantly affect the cation stoichiometry of RNiO<sub>3</sub> by increasing the scattering of ablated species in the background gas during pulsed laser deposition; the ablated elemental species experience different scattering conditions in different oxygen pressure as they propagates toward the substrate, influencing film cation stoichiometry, as well as oxygen deficiencies. Consequently, despite several

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**Figure 1.** Tensile-strain-induced oxygen deficiency in NNO epitaxial thin films. **(a)** In-plane lattice constants and lattice mismatch of NNO with LSAT and STO substrates. **(b,c)** XRD  $\theta$ - $2\theta$  scan of NNO epitaxial thin films on (001) LSAT **(b)** and (001) STO **(c)**, as a function of growth temperature ( $T_D$ ). Black vertical lines: (002) diffraction peak of substrates (dashed) and pseudocubic NNO epitaxial thin films (solid). **(d)** Quantitative EDS analysis for cation stoichiometry in NNO thin films grown at various  $T_D$ ; cation atomic ratios of all samples differ by less than the error bar.

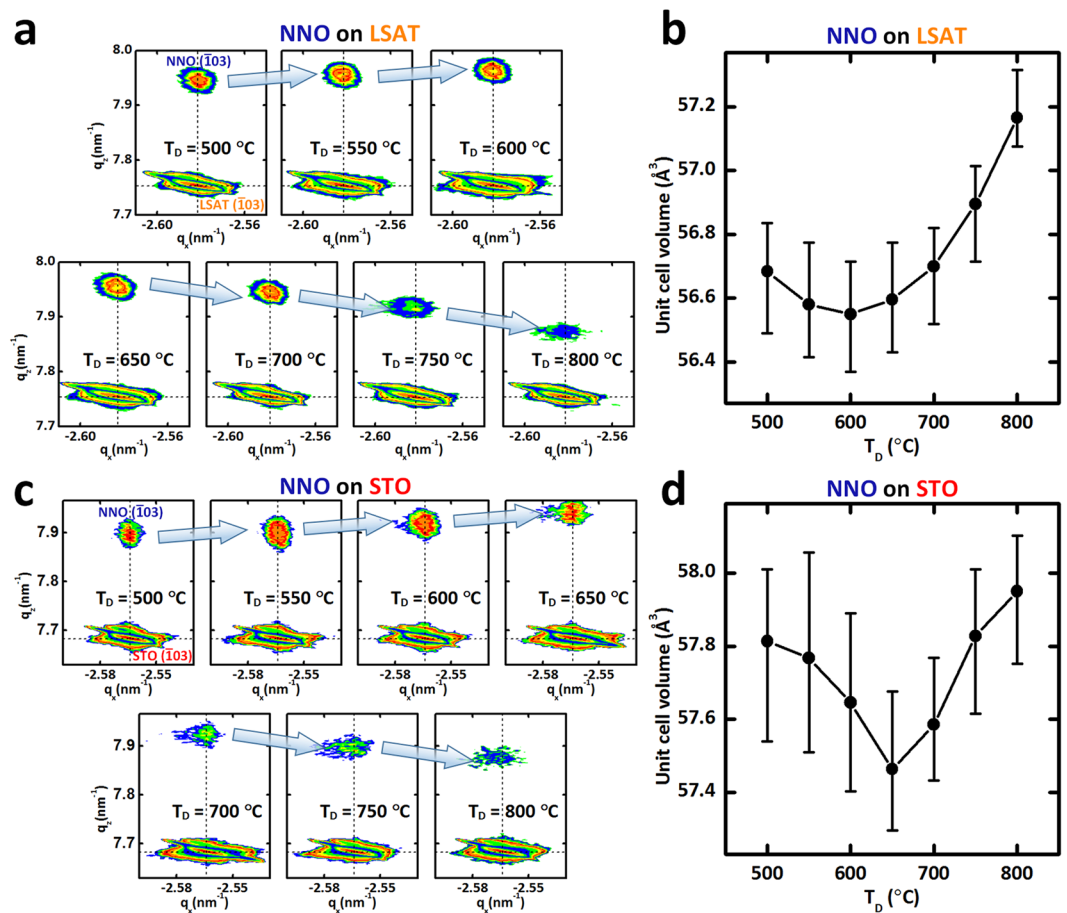
studies that used  $p(O_2)$  to modulate oxygen deficiency in  $RNiO_3$ , to the best of our knowledge, no report has been demonstrated on the tuning of  $T_{MI}$  in  $RNiO_3$  using oxygen deficiency as a single variable while maintaining cation stoichiometry. To investigate the effect of pure oxygen vacancy on  $T_{MI}$  in  $RNiO_3$  thin films, their effect must be isolated from other stimuli such as cation stoichiometry and strain states.

In this Article, we report systematic control of oxygen deficiency in NNO epitaxial thin films with in-plane tensile strain and investigate the influence of oxygen deficiency on the MIT characteristics. In particular, unlike the previous reports, oxygen deficiency was spontaneously generated in tensile-strained NNO thin films, and was finely adjusted by varying growth temperature  $T_D$  without modifying other growth parameters; this result allows tuning the degree of oxygen deficiency while maintaining other stimuli.  $T_{MI}$  was directly correlated with the unit-cell's structural expansion, which is a direct measure of oxygen deficiency. Consequently, these results demonstrate the influence of  $T_D$  on the formation of oxygen vacancies, and also provide insights into the effect of oxygen deficiency on the correlated phase in rare-earth nickelates.

## Results

Epitaxial NNO thin films were grown on (001)-oriented  $(LaAlO_3)_{0.3}-(SrAl_{0.5}Ta_{0.5}O_3)_{0.7}$  (LSAT,  $a = 3.868$  Å) or  $SrTiO_3$  (STO,  $a = 3.905$  Å) single-crystal substrate by pulsed laser deposition at  $500 \leq T_D \leq 800$  °C under  $p(O_2) = 300$  mTorr with a KrF excimer laser fluence of  $2.2$  J/cm<sup>2</sup>. Increase in in-plane tensile strain tends to destabilize the  $Ni^{3+}$  valence state, and as a result, increase oxygen deficiency<sup>22–24</sup>; therefore, in-plane tensile strain was strategically applied in pseudocubic NNO films on both cubic substrates (+1.58% for NNO on LSAT and +2.51% for NNO on STO, Fig. 1a) to systematically generate oxygen deficiency. To ensure that NNO layers were fully-strained on both substrates (Fig. 2a,c), the NNO films used were 30 nm thick on LSAT and 15 nm thick on STO.

$\theta$ - $2\theta$  X-ray diffraction (XRD) symmetrical scan of NNO thin films on (001) LSAT (Fig. 1b) and STO (Fig. 1c) substrates grown at  $500 \leq T_D \leq 800$  °C show a (002) NNO peak near the (002) substrate peak; this observation indicates that the NNO films grew epitaxially in the  $c$ -axis orientation regardless of the growth temperature ( $T_D$ ). In both cases, as  $T_D$  increased from 500 °C, NNO peak was moved to slightly higher scattering angle until to 600 °C for NNO on LSAT, and until 650 °C on STO, then shifted to the lower scattering angle as  $T_D$  increased further. The decrease of the scattering angle represents the expansion of the out-of-plane lattice constant of the NNO thin films. To check the possibility of cation non-stoichiometry as an origin of this expansion, all samples grown at different temperature were characterized by energy dispersive spectroscopy (EDS). The cation atomic ratios of all samples were within 2% of 50% (Fig. 1d); this observation indicates that cation non-stoichiometry is unlikely to be the origin of the structural expansion. Instead, oxygen deficiency can be only modulated by  $T_D$ , because the sticking coefficient and atomic mobility of light-weighted oxygen in the deposited film can be significantly influenced by  $T_D$  while the ablated elemental species (Nd, Ni) experience similar scattering conditions at the same laser fluence and  $p(O_2)$ . The presence of oxygen vacancies would modify the Ni valence from  $Ni^{3+}$



**Figure 2.** Structural modulation of tensile-strained NNO films as a function of oxygen deficiency. (a) Reciprocal space mapping (RSM) around the  $(\bar{1}03)$  LSAT Bragg peaks for 30-nm-thick NNO thin films grown at various  $T_D$ . (b) Unit-cell volume of NNO epitaxial thin films grown at various  $T_D$  on LSAT substrate. (c) RSM around the  $(\bar{1}03)$  STO Bragg peaks for 15-nm-thick NNO thin films grown at various  $T_D$ . (d) Unit-cell volume of NNO epitaxial thin films grown at various  $T_D$  on STO substrate.

( $r = 0.56 \text{ \AA}$ ) to  $\text{Ni}^{2+}$  ( $r = 0.69 \text{ \AA}$ ) by supplying electrons from the missing oxygen; this reduction process expands the crystal lattice by chemical expansion<sup>25–27</sup>.

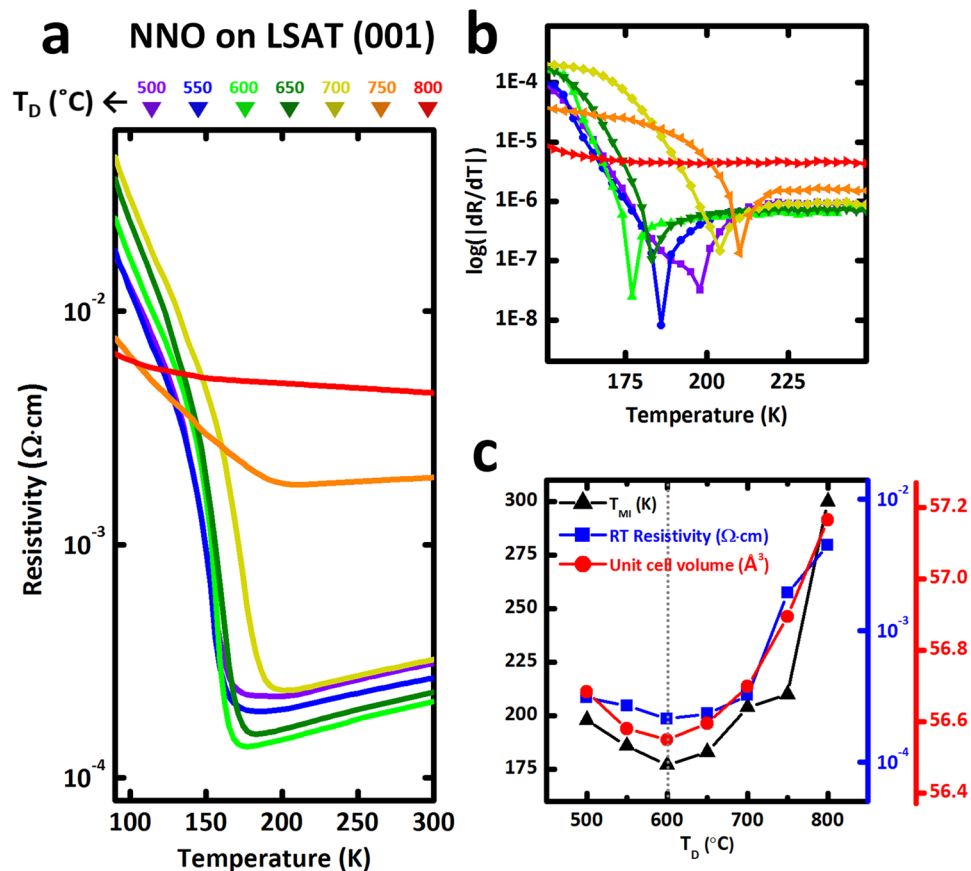
To further verify structural modulation as a function of oxygen deficiency ( $\delta$ ), the volume change of the unit-cell was monitored using reciprocal space mapping (RSM) around the  $(\bar{1}03)$  Bragg reflection of (001)-oriented NNO thin films grown at different  $T_D$  on LSAT (Fig. 2a) and STO (Fig. 2c) substrates.  $q_x$  was the same for all peaks from substrates and NNO films; this consistency means that all NNO thin films are fully-strained on both substrates (Fig. S1). The out-of-plane lattice constants extracted from RSM  $q_z$  values are consistent with those extracted from XRD  $\theta$ - $2\theta$  scans (Fig. S1); this agreement confirms that the lattice expansion caused by oxygen deficiency is accommodated by constraint-free out-of-plane lattice expansion, whereas in-plane axes are tightly constrained by interaction with the substrate in both cases (Fig. 2a,c).

When the films are strained under tensile strain, the purely strain-induced out-of-plane lattice constant  $c_{\text{NNO}}$  of the NNO thin film can be accurately calculated as

$$c_{\text{NNO}} = [(1 + \nu)a_0 - 2\nu a_{\text{NNO}}]/[1 - \nu], \quad (1)$$

where  $\nu = 0.35$  (ref. 28) is Poisson's ratio,  $a_0 = 3.806 \text{ \AA}$  is the unstrained lattice constant, and  $a_{\text{NNO}}$  is the in-plane lattice constant of NNO ( $3.868 \text{ \AA}$  for NNO/LSAT (same as LSAT) or  $3.905 \text{ \AA}$  for NNO/STO (same as STO)). Based on this relationship, estimated  $c_{\text{NNO}}$  under tensile strain were  $3.741 \text{ \AA}$  on NNO/LSAT and  $3.701 \text{ \AA}$  on NNO/STO. However, the smallest measured  $c_{\text{NNO}}$  were  $3.768 \text{ \AA}$  on NNO/LSAT grown at  $600 \text{ }^\circ\text{C}$  (i.e., 0.72% greater than calculated) and  $3.778 \text{ \AA}$  on NNO/STO grown at  $650 \text{ }^\circ\text{C}$  (i.e., 2.08% greater than calculated).

The calculations assume pure strain, so the observation that measured  $c_{\text{NNO}}$  exceeds calculated  $c_{\text{NNO}}$  indicates that, on both substrates, oxygen vacancies are incorporated into even optimized NNO films with the smallest unit-cell volume. Therefore, the experimental value of  $c_{\text{NNO}}$  cannot be simply understood by purely strain-induced Poisson-type contraction. Furthermore, the deviation from the pure strain effect (i.e., difference between measured  $c_{\text{NNO}}$  and calculated  $c_{\text{NNO}}$ ) becomes more pronounced in the case of more tensile-strained NNO films, i.e., NNO on STO. Previous researchers claimed that oxygen vacancies provide a channel that can relieve tensile strain

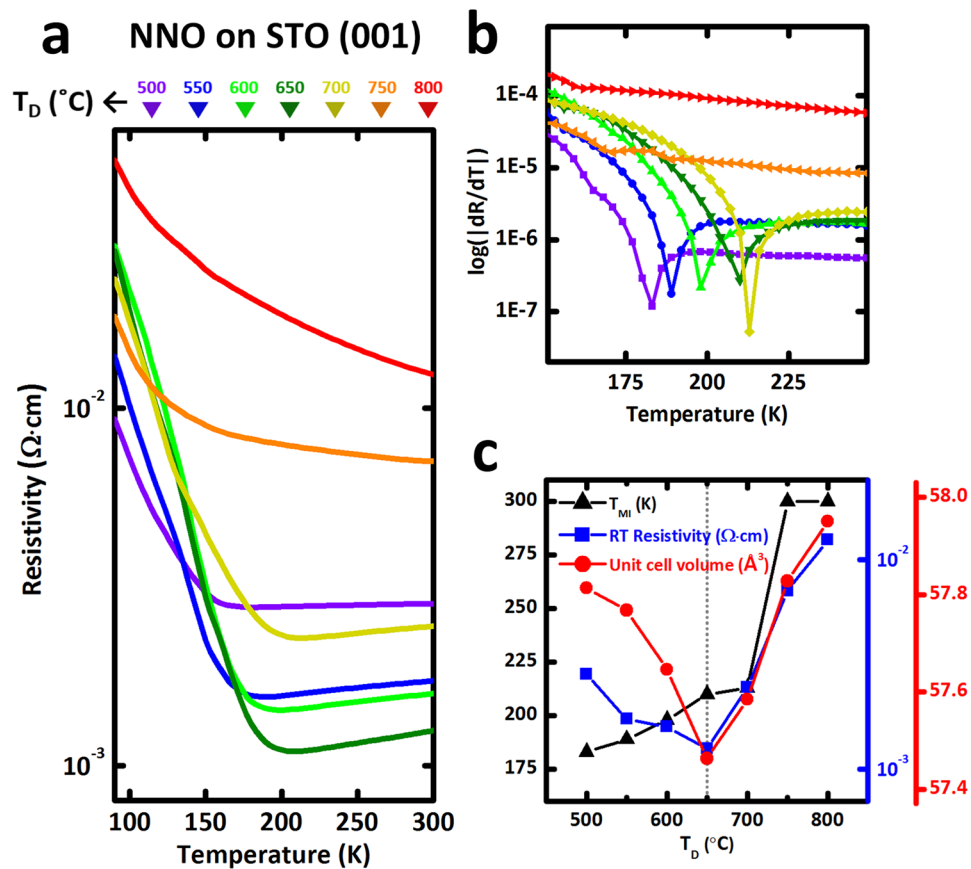


**Figure 3.** Correlation between MIT characteristics and structural modulation in oxygen-deficient  $\text{NdNiO}_{3-\delta}$  films on LSAT substrates. **(a)** Temperature ( $T$ )-dependent in-plane resistivity for the NNO films on LSAT grown at various  $T_D$ . **(b)** The first derivative of resistivity with respect to temperature ( $dR/dT$ ) as a function of growth temperature ( $T_D$ ) to clarify  $T_{MI}$ . **(c)** Correlation between  $T_{MI}$ , room temperature (RT) resistivity and unit-cell volume in oxygen-deficient  $\text{NdNiO}_{3-\delta}$  films. In case of NNO/LSAT grown at 800 °C, temperature dependence of resistivity curves show all insulating behaviors within the measured range (90 to 300 K), which indicates that  $T_{MI}$  is outside the plot window ( $T_{MI} > 300$  K).

in perovskite oxides grown on substrates with large lattice parameters; this claim suggests that the number of oxygen vacancies generated in NNO films to accommodate strain can increase as tensile strain increases<sup>9,22–24,29</sup>. Therefore, the deviation in unit-cell volume from that expected due to the pure strain effect is attributed to the formation of oxygen vacancies, which are increased by tensile strain.

In some perovskite oxides, e.g.,  $\text{CaMnO}_3$ , oxygen deficiency formed at high-temperature growth are unstable over time due to oxygen exchange with the environment, which is facilitated by tensile strain; when left exposed to ambient conditions, the films equilibrate and adopt a very small vacancy concentration without capping layer<sup>30</sup>. However, in our samples, the oxygen vacancies in NNO are not easily diffused out even under ambient conditions: To check time-dependent degree of oxygen deficiency in NNO films, we re-examined the change in the XRD peaks of all  $\text{NdNiO}_{3-\delta}$  samples after 10 months (Fig. S2). Not only did all samples maintain their XRD pattern, but the trend of volcano shapes in out-of-plane scattering angles was also unchanged even after 10 months. Thus, although high degree of oxygen deficiency generated at high temperature are in a metastable state, oxygen exchange with atmosphere appeared to be suppressed in our  $\text{NdNiO}_{3-\delta}$  films because of very limited kinetics of oxygen (or oxygen vacancies) at room temperature. Indeed, previous literature reported that the out-of-plane oxygen vacancy migration barrier of NNO is much larger than that of other perovskite oxide system, such as  $\text{CaMnO}_3$ <sup>31</sup>.

To probe how MIT characteristics can be modulated by adjusting  $\delta$  in  $\text{NdNiO}_{3-\delta}$  films, we measured temperature-dependent in-plane electrical resistivity for the NNO thin films on LSAT (Fig. 3a) and on STO (Fig. 4a) with different  $T_D$ . All samples showed a sharp MIT, except grown at  $T_D = 750$  and 800 °C, which shows degraded MIT characteristic. At room temperature (RT), the resistivity was higher in the optimized NNO on STO ( $>10^{-3} \Omega\cdot\text{cm}$  at  $T_D = 650$  °C) than in the optimized NNO on LSAT ( $<10^{-3} \Omega\cdot\text{cm}$  at  $T_D = 600$  °C); this difference indirectly confirms that the number of oxygen vacancies created to accommodate tensile strain in films increases as the amount of applied strain increases<sup>9,22–24,29</sup>.  $T_{MI}$  can be seen clearly in a plot of the first derivative of resistivity with respect to temperature ( $dR/dT$ ) as a function of growth temperature ( $T_D$ ) (Figs 3b and 4b). Interestingly, the  $T_{MI}$  and RT resistivity were quite correlated with unit-cell volume in NNO films grown on LSAT:



**Figure 4.** Correlation between MIT characteristics and structural modulation in oxygen-deficient  $\text{NdNiO}_{3-\delta}$  films on STO substrates. (a) Temperature ( $T$ )-dependent in-plane resistivity for the NNO films on STO grown at various  $T_D$ . (b) The first derivative of resistivity with respect to temperature ( $dR/dT$ ) as a function of growth temperature ( $T_D$ ) to clarify  $T_{MI}$ . (c) Correlation between  $T_{MI}$ , room temperature (RT) resistivity and unit-cell volume in oxygen-deficient  $\text{NdNiO}_{3-\delta}$  films. In case of NNO/STO grown at 750 and 800 °C, temperature dependence of resistivity curves show all insulating behaviors within the measured range (90 to 300 K), which indicates that  $T_{MI}$  is outside the plot window ( $T_{MI} > 300$  K). Unlike NNO/LSAT,  $T_{MI}$  variation at  $500 \leq T_D \leq 650$  °C does not correspond to the unit-cell volume change in NNO/STO.

both quantities were lowest at the minimum of unit-cell volume (Fig. 3c); this result indicates that the electrical properties of this correlated oxide are strongly influenced by the degree of oxygen deficiency ( $\delta$ ). This tendency is also true for NNO/STO samples: the RT resistivity was lowest at the minimum of unit-cell volume, although  $T_{MI}$  variation from  $T_D = 500$  to 650 °C does not correspond to the unit-cell volume change in NNO/STO.

## Discussion

The results of this experiment suggest that oxygen vacancies induced in NNO film increases as  $T_D$  deviates from the optimal  $T_D$  (600 °C for NNO/LSAT, 650 °C for NNO/STO). Oxygen vacancies increase the unit-cell volume of  $\text{RNiO}_3$ <sup>13,20</sup>, and also increase the RT resistivity of the metallic state<sup>6,7,13–15,32</sup>. The degree of oxygen deficiencies in our NNO films was quantitatively estimated by introducing a simple model in order to analyze the relationship between unit-cell volume and oxygen vacancies in Fig. 5a. According to the empirical model, the unit-cell volume  $V$  of perovskite  $\text{NdNiO}_{3-\delta}$  can be expressed as

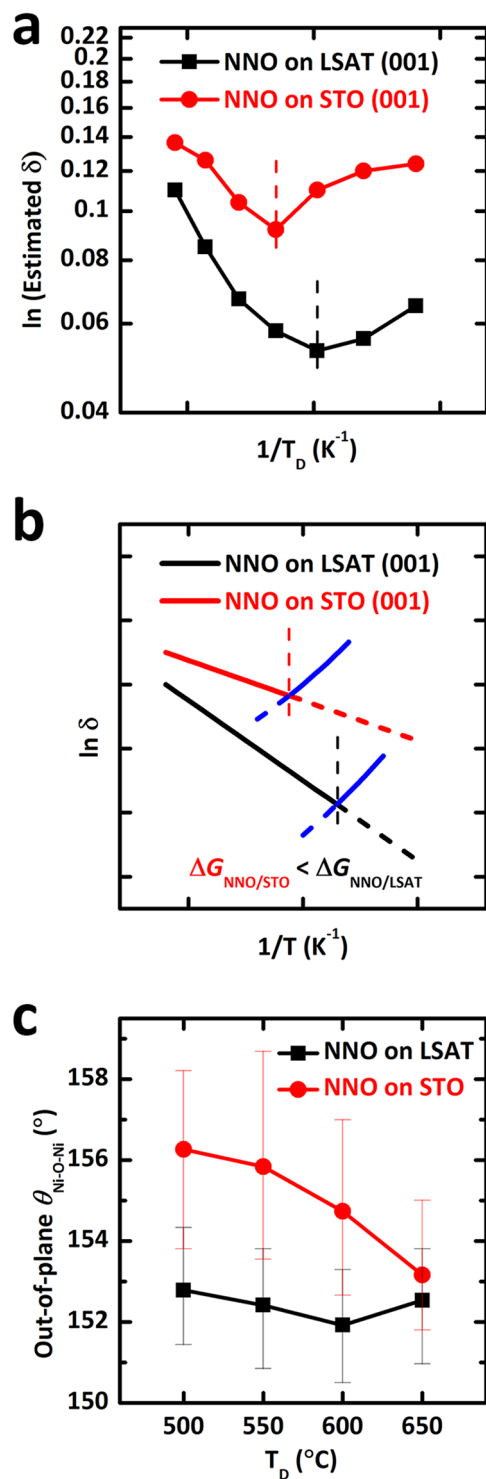
$$V = A^3(r_B + r_{anion})^3 \quad (2)$$

where  $r_B$  and  $r_{anion}$  are the ionic radii of the B-site cation (nickel site) and anion, respectively, and  $A$  is a constant close to 2.

By assuming oxygen-deficient  $\text{NdNiO}_{3-\delta}$  due to the unchanged cation ratio (Nd/Ni) characterized by EDS (Fig. 1d), the effective cation and anion radii could be defined as

$$r_B = (1 - 2\delta)r_{Ni^{3+}} + (2\delta)r_{Ni^{2+}} \quad (3)$$

$$r_{anion} = \left(\frac{3 - \delta}{3}\right)r_O + \left(\frac{\delta}{3}\right)r_V \quad (4)$$



**Figure 5.** Quantitative estimation of oxygen deficiency ( $\delta$ ) and octahedral rotation in oxygen-deficient  $\text{NdNiO}_{3-\delta}$  films. (a) Estimation of the oxygen deficiency ( $\delta$ ) in  $\text{NdNiO}_{3-\delta}$  films using quantitative model. (b) Thermodynamic and kinetic model of inverse temperature ( $1/T$ )-dependent  $\log \delta$  for NNO/LSAT and NNO/STO to explain quantitative estimation of  $\delta$ . (c) Out-of-plane Ni-O-Ni bond angles ( $\theta_{\text{Ni-O-Ni}}$ ) with respect to  $T_D$ . The Ni-O-Ni bond angle is calculated from  $\theta_{\text{Ni-O-Ni}} = 2 \arcsin(c/2d_{\text{Ni-O}})$ , where Ni-O bond length ( $d_{\text{Ni-O}}$ ) was assumed to be equal to that of bulk NNO (1.942 Å). Note that there exists drastic change of Ni-O-Ni bond angle in  $T_D = 500$  to  $650$  °C of NNO/STO.

where  $r_{\text{Ni}^{3+}}$ ,  $r_{\text{Ni}^{2+}}$ ,  $r_{\text{O}}$ , and  $r_{\text{V}}$  are the ionic radii of  $\text{Ni}^{3+}$ ,  $\text{Ni}^{2+}$ , oxygen, and oxygen vacancy, respectively<sup>27, 33</sup>. Consequently, the degree of oxygen deficiency ( $\delta$ ) (Fig. 5a) was roughly estimated to be  $\sim 0.05$  and  $\sim 0.09$  for NNO/LSAT grown at  $600$  °C and NNO/STO grown at  $650$  °C, respectively, from the measured unit-cell volume



(Fig. 2b,d) using this empirical model. This consistent trend on oxygen deficiency in NNO films grown on two different substrates, i.e., the volcano shape as a function of  $T_D$ , indicates that there should be different mechanisms to generate oxygen vacancies at high temperature (600 °C~800 °C for NNO/LSAT and 650 °C~800 °C for NNO/STO) and at low temperature (500 °C~600 °C for NNO/LSAT and 500 °C~650 °C for NNO/STO).

In case of high-temperature regime, a thermodynamic model of oxygen deficiency can be utilized to explain the mechanism based on given experimental evidences. The formation energy of oxygen vacancies ( $\Delta H$ ) decreases with the tensile strain<sup>23</sup> and the degree of oxygen deficiencies ( $\delta$ ) can be expressed as follows:

$$\delta \propto \exp\left(-\frac{\Delta G}{kT}\right) \quad (5)$$

where  $\Delta G = \Delta H - T\Delta S$  ( $\Delta H$  is the enthalpy of oxygen vacancy formation). If we ignore small entropy effects ( $\Delta G = \Delta H$ ), the inverse temperature ( $1/T$ )-dependent  $\log \delta$  for NNO/LSAT and NNO/STO are estimated as shown in Fig. 5b. Moreover, as tensile strain increases, i.e., NNO/STO, the formation energy of oxygen vacancies decreases<sup>23</sup>; this trend can affect oxygen vacancies under various growth temperature ( $T_D$ ), and the change of oxygen deficiency causes systematic variation in  $T_{MI}$  as a function of  $T_D$ . While this thermodynamic model provides a plausible explanation for the increase of oxygen deficiency in the high temperature regime (black and red solid lines in Fig. 5b), it does not account for the increase of oxygen deficiency in the low temperature regime.

In case of low-temperature regime, the insufficient mobility of oxygen in surface migration during the growth might lead to more oxygen-deficient NNO films on both substrates at lower temperature: Since oxygen atoms are unlikely to be supplied sufficiently at low temperature, the amount of oxygen incorporated into the thin film is kinetically limited, resulting in the increase of oxygen deficiency (blue solid lines in Fig. 5b) with decreasing the growth temperature. Considering higher concentration of oxygen vacancies in NNO/STO at low temperature, surface migration of oxygen was more suppressed in NNO/STO than in NNO/LSAT. Indeed, it was previously reported that the oxygen vacancy migration barrier of NNO/STO is larger than that of NNO/LSAT due to the tensile strain-induced suppression of oxygen vacancy migration to the surface<sup>31</sup>.

Despite the close relationship between unit-cell volume and  $T_{MI}$  above, the modulation of  $T_{MI}$  by unit-cell volume was inconsistent at  $500 \leq T_D \leq 650$  °C in highly tensile-strained NNO films on STO substrates. Beyond the explanation using unit-cell volume caused by oxygen vacancies, this can be explained by the influence of the octahedron distortion by tensile strain. In case of LaNiO<sub>3</sub>, in-plane tensile strain weakly modifies the out-of-plane Ni-O distance  $d_{Ni-O}$ , but the out-of-plane Ni-O-Ni bond angle  $\theta_{Ni-O-Ni}$  is highly sensitive to the strain<sup>34</sup>. Since our samples are all fully-strained state (RSM data, see Figs 2 and S1), it is reasonable to assume that in-plane Ni-O-Ni bond angle is fixed. Because  $d_{Ni-O}$  (out-of-plane) is almost insensitive to the strain,  $\theta_{Ni-O-Ni}$  (out-of-plane) can be estimated from the relationship  $\theta_{Ni-O-Ni} = 2\arcsin(c/(2d_{Ni-O}))$  based on an assumption that  $d_{Ni-O}$  (out-of-plane) is equal to that of bulk NNO (1.942 Å). By inserting the observed out-of-plane lattice parameters ( $c$ ), Ni-O-Ni bond angles ( $\theta_{Ni-O-Ni}$ ) was estimated as shown in Fig. 5c. In NNO/STO, Ni-O-Ni bond angles drastically increased up to 156° as  $T_D$  deviated from the optimal  $T_D \sim 650$  °C down to 500 °C; this trend is unlike that of NNO on LSAT. Because Ni-O-Ni bond angle determines the degree of overlap between Ni 3d and O 2p orbitals, orbital overlap increases and, therefore  $T_{MI}$  decreases despite the increase in unit-cell volume. Therefore, in highly tensile-strained NNO films on STO, the effect of octahedron distortion on  $T_{MI}$  is more influential than that of oxygen vacancies on  $T_{MI}$  in the range of  $T_D = 500$  to 650 °C.

## Conclusion

In summary, we report systematic control of metal-insulator transition (MIT) in tensile-strained NdNiO<sub>3- $\delta$</sub>  (NNO) epitaxial thin films by varying the degree of oxygen deficiency  $\delta$  naturally by applying in-plane tensile strain, and by adjusting  $\delta$  at growth temperature  $T_D$  through thermodynamic and/or kinetic control of defect formation. Unit-cell volume and MIT temperature  $T_{MI}$  both increased as  $\delta$  increased. Our study provides useful approaches to finely modulate  $\delta$  in correlated oxides, and provides insights on functional defects in rare-earth nickelate system.

**Experimental Methods.** Epitaxial NdNiO<sub>3</sub> (NNO) thin films were grown on (001)-oriented (LaAlO<sub>3</sub>)<sub>0.3-</sub>(SrAl<sub>0.5</sub>Ta<sub>0.5</sub>O<sub>3</sub>)<sub>0.7</sub> (LSAT) and SrTiO<sub>3</sub> (STO) single-crystal substrates (CrysTec GmbH, Germany) by pulsed laser deposition with a KrF excimer laser ( $\lambda = 248$  nm, Coherent Compex Pro 102 F). A polycrystalline stoichiometric NdNiO<sub>3</sub> ceramic target was used to fabricate the thin film. During deposition, the laser pulse frequency was 10 Hz, laser energy was 2.2 J/cm<sup>2</sup>, and  $p(O_2)$  was 300 mTorr. The growth (= substrate) temperature  $T_D$  was varied from 500 to 800 °C. After each deposition, the sample was cooled to room temperature for 30 min under  $p(O_2) = 300$  mTorr without any *in-situ* post-annealing.

Structural characteristics of the NNO films were measured using  $\theta$ - $2\theta$  scans and reciprocal space mapping (RSM) with a high-resolution X-ray diffractometer (XRD, Bruker D8 Discover X-ray diffractometer) with Cu K $_{\alpha 1}$  radiation ( $\lambda = 1.5406$  Å). Cation stoichiometry of the film was determined in high vacuum by using a field-emission scanning electron microscope (FE-SEM, JEOL JSM-7401F) coupled with an energy dispersive spectroscopy system (EDS, Oxford INCA). In-plane electrical transport was measured using the Van der Pauw method at  $90 \leq T \leq 300$  K during heating.

To check time-dependent change of oxygen deficiency in NNO films, we re-examined the change in the XRD peaks of all NdNiO<sub>3- $\delta$</sub>  samples used for the manuscript after 10 months. All samples were stored in the desiccator at room temperature for 10 months without any capping layer.

## References

1. Yang, Z., Ko, C. & Ramanathan, S. Oxide Electronics Utilizing Ultrafast Metal-Insulator Transitions. *Annu. Rev. Mater. Res* **41**, 337–367 (2011).
2. Ngai, J. H., Walker, F. J. & Ahn, C. H. Correlated Oxide Physics and Electronics. *Annu. Rev. Mater. Res* **44**, 1–17 (2014).
3. Imada, M., Fujimori, A. & Tokura, Y. Metal-insulator transitions. *Rev. Mod. Phys* **70**, 1039–1263 (1998).
4. Scherwitzl, R. *et al.* Electric-Field Control of the Metal-Insulator Transition in Ultrathin NdNiO<sub>3</sub> Films. *Adv. Mater* **22**, 5517–5520 (2010).
5. Liu, M. *et al.* Non-volatile ferroelastic switching of the Verwey transition and resistivity of epitaxial Fe<sub>3</sub>O<sub>4</sub>/PMN-PT (011). *Sci. Rep* **3**, 1876 (2013).
6. Xiang, P.-H. *et al.* Strain controlled metal-insulator transition in epitaxial NdNiO<sub>3</sub> thin films. *J. Appl. Phys* **114**, 243713 (2013).
7. Gayathri, N., Raychaudhuri, A. K., Xu, X. Q., Peng, J. L. & Greene, R. L. Electronic conduction in LaNiO<sub>3-δ</sub>: the dependence on the oxygen stoichiometry  $\delta$ . *J. Phys.-Condens. Mat* **10**, 1323 (1998).
8. Smolin, S. Y. *et al.* Static and Dynamic Optical Properties of La<sub>1-x</sub>Sr<sub>x</sub>FeO<sub>3-δ</sub>: The Effects of A-Site and Oxygen Stoichiometry. *Chem. Mater* **28**, 97–105 (2016).
9. Kubicek, M. *et al.* Tensile Lattice Strain Accelerates Oxygen Surface Exchange and Diffusion in La<sub>1-x</sub>Sr<sub>x</sub>CoO<sub>3-δ</sub> Thin Films. *ACS Nano* **7**, 3276–3286 (2013).
10. Petrie, J. R. *et al.* Enhanced Bifunctional Oxygen Catalysis in Strained LaNiO<sub>3</sub> Perovskites. *J. Am. Chem. Soc* **138**, 2488–2491 (2016).
11. Tsvetkov, N., Lu, Q., Sun, L., Crumlin, E. J. & Yildiz, B. Improved chemical and electrochemical stability of perovskite oxides with less reducible cations at the surface. *Nat. Mater* **15**, 1010–1016 (2016).
12. Torrance, J. B., Lacorre, P., Nazzari, A. I., Ansaldo, E. J. & Niedermayer, C. Systematic study of insulator-metal transitions in perovskites RNiO<sub>3</sub> (R = Pr, Nd, Sm, Eu) due to closing of charge-transfer gap. *Phys. Rev. B* **45**, 8209–8212 (1992).
13. Nikulin, I. V., Novojilov, M. A., Kaul, A. R., Mudretsova, S. N. & Kondrashov, S. V. Oxygen nonstoichiometry of NdNiO<sub>3-δ</sub> and SmNiO<sub>3-δ</sub>. *Mat. Res. Bull* **39**, 775–791 (2004).
14. Moriga, T. *et al.* Characterization of oxygen-deficient phases appearing in reduction of the perovskite-type LaNiO<sub>3</sub> to La<sub>2</sub>Ni<sub>2</sub>O<sub>5</sub>. *Solid State Ionics* **79**, 252–255 (1995).
15. Ha, S. D., Otaki, M., Jaramillo, R., Podpirka, A. & Ramanathan, S. Stable metal-insulator transition in epitaxial SmNiO<sub>3</sub> thin films. *J. Solid State Chem* **190**, 233–237 (2012).
16. Catalan, G., Bowman, R. M. & Gregg, J. M. Metal-insulator transitions in NdNiO<sub>3</sub> thin films. *Phys. Rev. B* **62**, 7892–7900 (2000).
17. Tiwari, A., Jin, C. & Narayan, J. Strain-induced tuning of metal-insulator transition in NdNiO<sub>3</sub>. *Appl. Phys. Lett* **80**, 4039–4041 (2002).
18. Liu, J. *et al.* Strain-mediated metal-insulator transition in epitaxial ultrathin films of NdNiO<sub>3</sub>. *Appl. Phys. Lett* **96**, 233110 (2010).
19. Breckenfeld, E., Chen, Z., Damodaran, A. R. & Martin, L. W. Effects of Nonequilibrium Growth, Nonstoichiometry, and Film Orientation on the Metal-to-Insulator Transition in NdNiO<sub>3</sub> Thin Films. *ACS Appl. Mater. Interfaces* **6**, 22436–22444 (2014).
20. Hauser, A. J. *et al.* Correlation between stoichiometry, strain, and metal-insulator transitions of NdNiO<sub>3</sub> films. *Appl. Phys. Lett* **106**, 092104 (2015).
21. Disa, A. S. *et al.* Phase diagram of compressively strained nickelate thin films. *APL Mater* **1**, 032110 (2013).
22. Kushima, A., Yip, S. & Yildiz, B. Competing strain effects in reactivity of LaCoO<sub>3</sub> with oxygen. *Phys. Rev. B* **82**, 115435 (2010).
23. Aschauer, U., Pfenninger, R., Selbach, S. M., Grande, T. & Spaldin, N. A. Strain-controlled oxygen vacancy formation and ordering in CaMnO<sub>3</sub>. *Phys. Rev. B* **88**, 054111 (2013).
24. Gan, L.-Y., Akande, S. O. & Schwingenschlogl, U. Anisotropic O vacancy formation and diffusion in LaMnO<sub>3</sub>. *J. Mater. Chem. A* **2**, 19733–19737 (2014).
25. Adler, S. B. Chemical Expansivity of Electrochemical Ceramics. *J. Am. Ceram. Soc* **84**, 2117–2119 (2001).
26. Grande, T., Tolchard, J. R. & Selbach, S. M. Anisotropic Thermal and Chemical Expansion in Sr-Substituted LaMnO<sub>3</sub> +  $\delta$ : Implications for Chemical Strain Relaxation. *Chem. Mater* **24**, 338–345 (2012).
27. Marrocchelli, D., Bishop, S. R., Tuller, H. L., Watson, G. W. & Yildiz, B. Charge localization increases chemical expansion in cerium-based oxides. *Phys. Chem. Chem. Phys* **14**, 12070–12074 (2012).
28. Heo, S. *et al.* Modulation of metal-insulator transitions by field-controlled strain in NdNiO<sub>3</sub>/SrTiO<sub>3</sub>/PMN-PT (001) heterostructures. *Sci. Rep* **6**, 22228 (2016).
29. Petrie, J. R., Jeon, H., Barron, S. C., Meyer, T. L. & Lee, H. N. Enhancing Perovskite Electrocatalysis through Strain Tuning of the Oxygen Deficiency. *J. Am. Chem. Soc* **138**, 7252–7255 (2016).
30. Chandrasena, R. U. *et al.* Strain-Engineered Oxygen Vacancies in CaMnO<sub>3</sub> Thin Films. *Nano Lett.* **17**, 794–799 (2017).
31. Mayeshiba, T. & Morgan, D. Strain effects on oxygen migration in perovskites. *Phys. Chem. Chem. Phys* **17**, 2715–2721 (2015).
32. Kaur, D., Jesudasan, J. & Raychaudhuri, P. Pulsed laser deposition of NdNiO<sub>3</sub> thin films. *Solid State Commun* **136**, 369–374 (2005).
33. Shannon, R. D. Revised effective ionic radii and systematic studies of interatomic distances in halides and chalcogenides. *Acta Cryst.* **A** **32**, 751–767 (1976).
34. May, S. J. *et al.* Quantifying octahedral rotations in strained perovskite oxide films. *Phys. Rev. B* **82**, 014110 (2010).

## Acknowledgements

The authors acknowledge support for this work by the Basic Science Research Program through the National Research Foundation of Korea (NRF) funded by the Ministry of Science, ICT & Future Planning (2017R1A2B2007819), and the IT R&D programme of MOTIE/KEIT (10045226). In addition, this study was partially supported by Brain Korea 21 PLUS project (Center for Creative Industrial Materials).

## Author Contributions

S.H. and J.S. conceived the idea. S.H. and C.O. grew the thin films and performed cation stoichiometry measurements. S.H. performed the x-ray diffraction, electrical transport measurements under the supervision of J.S. and H.M.J. S.H. and J.S. wrote the manuscript with the help from all authors.

## Additional Information

**Supplementary information** accompanies this paper at doi:10.1038/s41598-017-04884-2

**Competing Interests:** The authors declare that they have no competing interests.

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