

Table 1. NEUTRON CAPTURE IN THE ISOTOPES OF SAMARIUM

Samarium isotopes at 30 ( $\pm 7$ ) keV							
A	N(Atom%)	Class	$\sigma_c$ (mb)	$N_s\sigma_c$	$N_s$	$N_r$	A
144	2.87	<i>p</i> ( <i>m</i> )	119 $\pm$ 55				
147	14.94	<i>rs</i>	1,173 $\pm$ 192	(2,850)	2.4 $\pm$ 0.4	12.5 $\pm$ 0.4	147
148	11.24	<i>s</i> -0	258 $\pm$ 48	2,930 $\pm$ 540			
149	13.85	<i>rs</i>	1,622 $\pm$ 279	(2,850)	1.8 $\pm$ 0.3	12.1 $\pm$ 0.3	149
150	7.36	<i>s</i> -0	370 $\pm$ 72	2,770 $\pm$ 535			
152	26.00	<i>rs</i>	411 $\pm$ 71	(2,850)	6.9 $\pm$ 1.0	20.0 $\pm$ 1.0	152
154	22.84	<i>r</i> -0	325 $\pm$ 61	0	0	22.84	154

The production process (column 3) is from ref. 1. The *r*-process abundances are derived by assuming  $N_s\sigma_c$  constant for all the isotopes of samarium.

more sensitive indicator is found in the  $N_s\sigma_c(150)/N_s\sigma_c(148)$  ratio. First, the *r*-process is shielded, and secondly the abundance of samarium-150 is only about half that of samarium-149. The samarium-149 depletion is then:

$$\frac{\Delta N(149)}{N(149)} = \frac{7.36}{13.85} [(1.02 \pm 0.06) - 1]$$

which gives an upper limit of 4 per cent for selective depletion, applicable to solar system material as discussed here. It should be mentioned<sup>7</sup> that inclusion of minor secondary processes such as (*p*, $\gamma$ ) and ( $\gamma$ ,*n*) reactions and neutron resonance effects can increase the uncertainties in this limit by perhaps a factor of two.

In summary the following conclusions may be drawn from the samarium results.

(1) The near equivalence of  $N_s\sigma_c$  for samarium-148 and samarium-150 is in excellent agreement with the predictions of Burbidge *et al.*<sup>1</sup>.

(2) At the same time these results provide some constraints for nucleosynthesis models<sup>1,2</sup>.

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<sup>1</sup> Burbidge, E. M., Burbidge, G. R., Fowler, W. A., and Hoyle, F., *Revs. Mod. Phys.*, **29**, 547 (1957).

<sup>2</sup> Fowler, W. A., Greenstein, J. L., and Hoyle, F., *Geophys. J. Roy. Astron. Soc.*, **6**, 148 (1962).

<sup>3</sup> Macklin, R. L., Inada, T., and Gibbons, J. H., *Nature*, **194**, 1272 (1962).

<sup>4</sup> Clayton, D. D. (private communication).

<sup>5</sup> Murthy, V. R., and Schmitt, R. A., *J. Geophys. Res.* (in the press).

<sup>6</sup> Pattenden, N. J. (unpublished results).

<sup>7</sup> Fowler, W. A. (private communication).

## A New ${}^2\Pi-X^2\Pi$ Band System of Nitrogen Sulphide

The spectrum of nitrogen sulphide is excited by means of a micro-wave oscillator of 2,450 Mc/s in sealed quartz discharge tubes containing a few mm of nitrogen and traces of sulphur. Besides the  $\beta$ - and  $\gamma$ -bands, a number of hitherto unassigned bands, degraded to the red, are recorded in the region 2650–2850 Å. These are photographed on a 6.6-m concave grating spectrograph in the second order at a dispersion of 0.57 Å/mm. Rotational structure of seven bands has been analysed and the results are reported in the present communication.

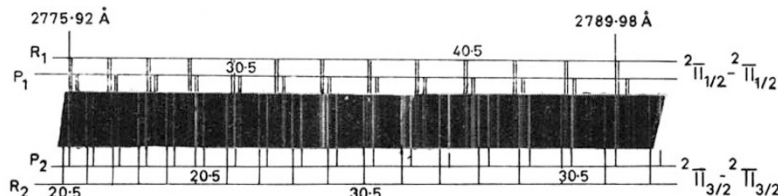


Fig. 1. ( $v' + 1$ ) - 0 band of the new  ${}^2\Pi-X^2\Pi$  system

Each band consists of two sub-bands with an average separation of 107  $\text{cm}^{-1}$ . One such band at 2775 Å is shown in Fig. 1, where the rotational structure of the two sub-bands extends towards the longer wave-lengths as indicated by leading lines. The sub-band shown at the

top in the spectrogram has two *P* and two *R* branches while the other sub-band with leading lines drawn below has only single series of *P* and *R* branches. Such a doubling in the case of *P* and *R* branches of only one sub-band is obtained in the  ${}^2\Pi_{1/2}-{}^2\Pi_{1/2}$  component of a  ${}^2\Pi-{}^2\Pi$  transition with an appreciable  $\Lambda$ -type doubling in the  ${}^2\Pi_{1/2}$  states while in the  ${}^2\Pi_{3/2}-{}^2\Pi_{3/2}$  component the  $\Lambda$ -doubling is negligible.

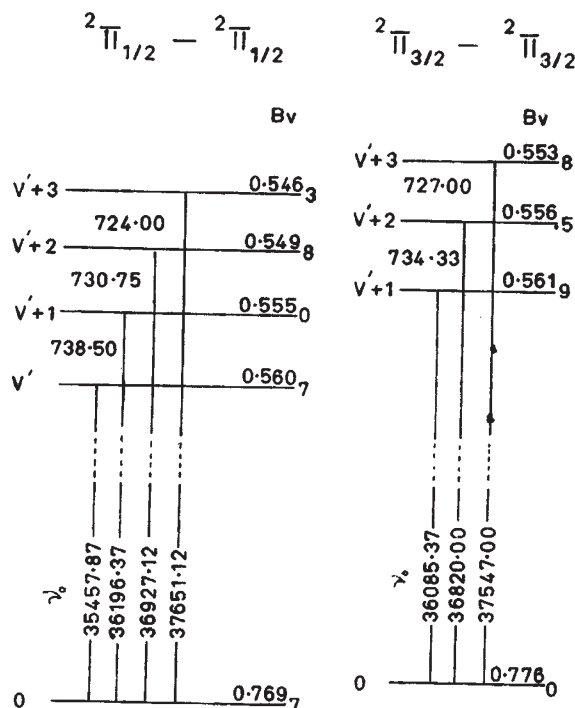


Fig. 2. Energy-level diagram of the vibrational levels of the  ${}^2\Pi-X^2\Pi$  band system

Four bands of  ${}^2\Pi_{1/2}-{}^2\Pi_{1/2}$  and three of  ${}^2\Pi_{3/2}-{}^2\Pi_{3/2}$  transitions have been analysed and their respective vibrational levels are given in an energy-level diagram in Fig. 2. It is found that all the bands arise from different initial vibrational levels to the  $v'' = 0$  level of the final  ${}^2\Pi$  state. The band system belongs to a  ${}^2\Pi(a)-{}^2\Pi(a)$  transition with large spin-splitting in both the states. A comparison of the  $\Delta_2 F''_1(J + 1/2)$  and  $\Delta_2 F''_2(J + 1/2)$  values, that is,  $R(J - 1) - P(J + 1)$ , of the  ${}^2\Pi_{1/2}$  and  ${}^2\Pi_{3/2}$  sub-states of the final  ${}^2\Pi$  level with those obtained for the  $v'' = 0$  of final levels of the  $\beta$ - and  $\gamma$ -systems<sup>1,2</sup> shows good agreement, thereby indicating that all the three band systems have the  ${}^2\Pi$  ground state as the common final level.

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<sup>1</sup> Barrow, R. F., Drummond, G., and Zeeman, P. B., *Proc. Phys. Soc.*, **67**, A, 365 (1954).

<sup>2</sup> Zeeman, P. B., *Canad. J. Phys.*, **29**, 174 (1951).