



THEORETICAL DECAY CURVES FOR DIFFERENT VALUES OF THE INITIAL INTENSITY AND FOR DIFFERENT TEMPERATURES

I' is strongly temperature-dependent and may be called phosphorescence in agreement with a suggestion put forward by Pringsheim⁵. I's is then the phosphorescence intensity extrapolated to t = 0. r decreases during the decay to its asymptotic value R. At very low densities of excitation r_0 may be equal to R. The decay then reduces to one bimolecular curve. With increasing densities of excita-tion I_0 approaches an upper limit corresponding to $l_0 = L$ and $r_0 = 1$. The phosphorescence is then saturated. At very low temperatures $I_0' = 0$. All traps are filled during excitation and the phosphorescence is then frozen in. The decay is bimolecular initially if $L \ll n$ but always changes to an exponential curve $(n - L \ll n)$. At very high temperatures when $\gamma \gg \beta L$, the whole decay is given by

by

$$I = \frac{\beta n_0^2}{(1 + \beta n_0 t)^2} \quad . \quad . \quad . \quad (13)$$

This corresponds to the "Obere Momentanzustand" of Lenard. It can be seen from formulæ (8) and (1) that this occurs if

$$T > \frac{E}{k} \quad \frac{1}{\ln s/\beta L}, \quad \ldots \quad \ldots \quad (14)$$

in our example if $T > 400^{\circ}$ K. At room temperature no phosphorescence will be observed if $E \ll kT$ in $si\beta L$. In our example this would have occurred for trap depths less than 0.35 eV. Fuller details will be published elsewhere. We wish to thank Mr. van Moll and the directors of Philips Lamps, Ltd., for permission to publish this work. H. A. KLASENS

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¹ Randall, J. T., and Wilkins, M. H. F., Proc. Roy. Soc., A., 184, 366

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Diffraction of Light by Ultra-sonic Waves of Very High Frequencies

USING a specially constructed rectifier giving 1,100 volts and a Taylor T 55 valve with amphenol bases and special inductances, frequencies up to 100 megacycles per second have been produced. A tourmaline plate prepared in this laboratory with a thickness of 2 mm. and a fundamental of about 2 Mc./sec. is made to oscillate up to its 54th harmonic, and at all stages it could be employed to maintain stationary waves in a column of water in the usual manner. Diffraction patterns at almost all the frequencies in the range 2 to 100 Mc./sec. could be observed. The highest frequency so far adopted⁴ for such work is only 52.5 Mc./sec. In order to detect dispersion, if

any, of ultra-sonic velocity in water, the crystal has been simultaneously excited by us at two frequencies and both patterns photographed on the same plate at the same instant. Frequency measurement has been effected by beating the oscillator with a standard Philips heterodyne wave-meter, using an audio amplifier for hearing the beat note. Results for two frequencies are given below, the temperature of water at the time of the experiment being 32.6° C.

Order of harmonic	y Frequency in megacycles per second	d Fringe width in cm.	v/d	Velocity calculated in metres per second
5	9·465	0·3655	25.89	$\begin{array}{r}1526\\1524\end{array}$
49	92·28	3·569	25.86	

Results may be taken as indicating that there is no dispersion of ultra-sonic velocity in distilled water. Bär's values when extrapolated to 32.6°C. give a velocity of 1,523 metres per second. S. BHAGAVANTAM

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Symbiotic Aspects of Nitrification

Synbotic Aspects of Nitrification